

Photoproduction of J/ψ in ultra-peripheral collisions

David Grund

Supervisor: Prof. J. G. Contreras

WEJCF 2020, Bílý Potok pod Smrkem

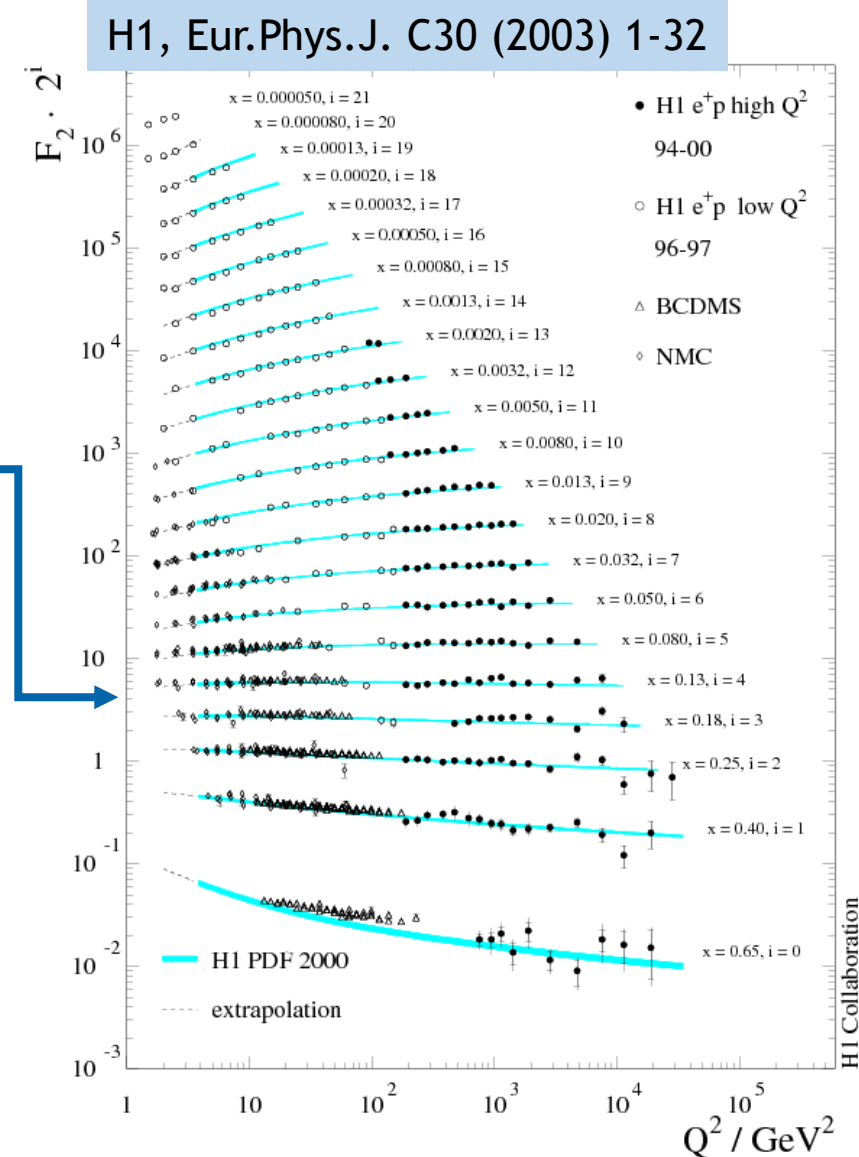
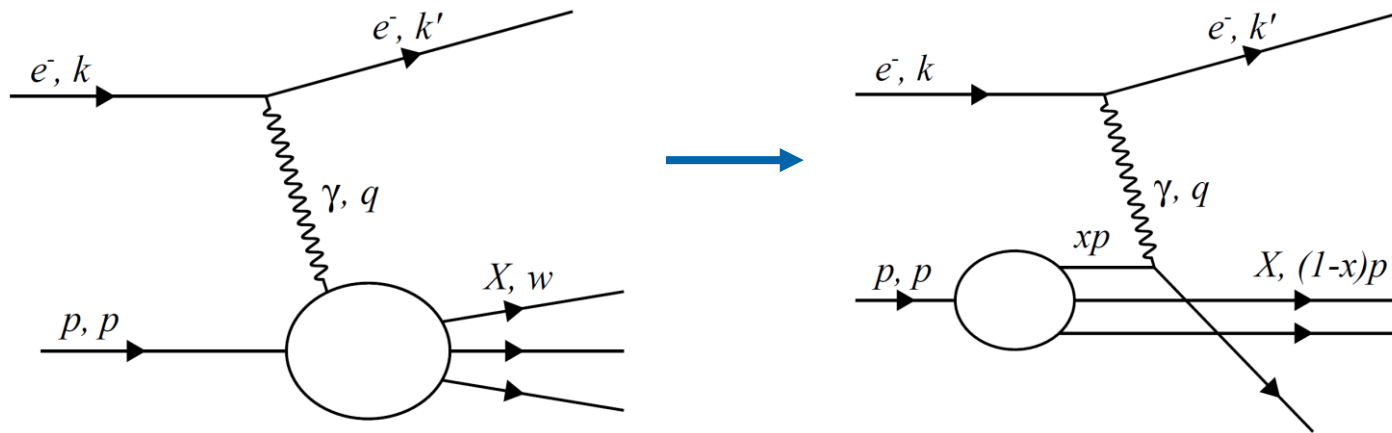
- Historical background
 - Deeply inelastic scattering
 - Structure of hadrons
- Ultra-peripheral collisions and photoproduction of vector mesons
- The cross section of J/ψ photoproduction
- An energy-dependent hot-spot model
 - In p-Pb collisions
 - In Pb-Pb collisions
- Summary and outlook

Deeply inelastic scattering

- The end of 1960s: first experiments
- From the Rosenbluth formula → the DIS cross section:

$$\frac{d^2\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[\left(1 - y - \frac{M_p xy}{s}\right) \frac{F_2(x, Q^2)}{x} + y^2 F_1(x, Q^2) \right]$$

- High- x region ($x \approx 0.1$): Bjorken scaling
- 1969 (Feynman, Bjorken): the parton model



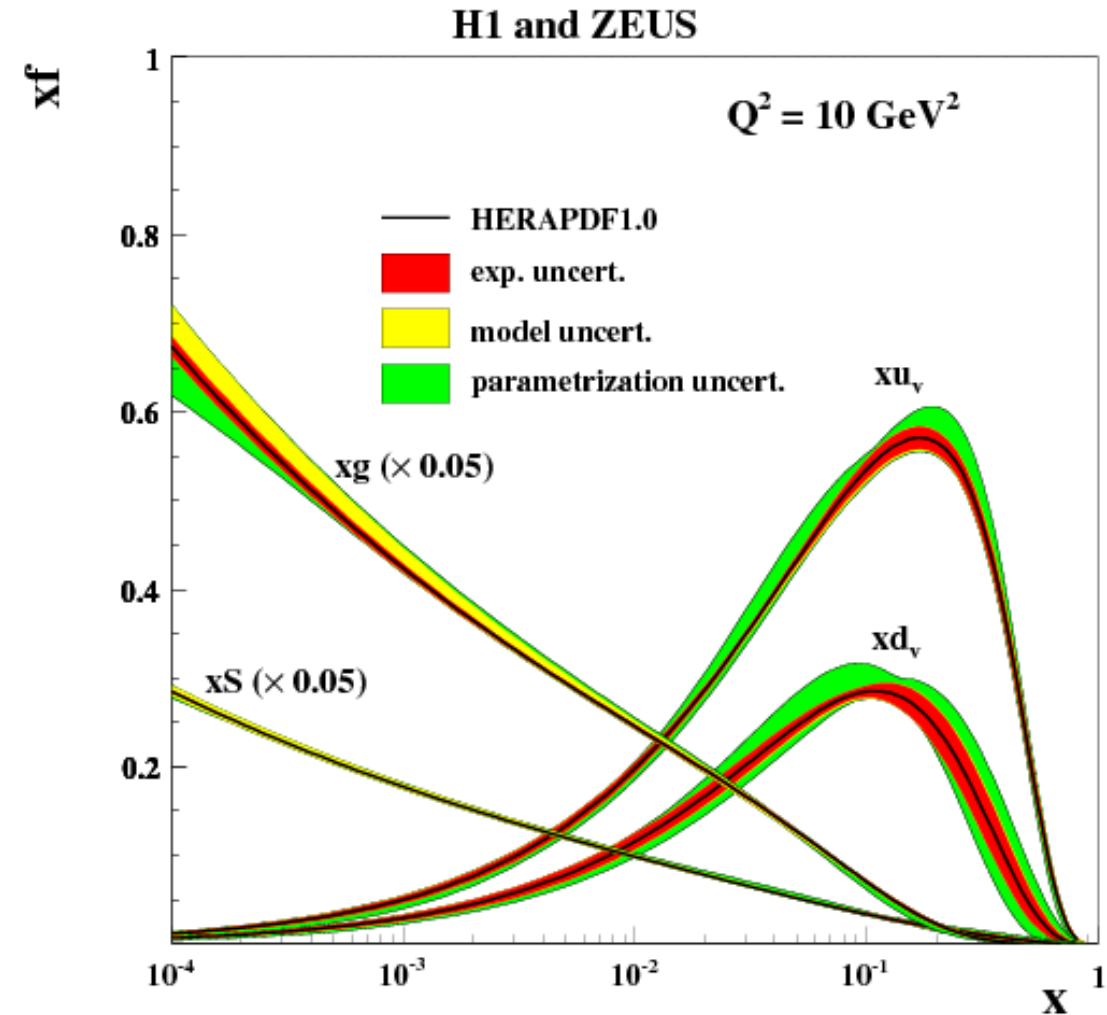
The structure function of a proton

- Callan-Gross relation: $F_2(x) = 2xF_1(x)$
- Parton-quark identification:

$$F_2(x) = x \sum_f e_f^2 [f(x) + \bar{f}(x)]$$

($x = p_{\text{parton}}/p_{\text{hadron}}$)

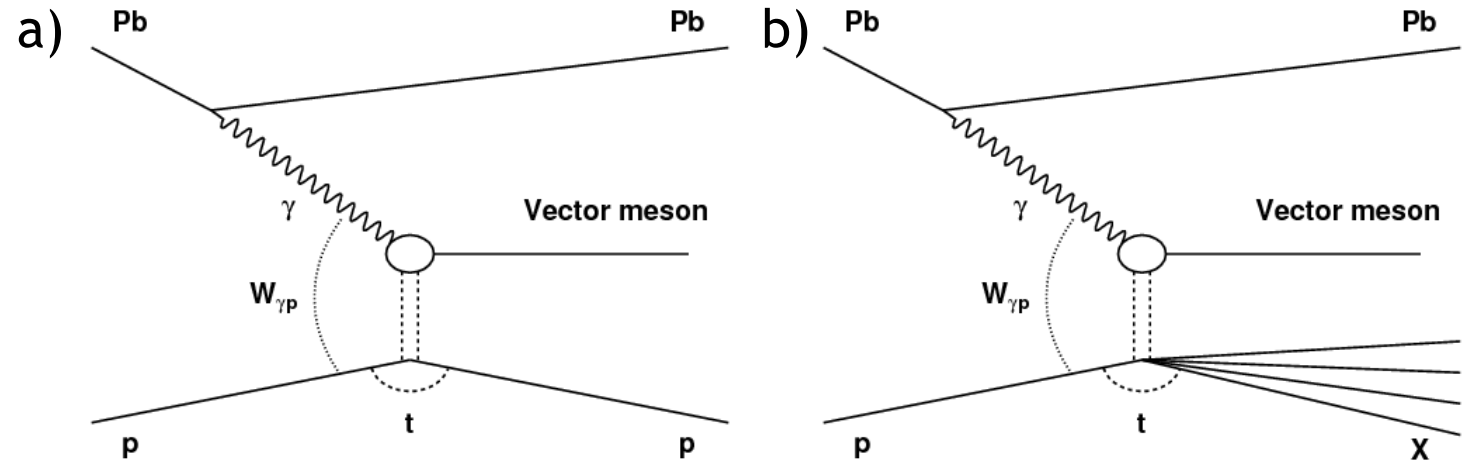
- Valence and sea quarks, gluons: parton distribution functions (PDFs)
- Bjorken scaling only approximate, structure becomes finer at the low- x region
- At even smaller x : **gluon saturation** expected



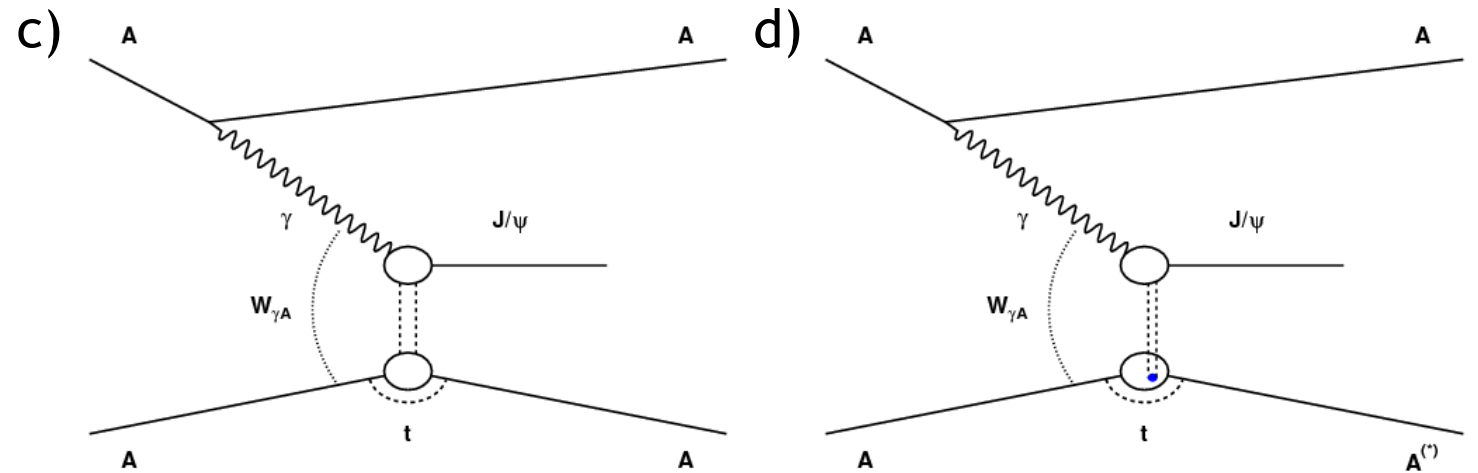
H1 and ZEUS, JHEP 1001 (2010) 109

Ultrapерipheral collisions: vector mesons photoproduction

- UPCs: when $b >$ sum of nuclear/proton radii
- Interaction induced by a photon (QCD short-ranged)
- Creation of a lepton pair vs. vector meson production:
 - p-Pb:
 - a) Exclusive
 - b) Dissociative
 - Pb-Pb:
 - c) Coherent
 - d) Incoherent
- Other possibilities (e.g. 2-jet diffractive production) not yet studied by ALICE



PLBB 766 (2017) 186



PRC 97 (2018), 024901

The J/ψ photoproduction

- **Experimental observables:**

- Rapidity of the $J/\psi \leftrightarrow$ centre-of-mass energy: $W_{\gamma A}^2 = \sqrt{s_{NN}} M_{J/\psi} e^{-y}$
- Transverse momentum of the J/ψ : $p_T^2 = -t$
 - Coherent: $p_T \sim$ a few tens of MeV
 - Incoherent: $p_T \sim$ a few hundred of MeV

F.T.

Distribution of a target matter
in the impact parameter plane
(gluon distribution)

- **The cross section – factorisation:**

- The photon emission (QED process)
- Interaction with a target, QCD enters here (e.g. colour dipole model)

Pb-Pb collision:

$$\frac{d\sigma_{\text{PbPb}}(y)}{dy} = N_{\gamma/\text{Pb}}(y, M)\sigma_{\gamma\text{Pb}}(y) + N_{\gamma/\text{Pb}}(-y, M)\sigma_{\gamma\text{Pb}}(-y),$$

p-Pb collision:

$$\frac{d\sigma_{\text{pPb}}(y, M)}{dy} \approx N_{\gamma/\text{Pb}}(y, M)\sigma_{\gamma\text{p}}(y).$$

Photon flux
 $N_{\gamma} \propto Z^2$

The colour dipole model

- Photon fluctuates into a quark-antiquark pair
- Good-Walker formalism:

γ -p and γ -A cross sections proportional to:

Type of a process	Sensitive to
Exclusive/ coherent	The average over target configurations
Dissociative/ incoherent	The variance over target configurations

$$\begin{aligned} &\longrightarrow \left| \langle A(x, Q^2, \vec{\Delta})_{T,L} \rangle \right|^2 \\ &\longrightarrow \left(\langle |A(x, Q^2, \vec{\Delta})_{T,L}|^2 \rangle - \left| \langle A(x, Q^2, \vec{\Delta})_{T,L} \rangle \right|^2 \right) \end{aligned}$$

• Amplitude:
$$A(x, Q^2, \vec{\Delta})_{T,L} = i \int d\vec{r} \int_0^1 \frac{dz}{4\pi} (\Psi^* \Psi_V)_{T,L} \int d\vec{b} e^{-i(\vec{b} - (1-z)\vec{r}) \cdot \vec{\Delta}} \frac{d\sigma_{\text{dip}}}{d\vec{b}}$$

Transverse momentum of the J/ψ

Distance between the quark and the antiquark

Photon-dipole and vector meson wave functions

The dipole-proton or the dipole-nucleus cross section

Carries information about the **transverse structure of a target!** $\rightarrow T(\vec{b})$

Energy-dependent hot-spot model: γ -p collisions



- A fluctuating proton profile = the sum of high-gluon-density hot spots with gaussian distributions

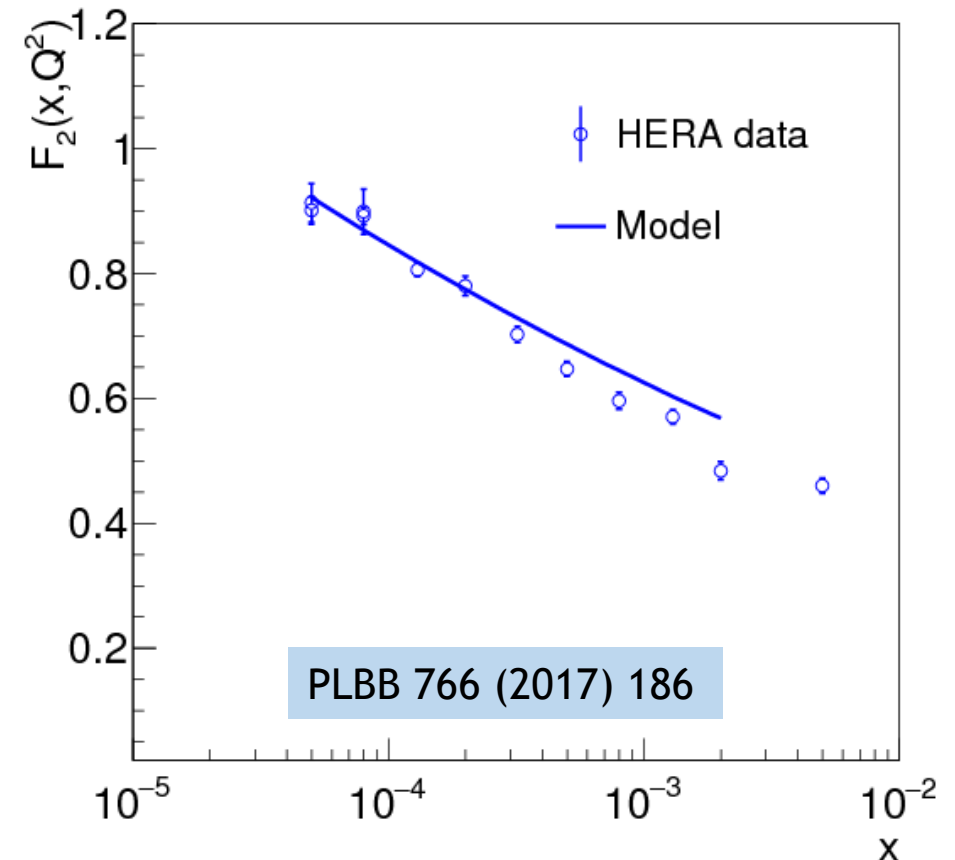
From a 2D gaussian distribution centered at $\vec{b} = (0,0)$

$$T(\vec{b}) = \frac{1}{N_{hs}} \sum_{i=1}^{N_{hs}} T_{hs}(\vec{b} - \vec{b}_i), \quad T_{hs}(\vec{b} - \vec{b}_i) = \frac{1}{2\pi B_{hs}} e^{-\frac{(\vec{b} - \vec{b}_i)^2}{2B_{hs}}}$$

- Number of hot spots grows with energy:

$$N_{hs}(x) = p_0 x^{p_1} (1 + p_2 \sqrt{x})$$

- Data from H1 (HERA, e-p) and ALICE (p-p): energy dependence of the J/ψ cross section
- But first: the model describes also the $F_2(x, Q^2)$ data quite well!



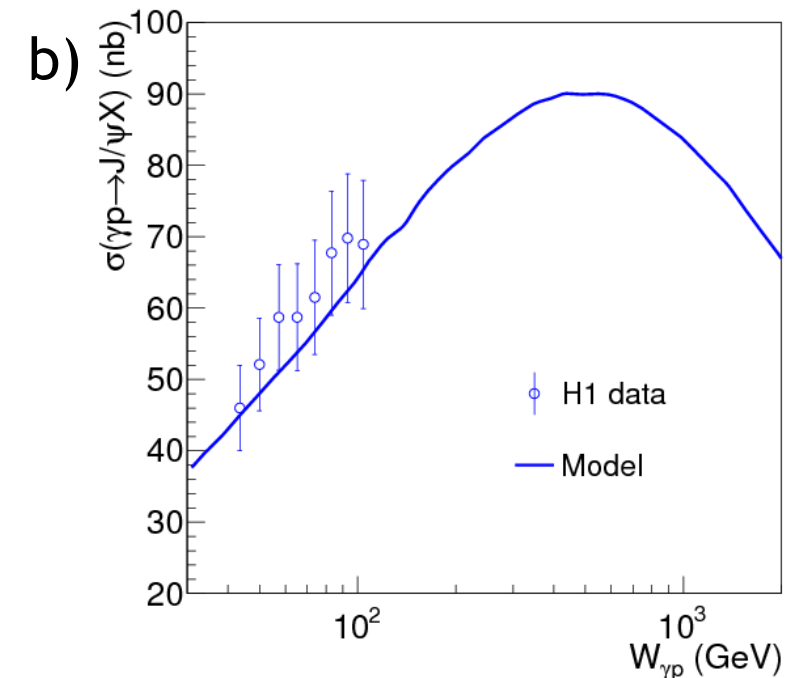
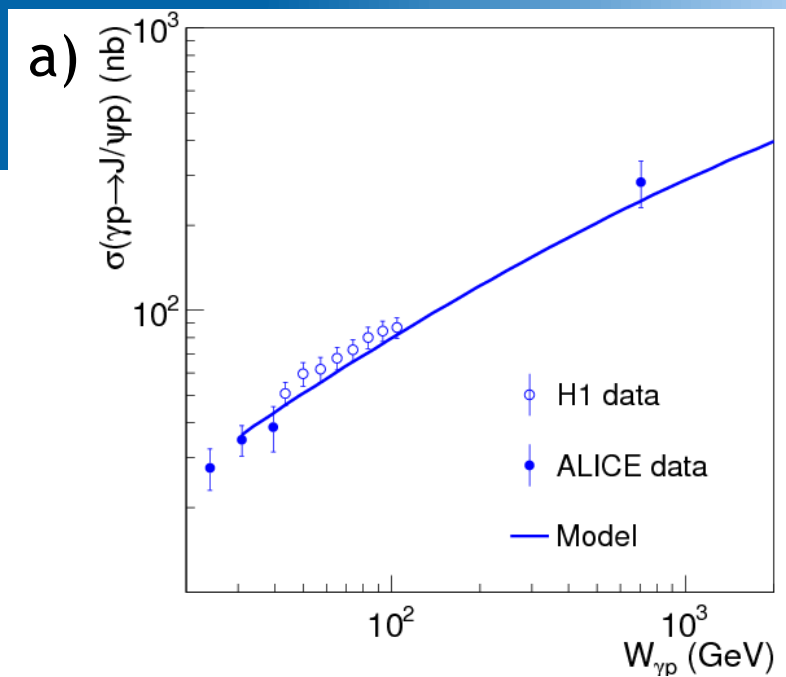
γ -p collisions: The total cross sections



a) The exclusive production

b) The dissociative production:

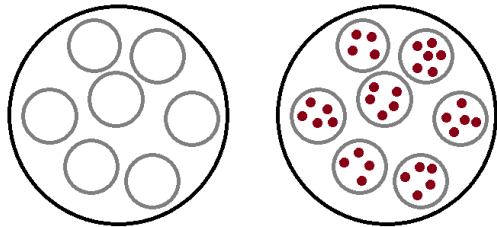
- Prediction of a maximum at $W_{\gamma p} \approx 500$ GeV
 - Ca. 10 hot spots (a sizeable overlap)
 - A decrease in variance
- A steep decrease for $W_{\gamma p} > 500$ GeV
- LHC energies
 - => a way to examine the gluon saturation?
- The t-dependence (not shown): model also in a good agreement



Energy-dependent hot-spot model: Pb-Pb collisions



- Extension to photonuclear interactions:
 - Glauber-Gribov formalism (GG)
 - Geometric scaling (GS)
- Models of the transverse nuclear profile:
 - Nucleons (n)
 - Nucleons made of **hot spots** (hs)
 - Inclusion of subnucleonic degrees of freedom



$$\left(\frac{d\sigma_{dA}}{d\vec{b}}\right)_j = 2 \left[1 - \exp\left(-\frac{1}{2}\sigma_{dp}(x, r)T_A^j(\vec{b})\right) \right]$$

$$\left(\frac{d\sigma_{dA}}{d\vec{b}}\right)_j = \sigma_0^A \left[1 - \exp\left(-r^2 Q_{A,s}^2(x)/4\right) \right] T_A^j(\vec{b})$$

$$T_A^j(\vec{b}) = \frac{1}{2\pi B_p} \sum_{i=1}^A \exp\left(-\frac{(\vec{b} - \vec{b}_i^j)^2}{2B_p}\right)$$

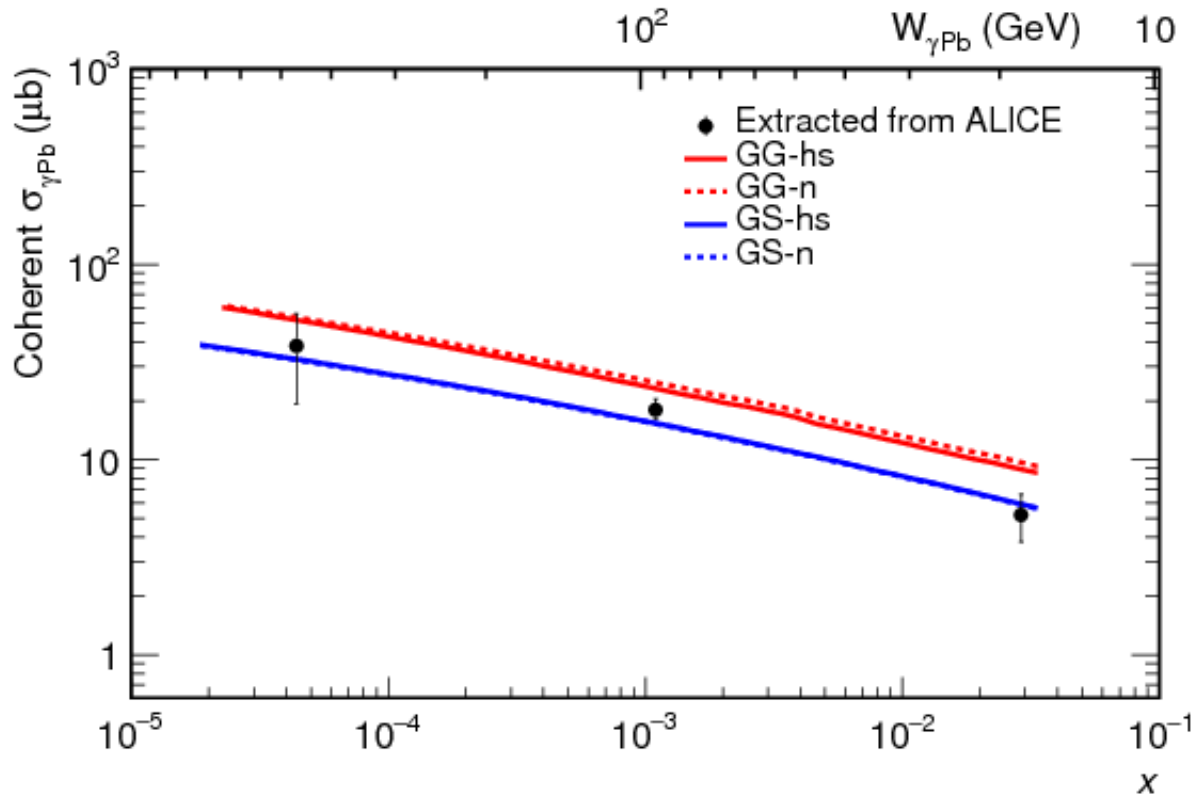
$$T_A^j(\vec{b}) = \frac{1}{2\pi B_{hs}} \sum_{i=1}^A \frac{1}{N_{hs}} \sum_{k=1}^{N_{hs}} \exp\left(-\frac{(\vec{b} - \vec{b}_i^j - \vec{b}_k^j)^2}{2B_{hs}}\right)$$

- Again $\langle N_{hs}(x) \rangle = p_0 x^{p_1} (1 + p_2 \sqrt{x})$ (mean value of the Poisson distribution)
- Comparison with data from RHIC and the LHC Run 1 (ALICE PCs and UPCs)

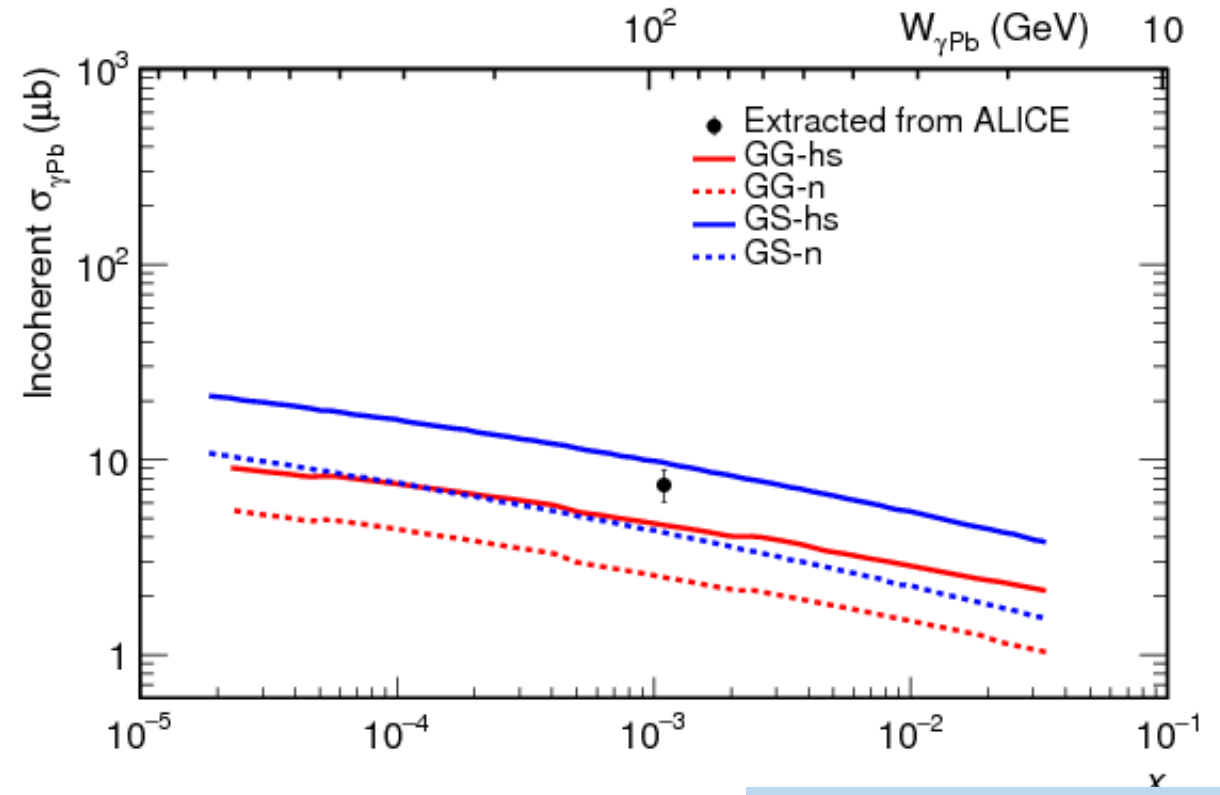
Pb-Pb collisions: (In)coherent cross sections



- The coherent case:
 - Similar predictions



- The incoherent case:
 - Looks like the model prefers hot-spot structure

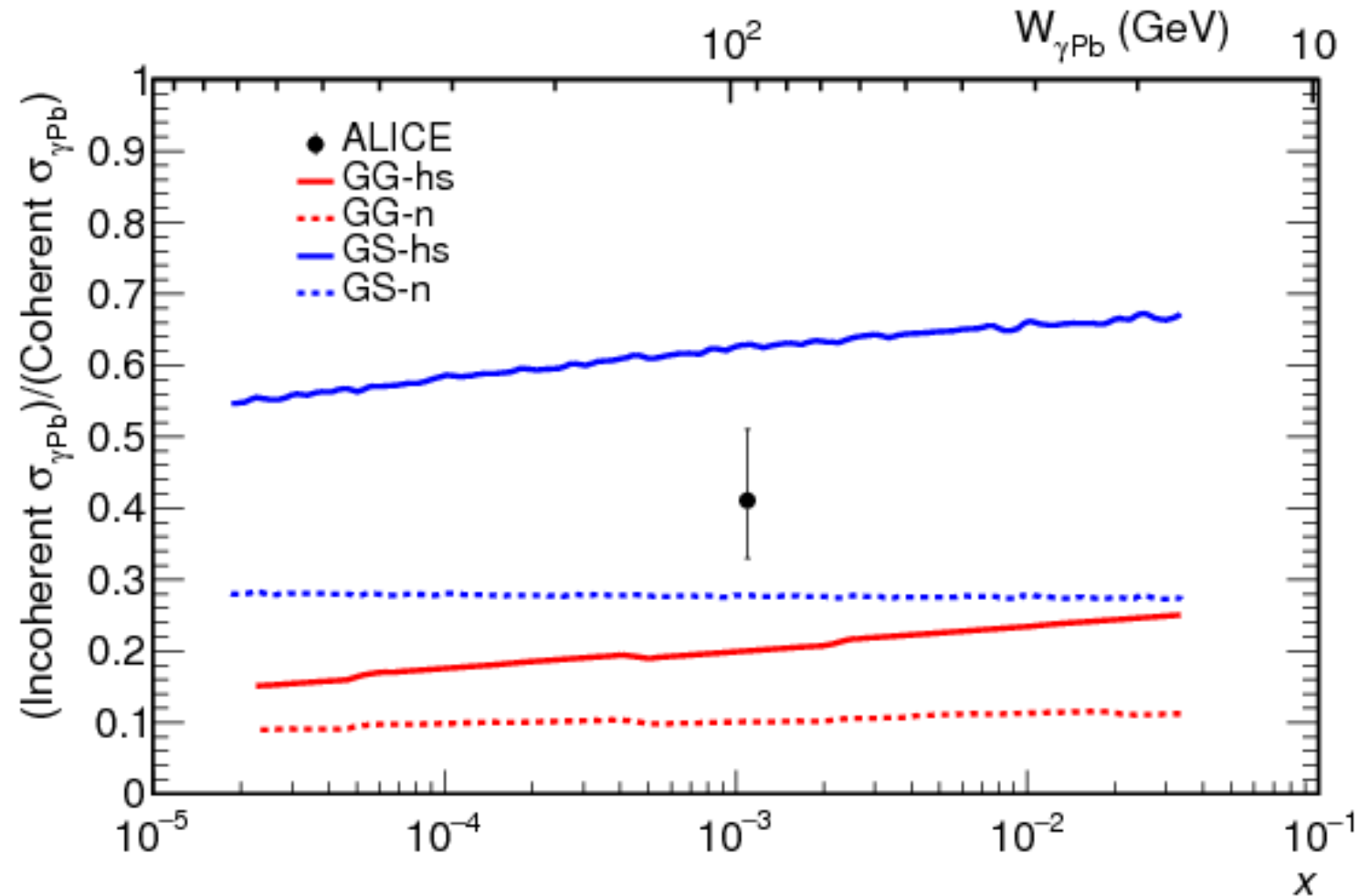


PRC 97 (2018), 024901

Pb-Pb collisions: Ratio of the coherent to the incoherent cross section



- Subnucleonic degrees of freedom introduce the x -dependence
 - Supported by data from PHENIX ($x \approx 0.015$)
- Motivation for new data with smaller uncertainties

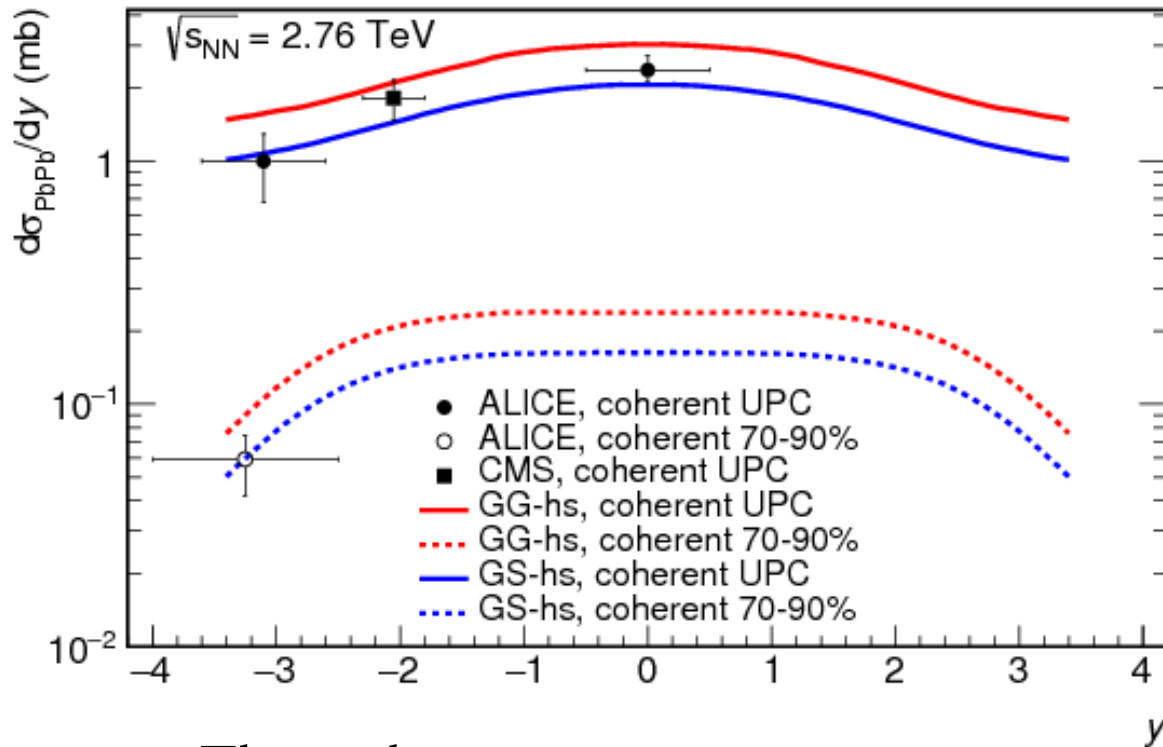


PRC 97 (2018), 024901

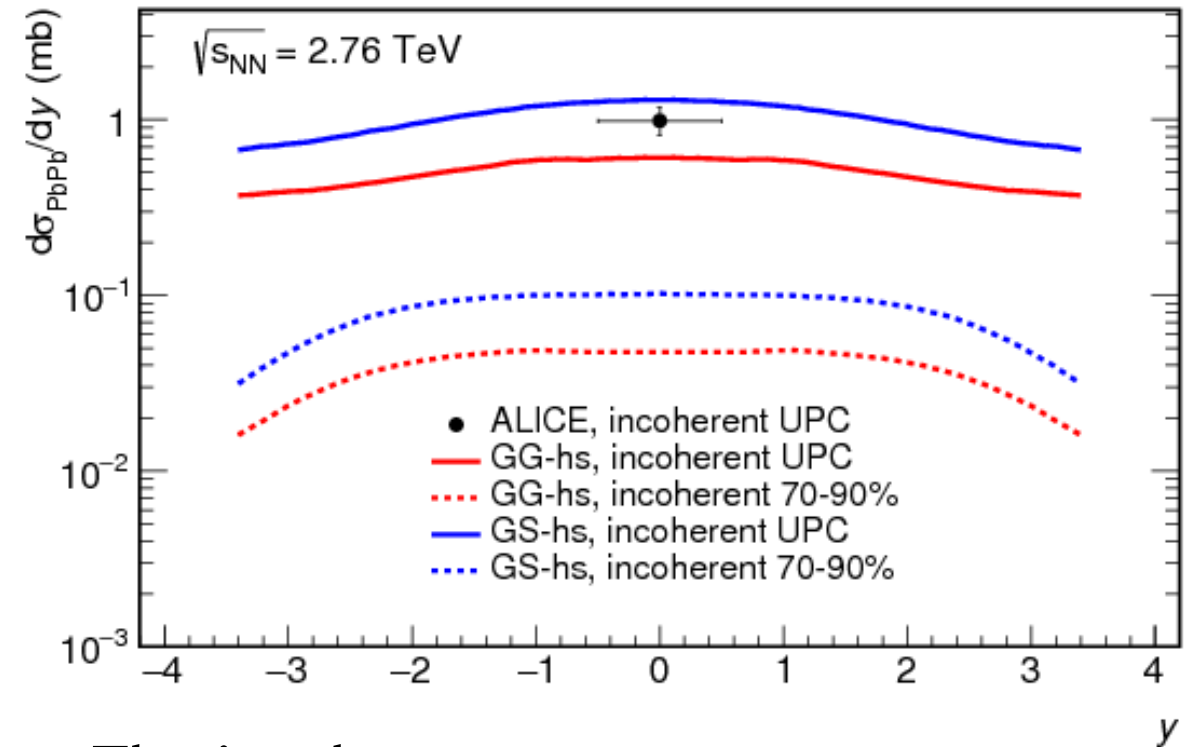
Pb-Pb collisions: The y dependence of the cross section



Note: only hot-spot transverse structure used



- The coherent case
 - GG model slightly overestimates the data



- The incoherent case
 - New data needed

PRC 97 (2018), 024901

- **Ultra-peripheral collisions** = a way to examine the QCD physics in photon-induced processes
 - ALICE: measurements of PDFs in long x -range (ca. from 10^{-2} to 10^{-5})
 - Can provide evidence for gluon saturation?
- **Energy-dependent hot-spot model**
 - Fluctuating proton/nuclear transverse structure
 - Energy-dependent number of hot spots
 - Seems to be in a noteworthy agreement with the data!
 - Has a potential to predict gluon saturation
- My task is to measure the energy dependence of the incoherent cross section using ALICE data

- J.G. Contreras, J. D. Tapia Takaki: Ultra-peripheral heavy-ion collisions at the LHC, International Journal of Modern Physics A 30(8), 1542012 (2015).
- J. Cepila, J. G. Contreras, J. D. Tapia Takaki: Energy dependence of dissociative J/ψ photoproduction as a signature of gluon saturation at the LHC, Phys. Lett. B766, 186-191 (2017).
- J. Cepila, J. G. Contreras, M. Krelina: Coherent and incoherent J/ψ photonuclear production in an energy-dependent hot-spot model, Phys. Rev. C97(2), 024901 (2018).