

Flow

selected topics

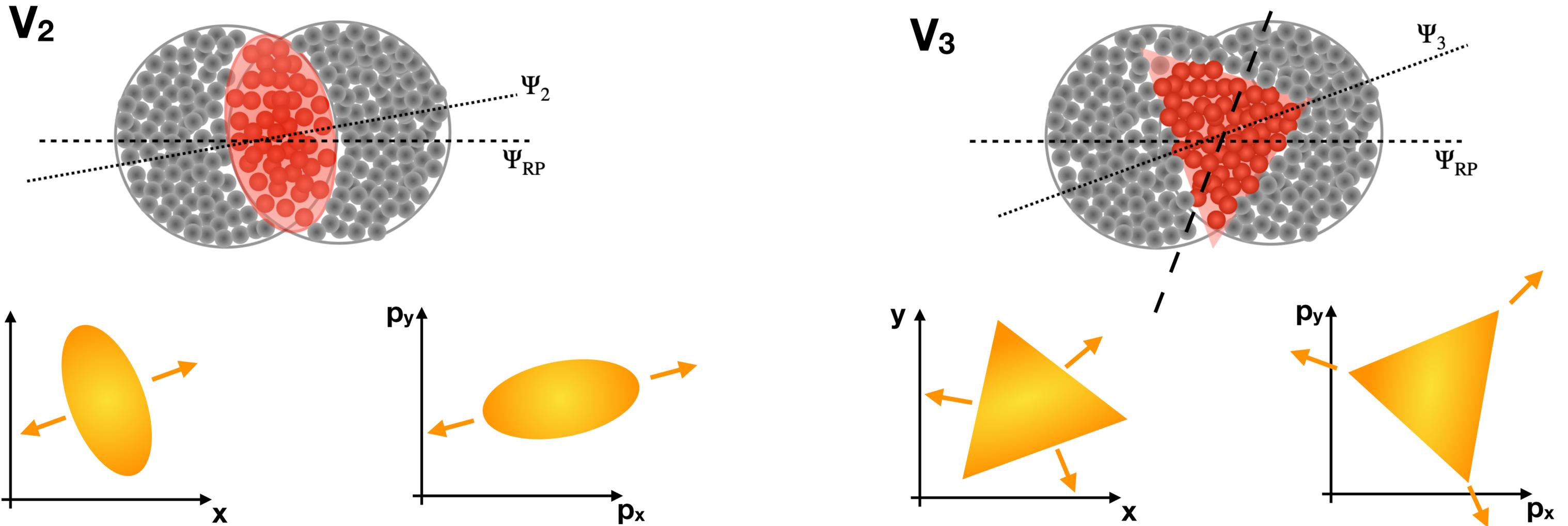
Katarína Křížková Gajdošová
Czech Technical University in Prague

Third workshop on diffraction and ultra-peripheral collisions
16th September 2020



Introduction

flow vector $V_n = v_n e^{in\Psi_n}$
flow magnitude v_n
flow angle: azimuth of V_n in momentum space Ψ_n

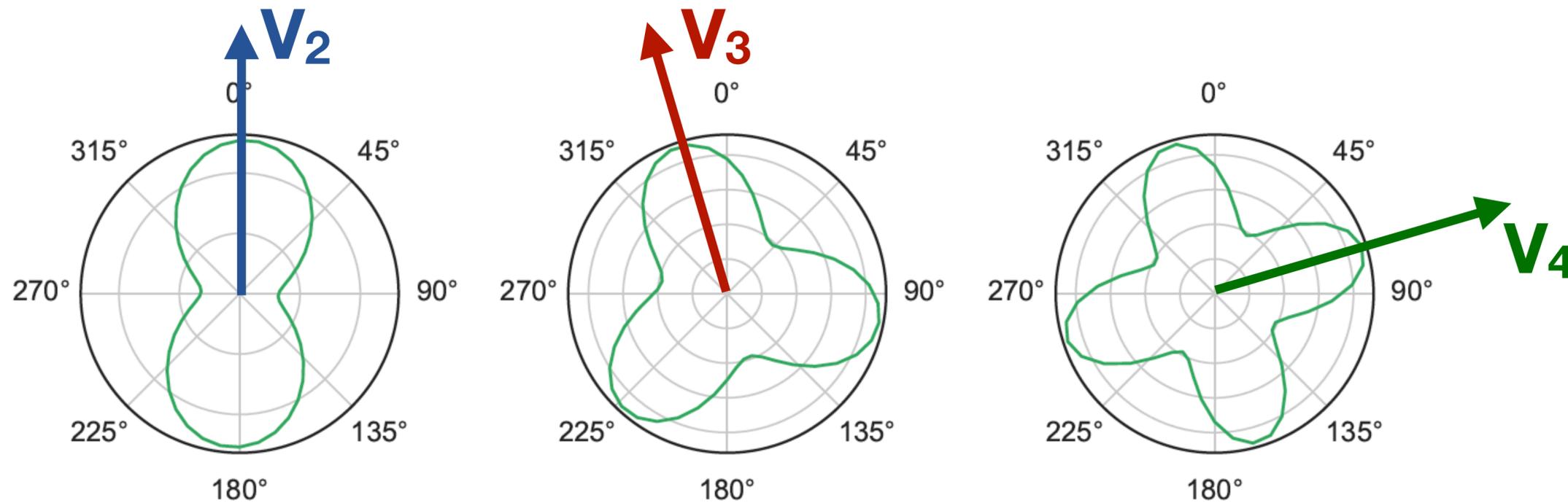


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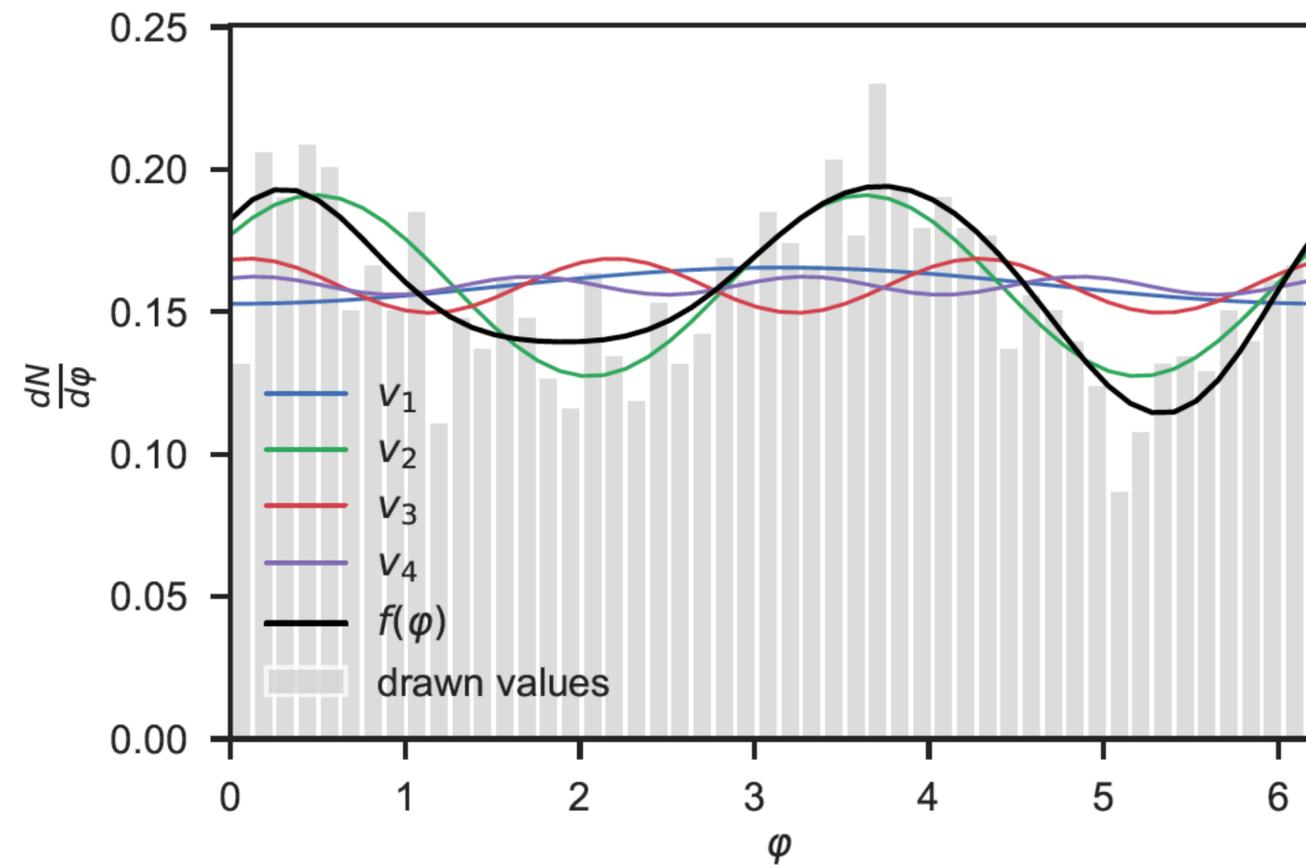
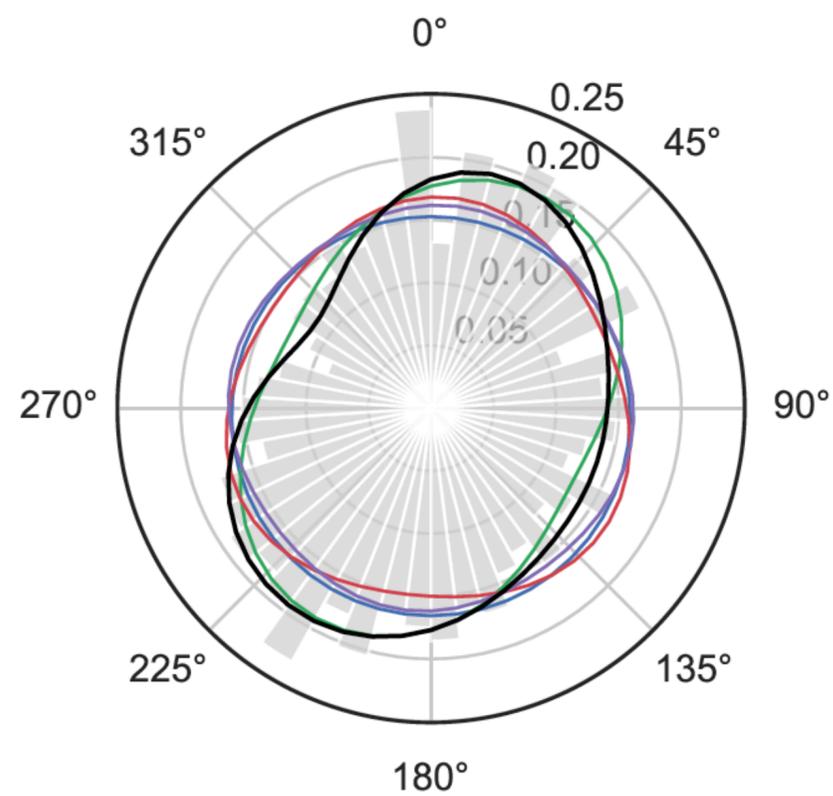


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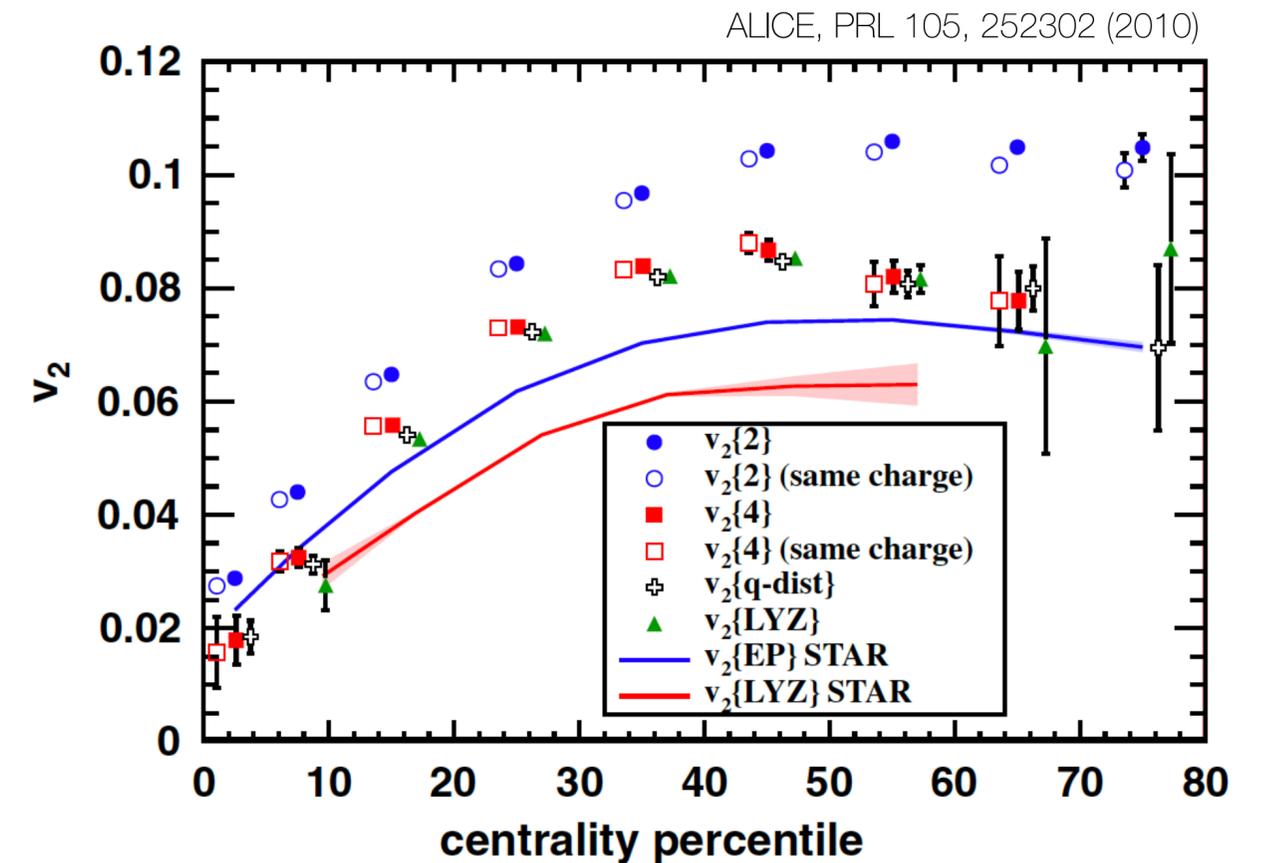
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- We usually measure magnitude of the flow vector = flow coefficient v_n

- As a function of:
 - centrality
 - p_T
 - pseudorapidity η
 - different collision systems



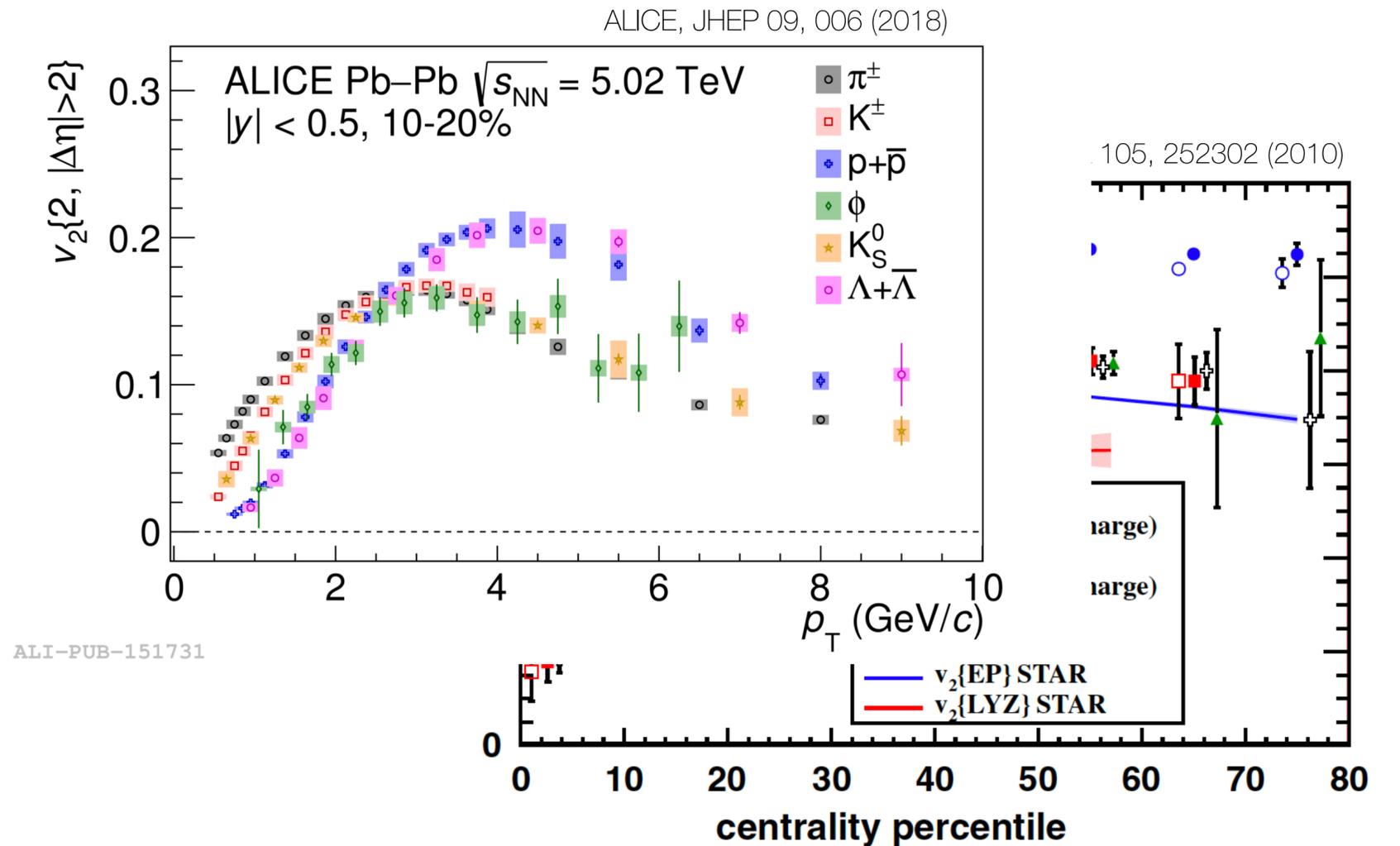
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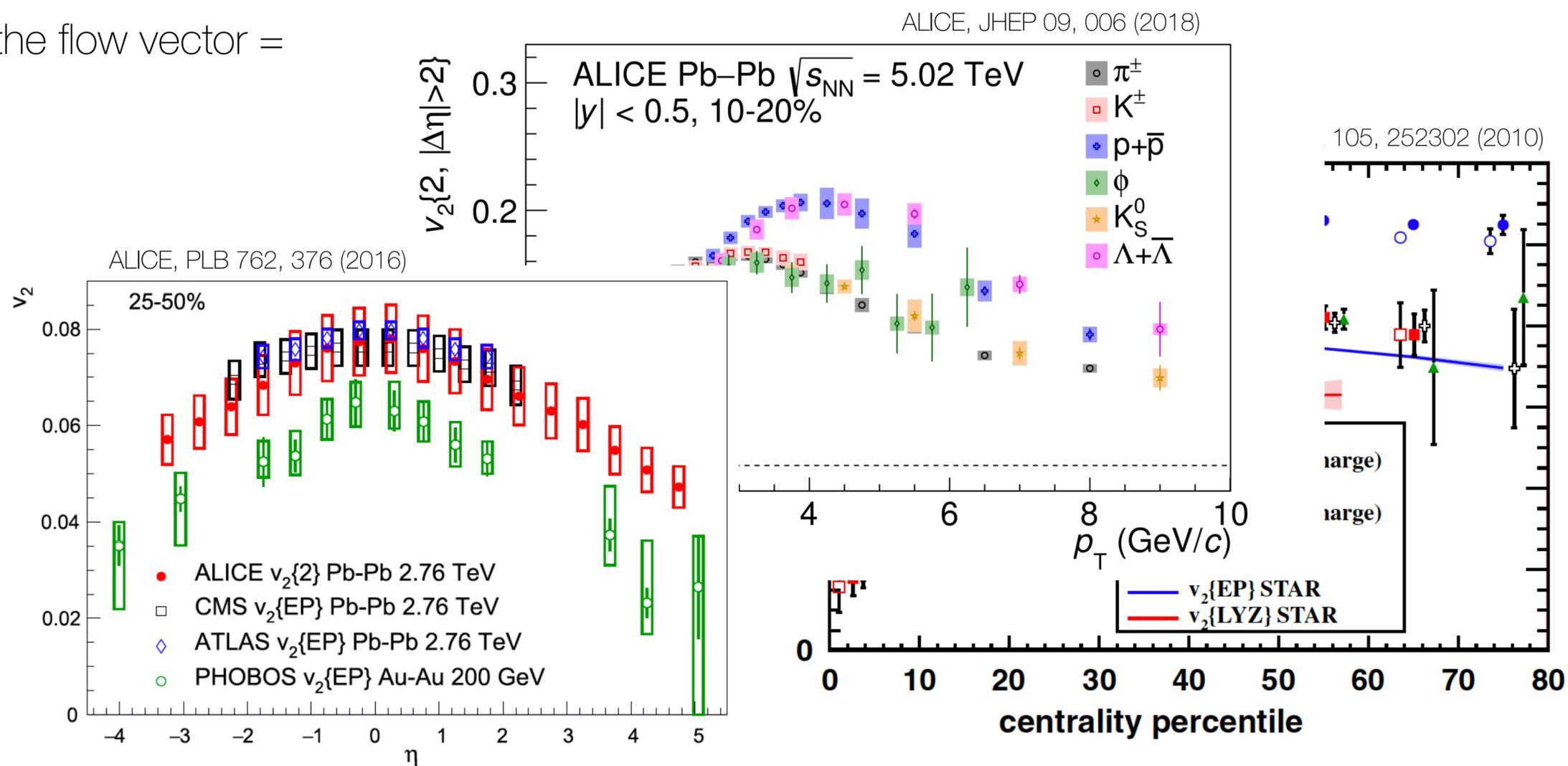
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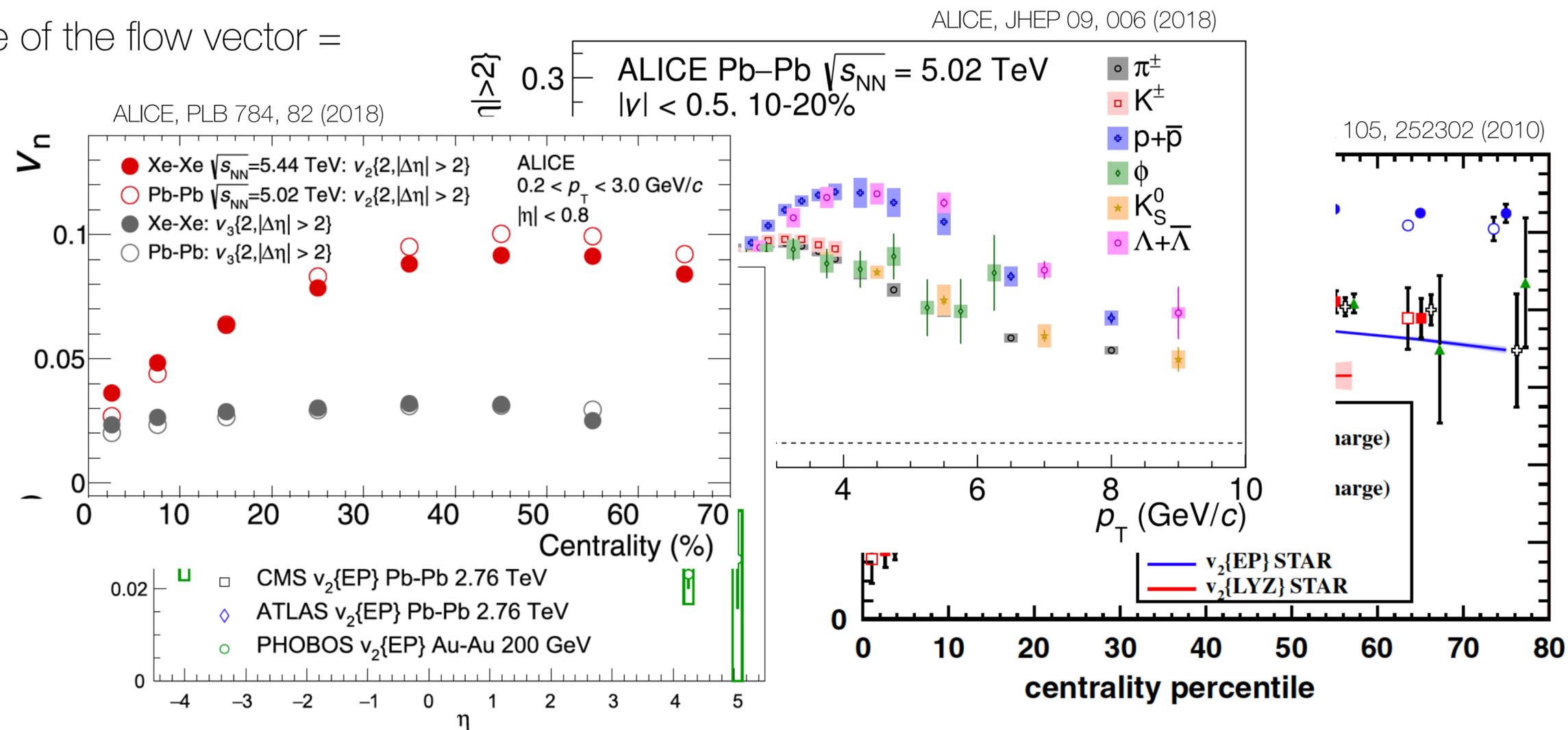
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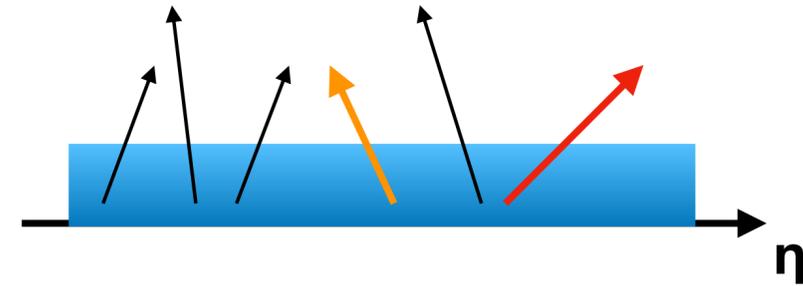
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$$v_n = \langle \cos n(\varphi - \Psi_n) \rangle$$

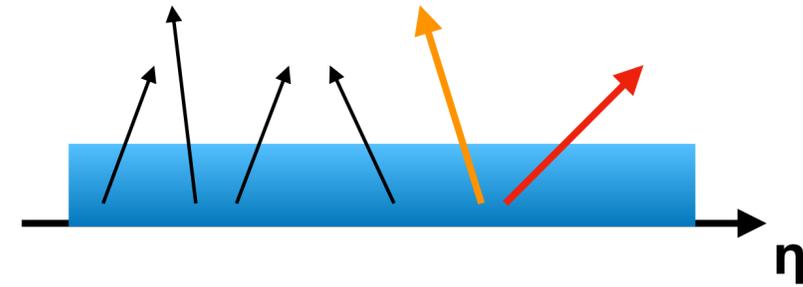


- Instead of the above we measure 2-particle correlations:

$$\langle\langle 2 \rangle\rangle = \langle\langle \cos n(\varphi_1 - \varphi_2) \rangle\rangle = \langle v_n^2 \rangle$$

$$v_n\{2\} = \sqrt{\langle\langle 2 \rangle\rangle}$$

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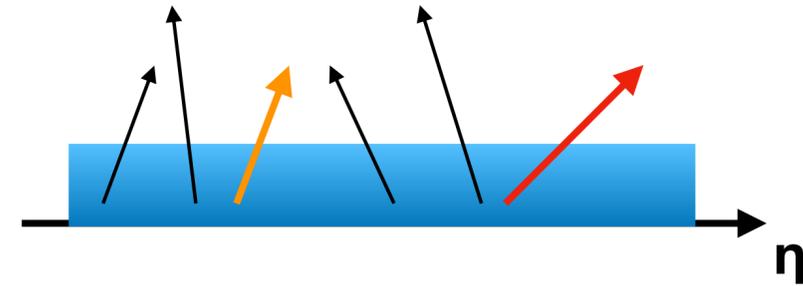


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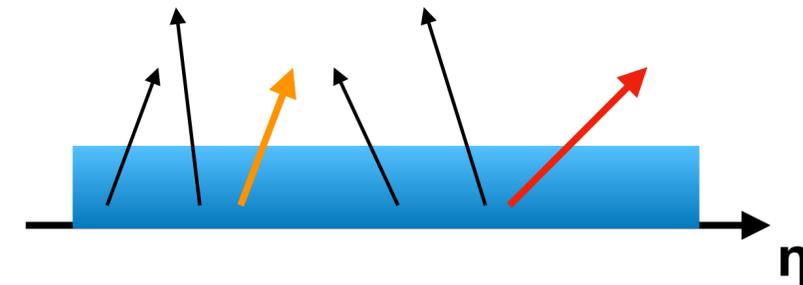


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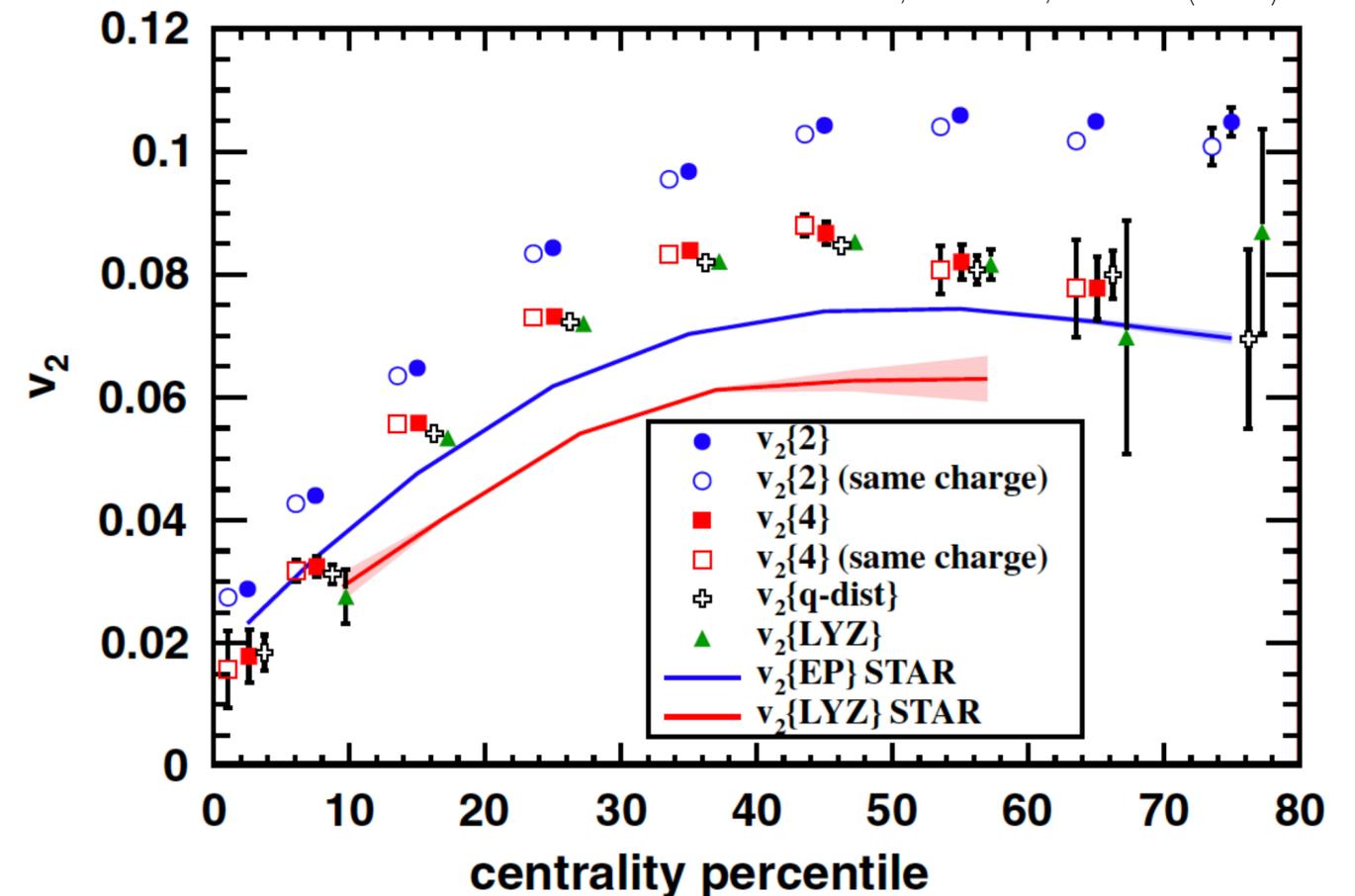
$$v_n\{2\} = \sqrt{\langle\langle 2 \rangle\rangle}$$

- And correlations of higher orders

$$\langle\langle 4 \rangle\rangle = \langle\langle \cos n(\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4) \rangle\rangle = \langle v_n^4 \rangle$$

$$v_n\{4\} = \sqrt[4]{2\langle\langle 2 \rangle\rangle^2 - \langle\langle 4 \rangle\rangle}$$

ALICE, PRL 105, 252302 (2010)



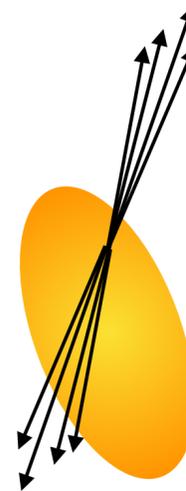
- Collisions of small systems with large multiplicities surprisingly exhibit flow-like patterns, only seen in AA collisions and understood as a manifestation of collective flow of a hot and dense medium
- Investigated via measurements of azimuthal correlations (similarly as in AA collisions)

See talk of Daniel Mihatsch
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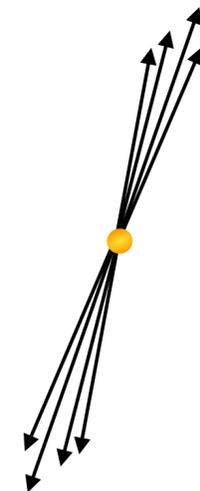
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 - Overall ~ 10 000 particles, dominated by the bulk



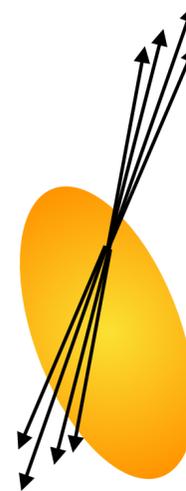
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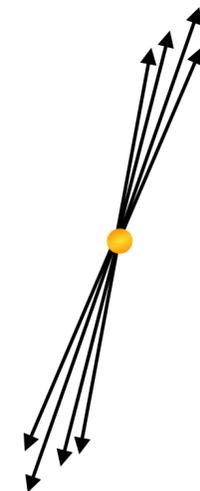
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can we still safely
do correlations ?

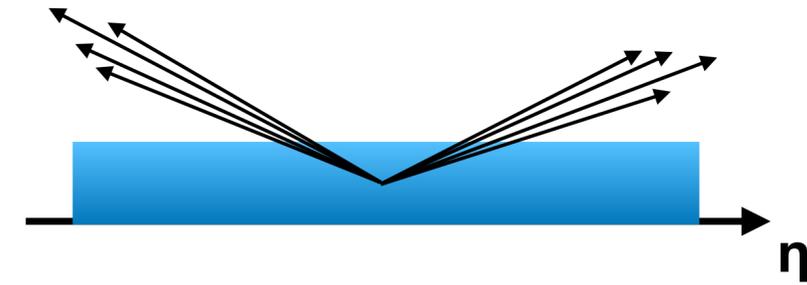
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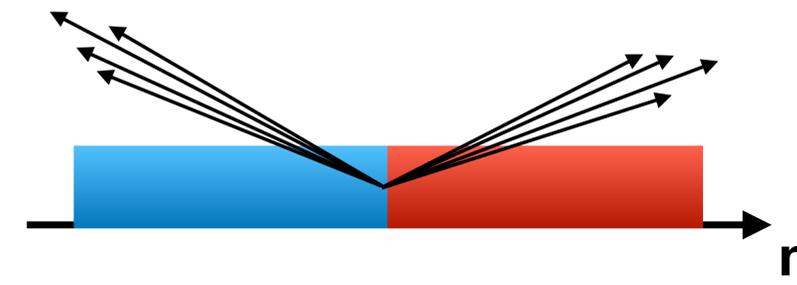
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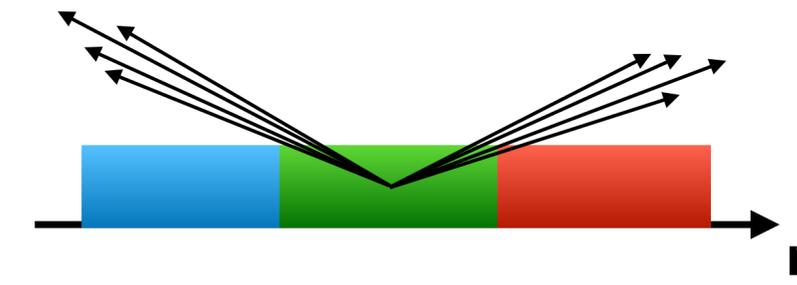
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we include whole jet in correlations

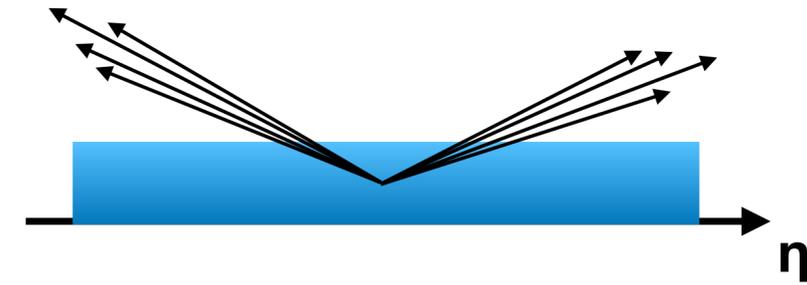


we do not correlate particles from the same jet cone
but we still correlate particles from opposite jet cones



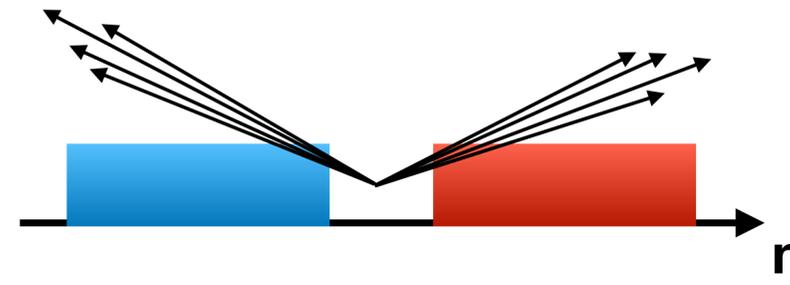
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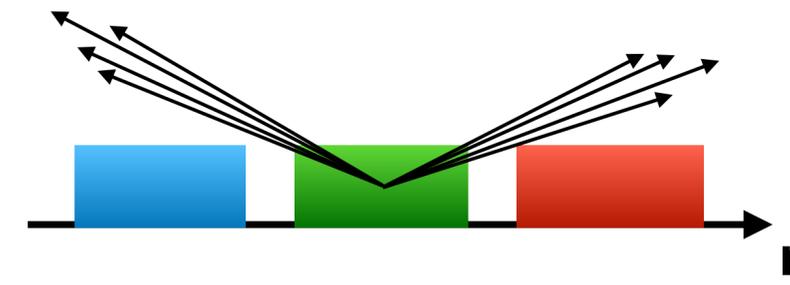
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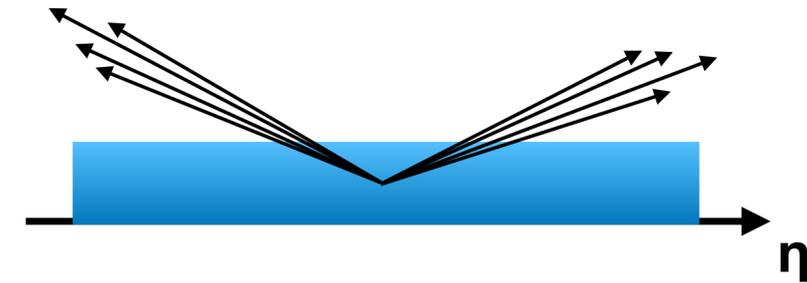
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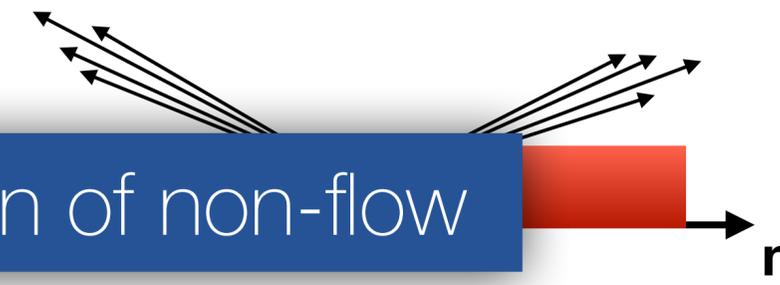
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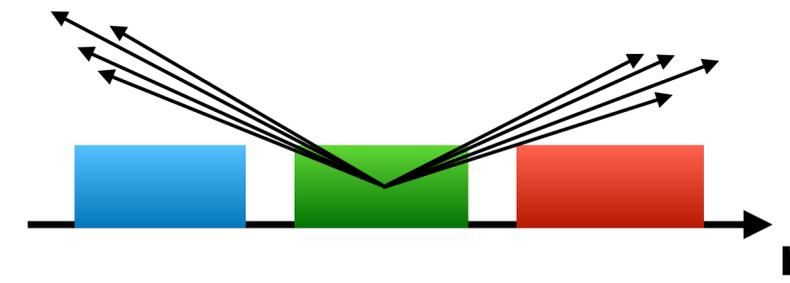
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$$\langle \cos n \text{ the larger the } \eta \text{ gap, the better suppression of non-flow} \rangle$$



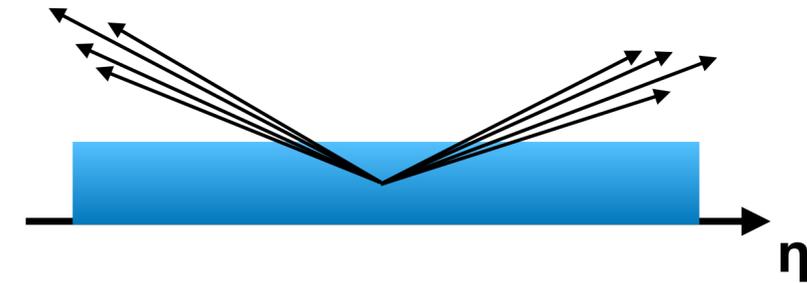
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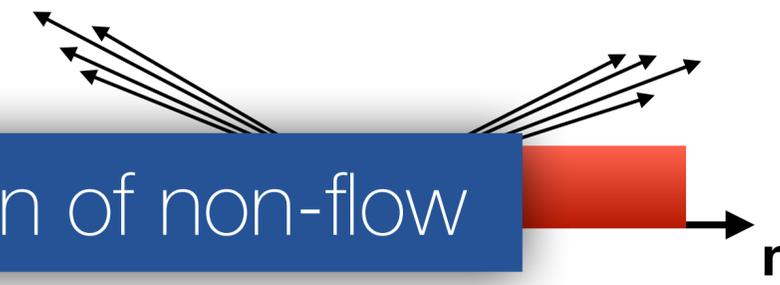


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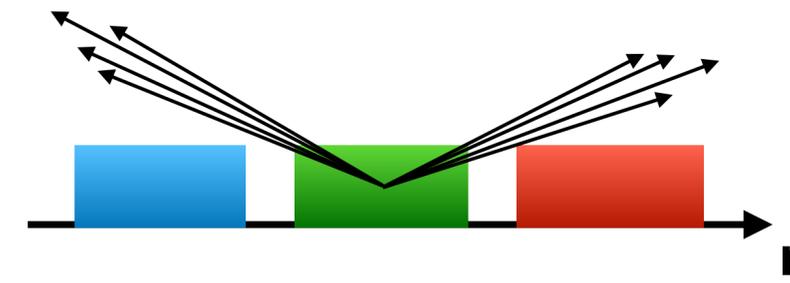
the larger the η gap, the more suppression of non-flow

Not necessarily true



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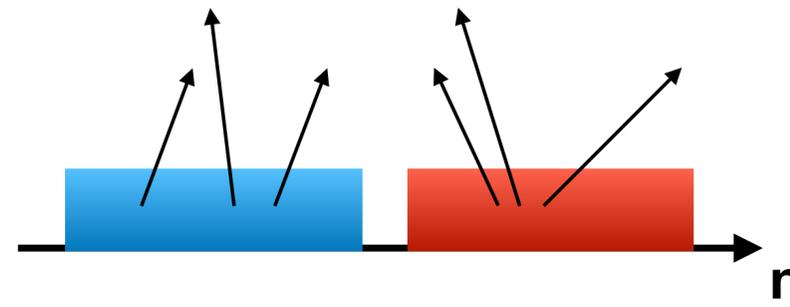
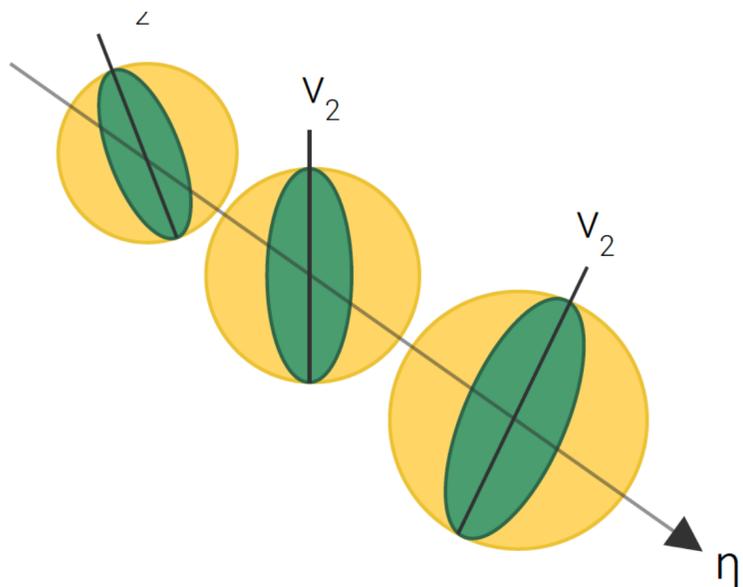
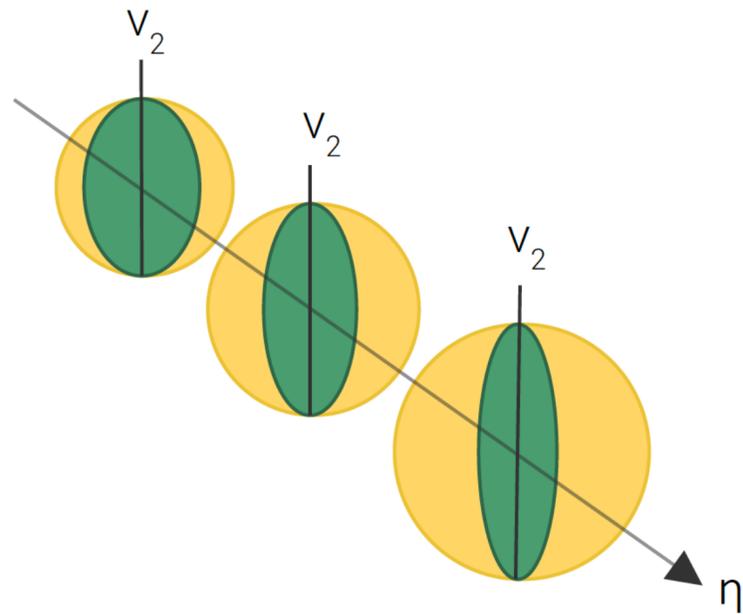
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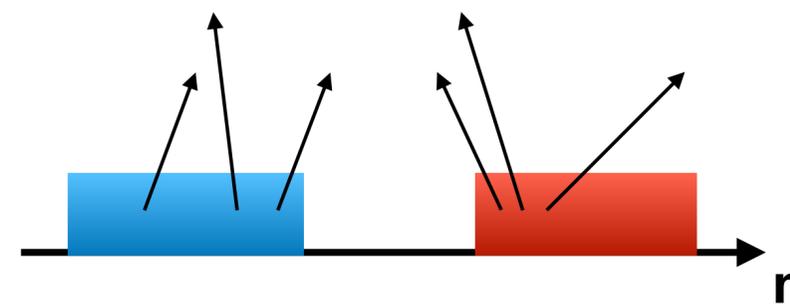
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$$V_n = v_n e^{in\Psi_n}$$

- Flow vector fluctuates in p_T and η
 - Flow magnitude fluctuations
 - Flow angle fluctuations
- Fluctuations in η (=longitudinal fluctuations) pose a danger to the subevent method



small η gap:
large non-flow
small decorrelation



large η gap:
small non-flow
large decorrelation

- Interesting measurements due to their link to the event-by-event fluctuations of the initial energy density (distribution of the participating nucleons in the initial overlap region)
 - Found by hydrodynamic calculations
- Interesting measurements for theorists
 - Allows to constrain 3D modelling of the evolution of a heavy-ion collision

η - dependent flow vector fluctuations

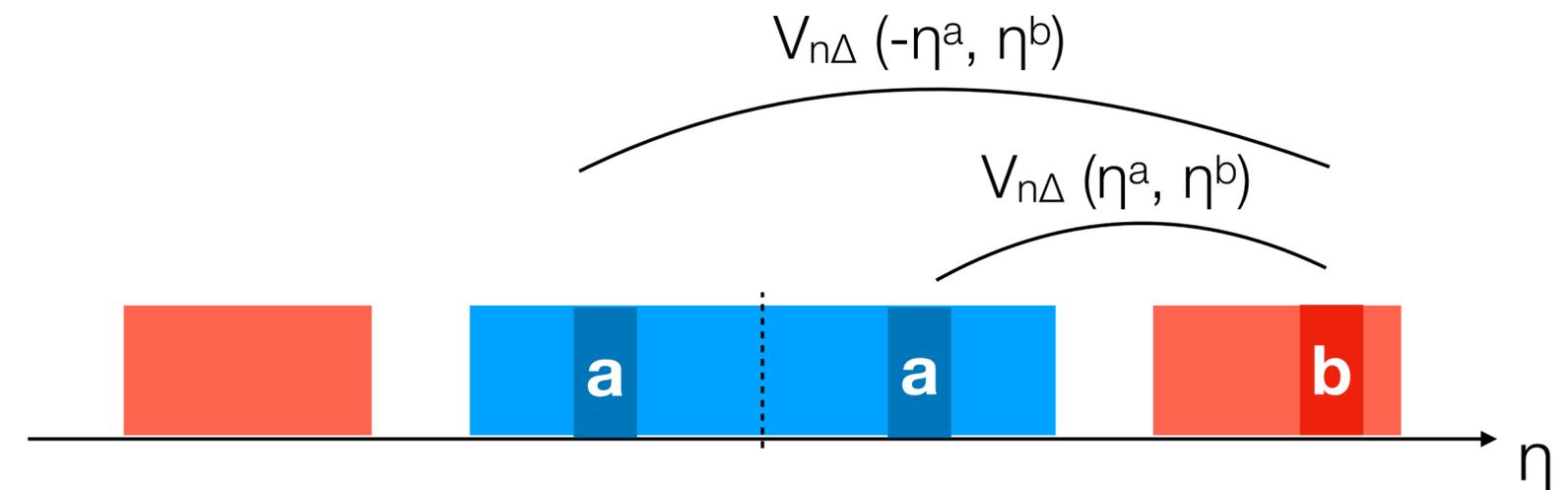
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measures relative flow angle fluctuations

$$r_n(\eta^a, \eta^b) = \frac{\langle v_n(-\eta^a) v_n(\eta^b) \cos\{n[\Psi_n(-\eta^a) - \Psi_n(\eta^b)]\} \rangle}{\langle v_n(\eta^a) v_n(\eta^b) \cos\{n[\Psi_n(\eta^a) - \Psi_n(\eta^b)]\} \rangle}$$

$$r_n = 1 \quad \dots \text{no fluctuations} \quad \Psi_n(\eta^a) = \Psi_n(\eta^b)$$

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Flow vector decorrelations

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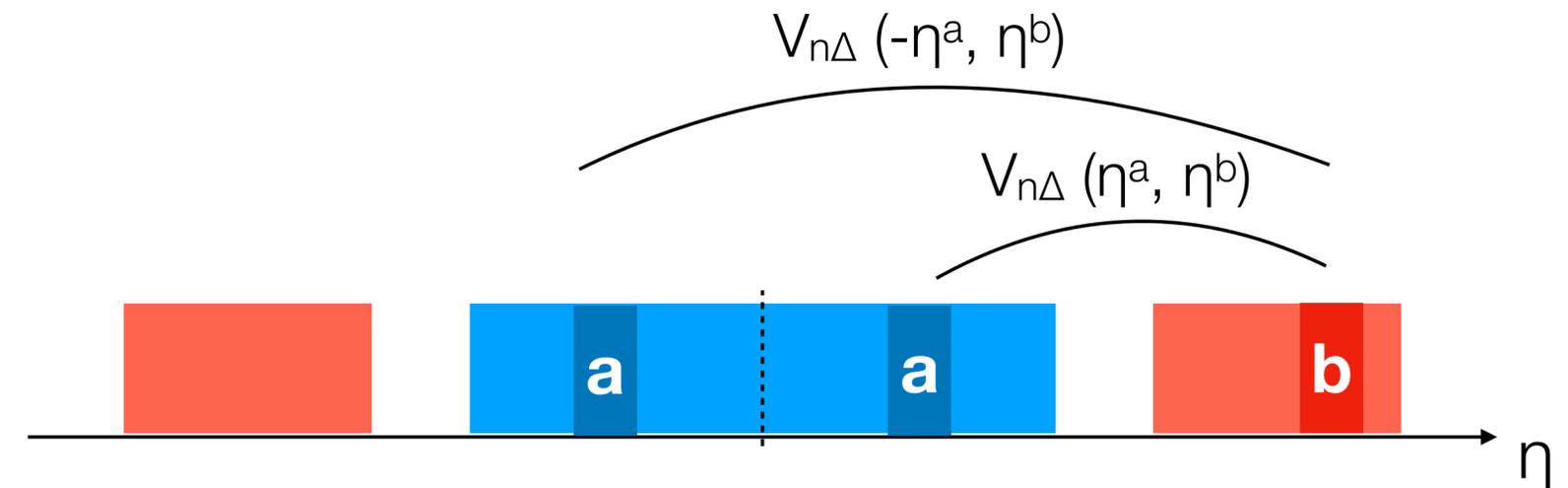
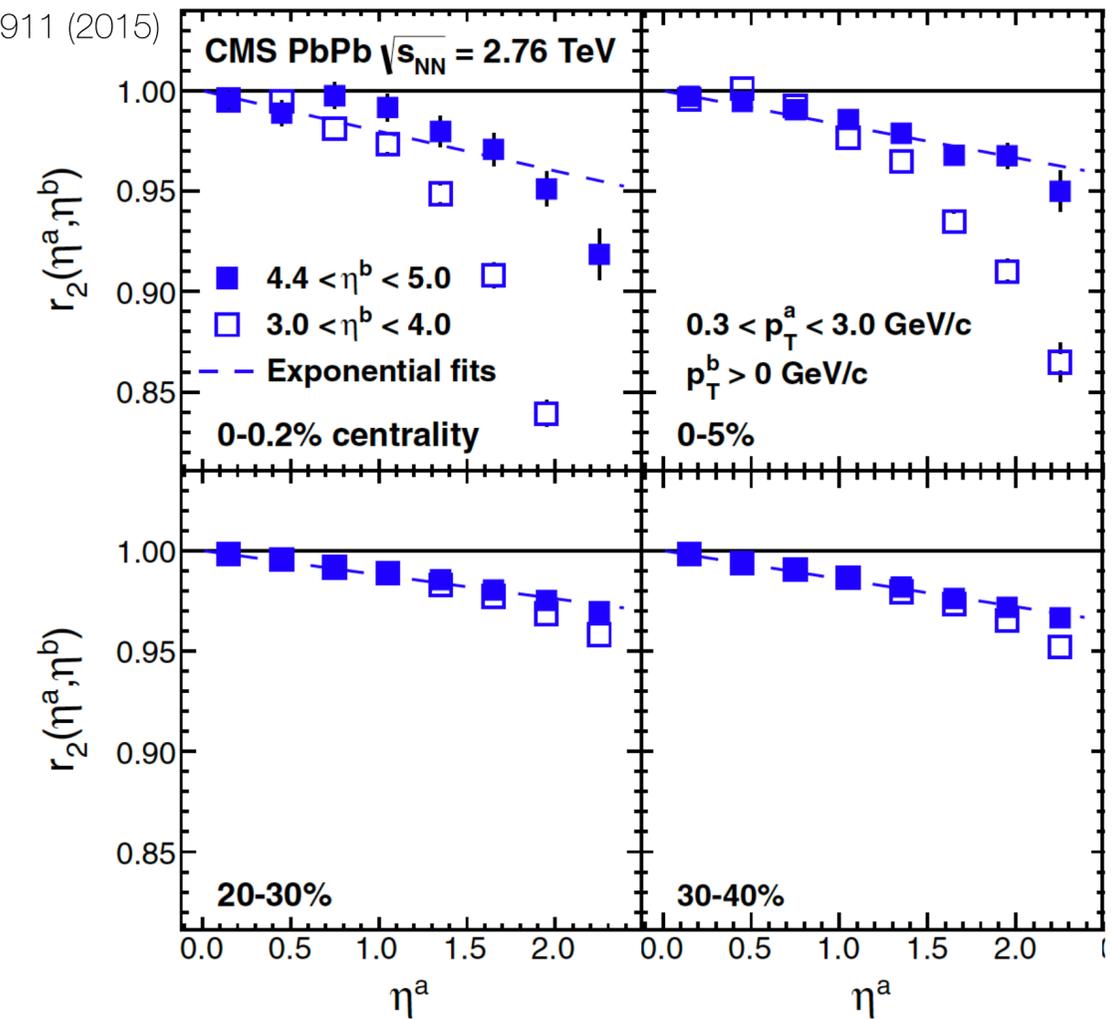
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CMS, PRC 92, 034911 (2015)



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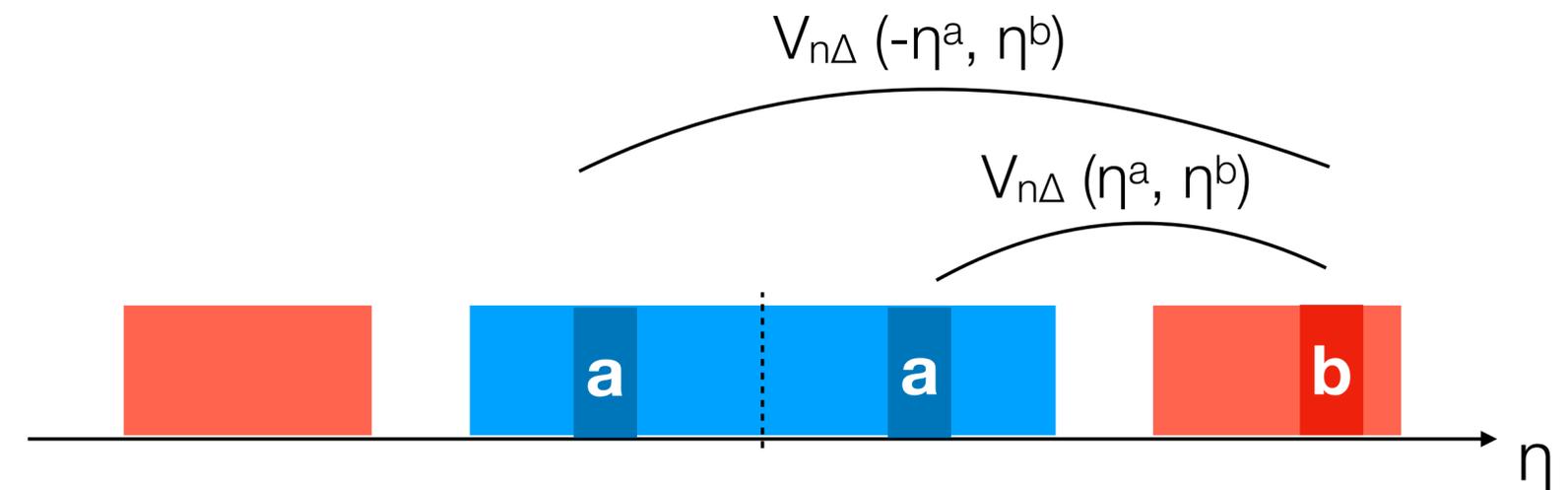
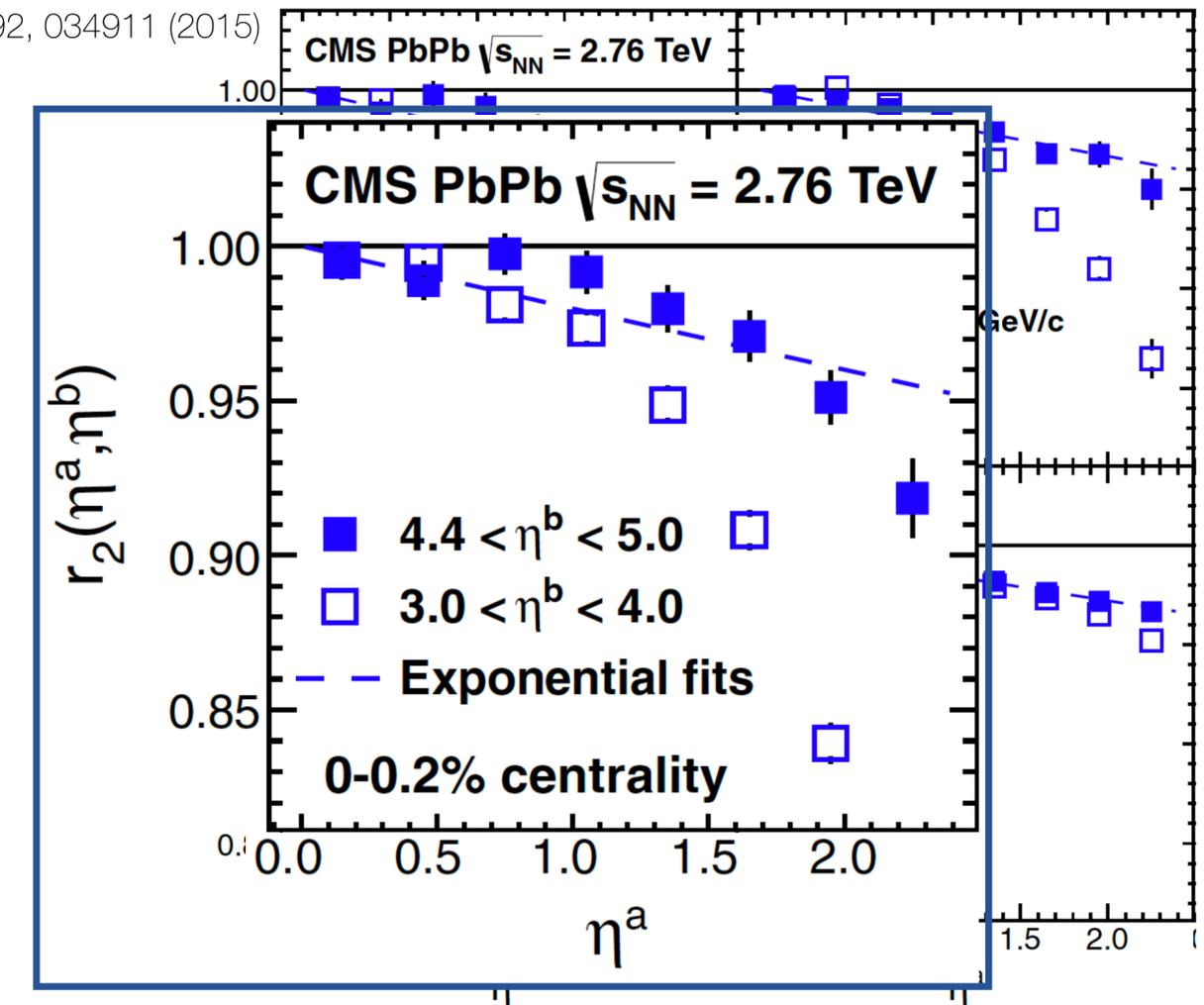
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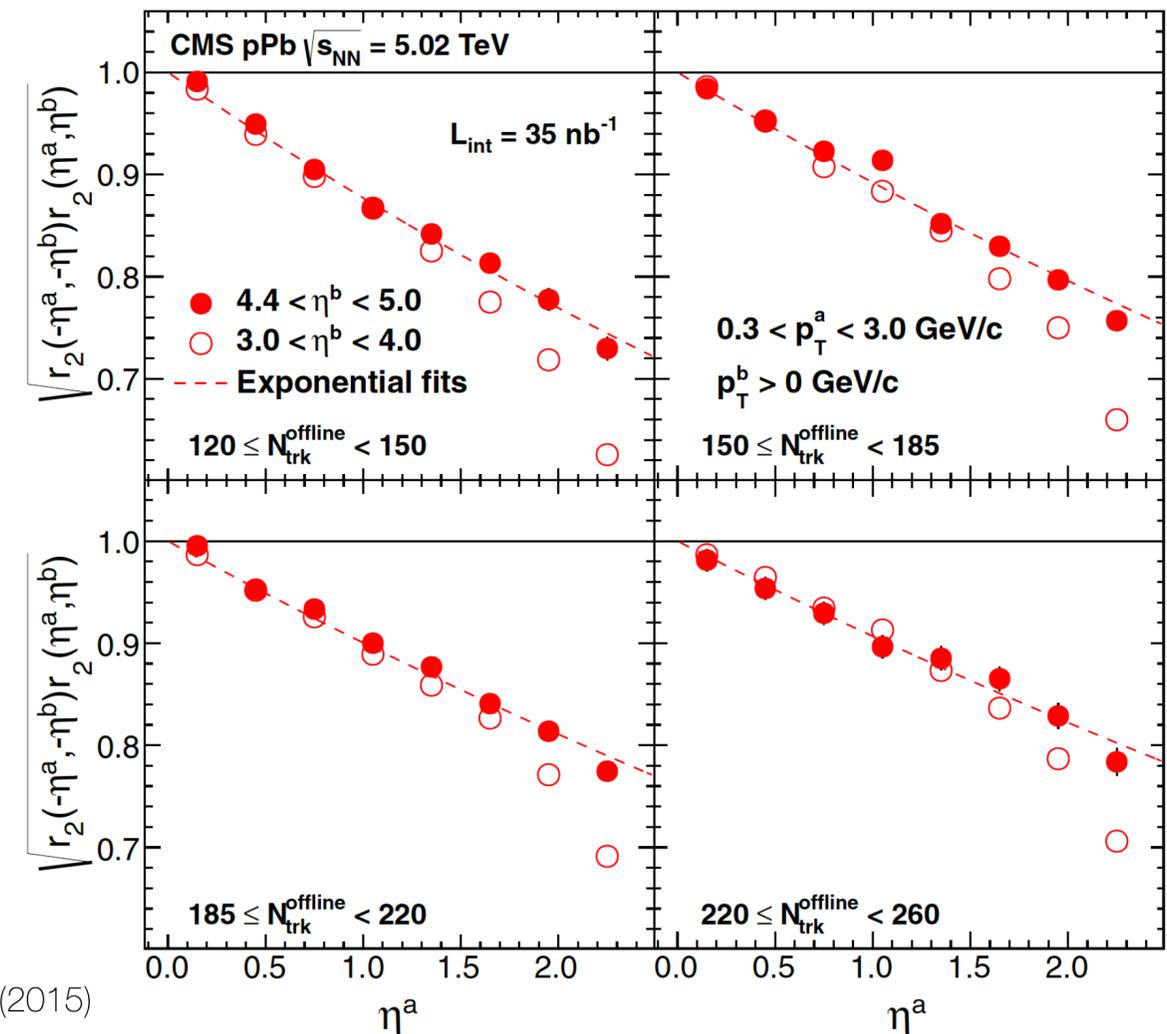
CMS, PRC 92, 034911 (2015)



- It is desirable to measure such observables in small systems
- If they indeed develop a hydrodynamic flow, then they should also exhibit flow vector decorrelations
 - We would expect stronger decorrelation because small systems are dominated by initial state fluctuations rather than the overall geometry

Decorrelations in small systems

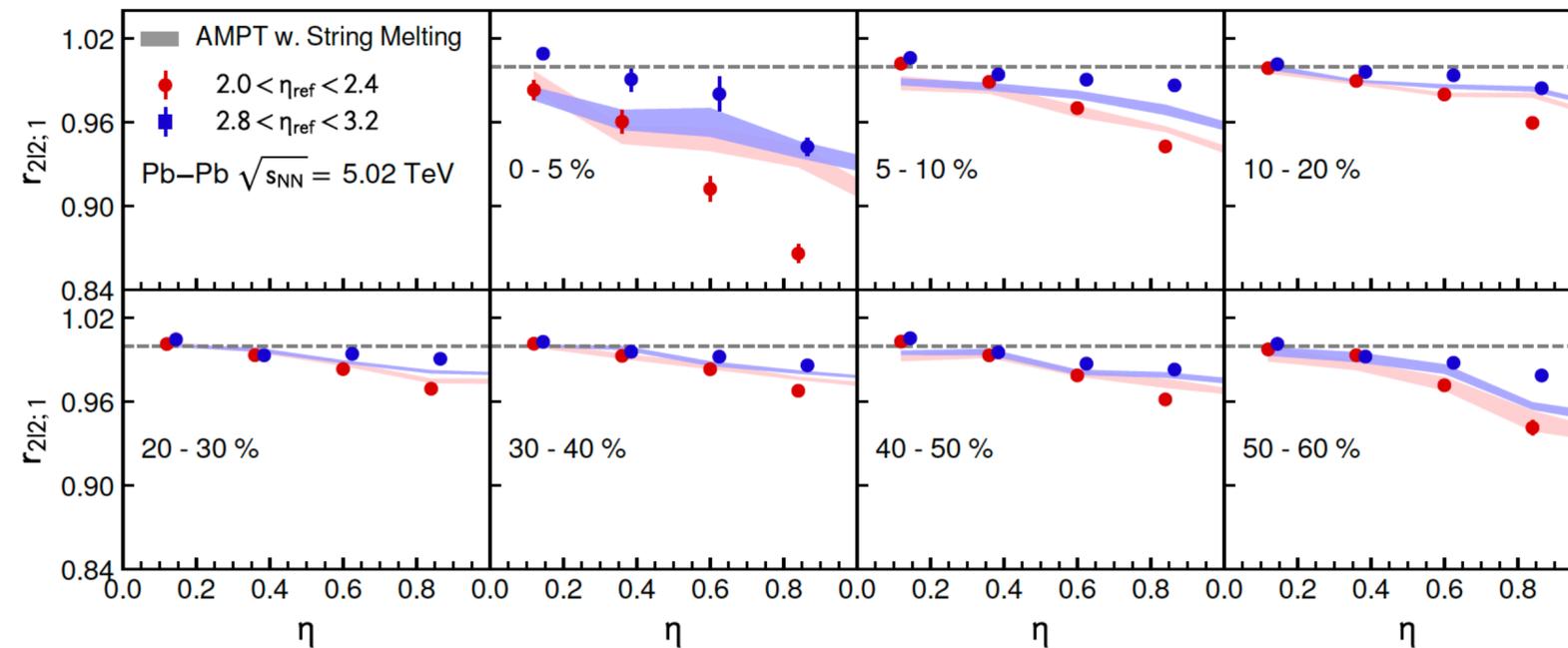
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- If they indeed develop a hydrodynamic flow, then they should also exhibit flow vector decorrelations
 - We would expect stronger decorrelation because small systems are dominated by initial state fluctuations rather than the overall geometry
- So far only measured by CMS in p-Pb collisions



CMS, PRC 92, 034911 (2015)

What we already measured

- ALICE measured this using the “old” detector set-up in Pb-Pb collisions (paper in preparation now)
 - Used the TPC + FMD, which didn't have tracking, measurements were affected by large contamination from secondaries



What we could measure in the future

- Decorrelations in AA collisions using the new MFT detector
- Decorrelations in small systems (including pp ?)

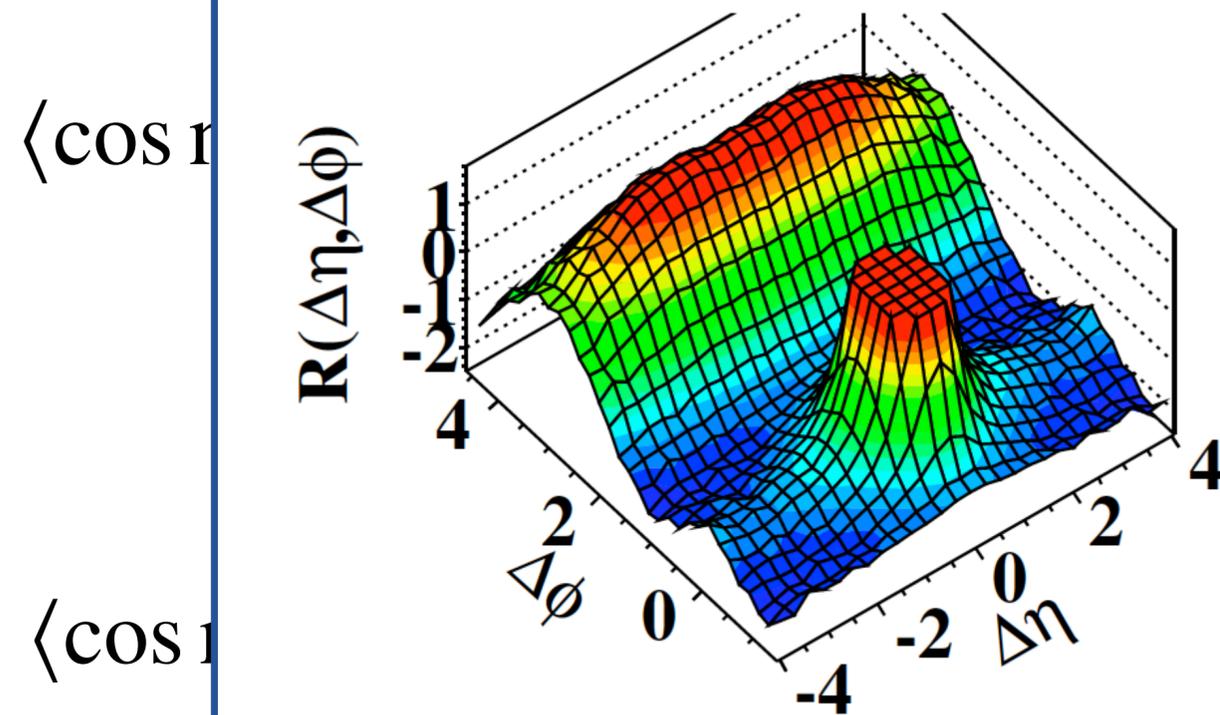
Summary

- Stay tuned for the (hopefully) exciting new results in Run 3 :)

$$\langle \cos n(\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4) \rangle$$

Side note

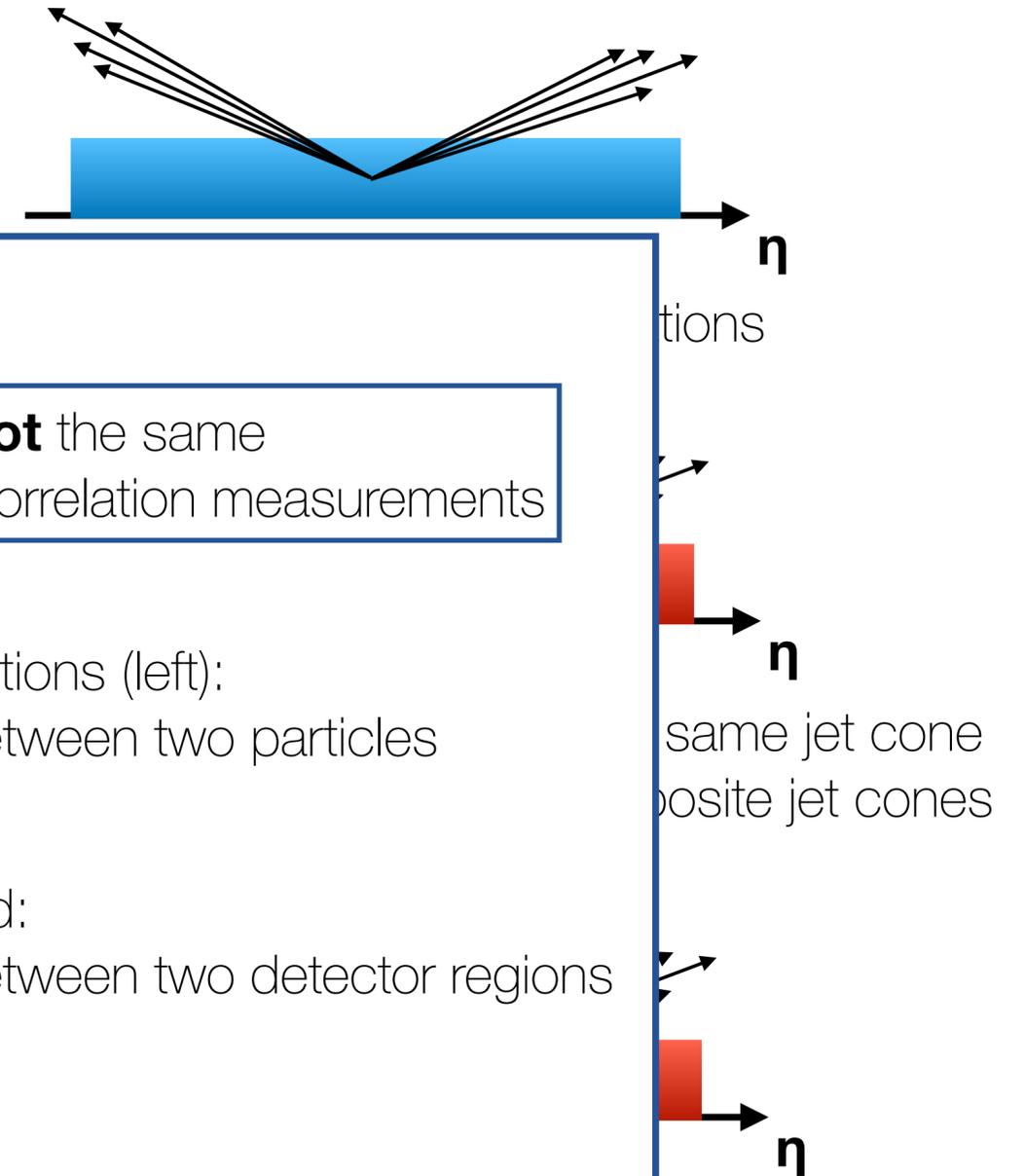
(d) CMS $N \geq 110$, $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$



This eta gap is **not** the same as in di-hadron correlation measurements

Di-hadron correlations (left):
eta separation between two particles

Subevent method:
eta separation between two detector regions



we do not correlate particles from the same jet cone
nor from opposite jet cones

p_T - dependent flow vector fluctuations

$$r_n(p_T^a, p_T^b) \equiv \frac{V_{n\Delta}(p_T^a, p_T^b)}{\sqrt{V_{n\Delta}(p_T^a, p_T^a) V_{n\Delta}(p_T^b, p_T^b)}}$$

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Flow vector decorrelations

p_T - dependent flow vector fluctuations

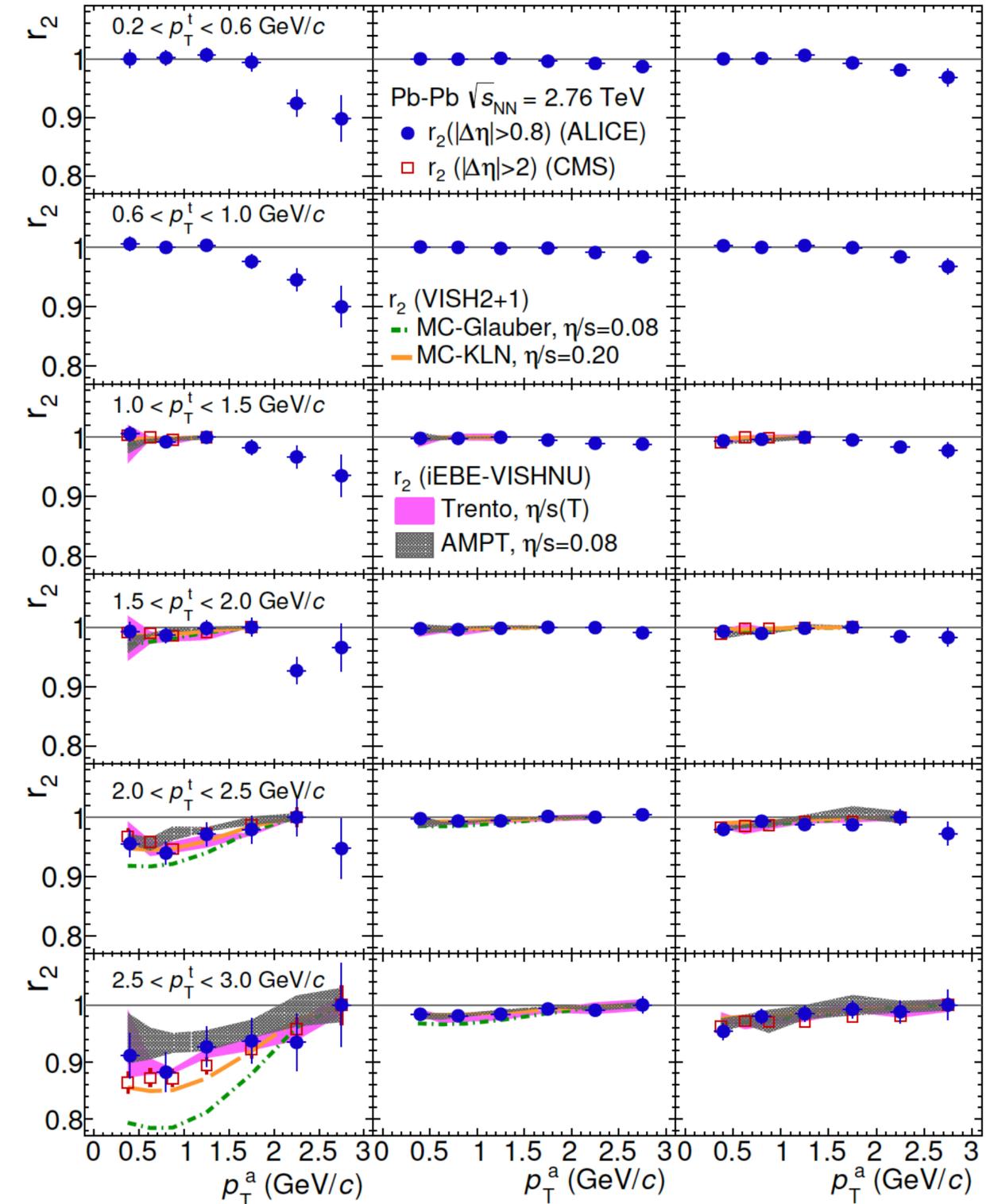
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ALICE, JHEP 09, 032 (2017) 0-5% 20-30% 40-50%



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$$r_n = 1 \quad \dots \text{no fluctuations} \quad \Psi_n(p_T^a) = \Psi_n(p_T^b)$$

$$r_n < 1 \quad \dots \text{fluctuations} \quad \Psi_n(p_T^a) \neq \Psi_n(p_T^b)$$

ALICE, JHEP 09, 032 (2017) 0-5% 20-30% 40-50%

