

CZECH TECHNICAL UNIVERSITY IN PRAGUE
Faculty of Nuclear Sciences and Physical Engineering



**Study of $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^+\pi^-\pi^-$
production in central exclusive processes
with the STAR detector at RHIC**

This work is a shortened version of the author's master's thesis, which
has been reduced to better meet the requirements of 11.
Česko-Slovenská studentská vědecká konference ve fyzice.

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Název práce: **Studie produkce $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ a $\pi^+\pi^+\pi^-\pi^-$ v centrálních exkluzivních procesech s detektorem STAR na RHIC**

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Abstrakt: Tato práce se zabývá prvním měřením centrální exkluzivní produkce v proton-protonových srážkách při energii $\sqrt{s} = 510$ GeV naměřených na experimentu STAR. Při této energii dominuje dvojitá výměna pomeronů. Z toho důvodu jsou tyto srážky velmi slibné pro hledání vázaných stavů gluonů, glueballů. Experimentální potvrzení jejich existence by silně podpořilo platnost kvantové chromodynamiky. Difrakčně rozptýlené protony pohybující se neporušeny uvnitř svazkové trubice byly změřeny pomocí detektorů Roman Pot, což umožnilo plnou kontrolu nad kinematikou interakce a tedy ověření exkluzivity. První výsledky ve formě distribuce invariantní hmotnosti centrálně produkovaných $\pi^+\pi^-$, K^+K^- a $p\bar{p}$ párů jsou prezentovány v této práci.

Klíčová slova: Pomeron, Glueball, RHIC, STAR, Roman Pot

Title: **Study of $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^+\pi^-\pi^-$ production in central exclusive processes with the STAR detector at RHIC**

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Abstract: This thesis is dedicated to the first measurement of the central exclusive production process in proton-proton collisions at RHIC with the STAR detector at $\sqrt{s} = 510$ GeV. At this energy, this process is dominated by a Double Pomeron Exchange mechanism. Therefore, it is very promising for the search of gluon bound states, glueballs. The experimental confirmation of their existence would be yet another strong support for the validity of the quantum chromodynamics theory. The diffractively scattered protons, moving intact inside the RHIC beam pipe after the collision, were measured in the Roman Pots system allowing full control of the interaction's kinematics, and thus verification of the exclusivity. The preliminary results on the invariant mass distributions of centrally exclusively produced $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs measured within the STAR acceptance are presented in this thesis.

Key words: Pomeron, Glueball, RHIC, STAR, Roman Pot

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Introduction

Everything around us is composed of elementary particles such as quarks, leptons and bosons. The standard model of particle physics describes how the elementary particles form the matter around us using three of the four fundamental forces: electromagnetic, weak and strong. The strong interaction is described by the quantum chromodynamics (QCD), a theory using gluons as the exchange particles. Together gluons and quarks are forming bound states e.g. proton or neutron which are a part of each atom. Among them are particles that are predicted by the QCD but not yet fully experimentally confirmed, like the gluon bound states called glueballs or exotic mesons. The experimentally measured glueball candidates like the $f_0(1500)$ and the $f_0(1710)$ are expected to decay into $\pi\pi$, $\pi\pi\pi\pi$ and KK states, respectively. Hence, it is interesting and important to perform a search for all these final states at every new opportunity. Such is for example, the STAR experiment at the Relativistic Heavy Ion Collider (RHIC) at the Brookhaven National Laboratory in the USA.

High energy proton-proton (pp) colliders like RHIC and the Large Hadron Collider (LHC) in the European Organization for Nuclear Research (CERN) complex provide a unique opportunity to study diffractive processes such as central exclusive production (CEP). In CEP, the protons stay intact and central state is produced with quantum numbers of vacuum. It is believed that the Double Pomeron Exchange (DIPE) is a dominant CEP mechanism at high energies. Since the Pomeron is in the lowest order in QCD described as a pair of gluons, the DIPE process provides a gluon-rich environment, which is considered to be a potential source of glueballs. The experimental confirmation of the existence of the glueball would be yet another strong support for the validity of the QCD theory. Despite its theoretical predictions, the existence of a glueball has not been unambiguously confirmed yet.

The first chapter of this thesis is a brief theoretical introduction to the Regge framework and CEP. The second chapter the newest results in the CEP process measured at STAR experiment. The following chapter describes the RHIC accelerator complex and the STAR experiment with its sub-detectors: the Time Projection Chamber, the Time of Flight detector and Roman Pots. The main goal of this thesis was to analyse the data from pp collisions at $\sqrt{s} = 510$ GeV collected by the STAR experiment in 2017. Since the data contain triggers from DIPE in CEP, they are very promising for a glueball search. The current state of the analysis is described in chapters 4 - 7. In particular, distributions of invariant mass of centrally exclusively produced $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^-\pi^+\pi^-$ are presented. Finally, in the eighth and last chapter of this thesis the results are compared to Graniitti Monte Carlo event generator.

Chapter 1

Central exclusive production

Before central exclusive production process is described, it is necessary to define the common kinematic variables and introduce theoretical framework of diffraction processes in proton-proton collisions.

1.1 Kinematics and cross-section

Mandelstam variables are commonly used to describe two-body scattering processes of two particles to two particles which can be seen in Fig. 1.1. There are three Mandelstam variables defined as:

$$\begin{aligned}s &= (p_1 + p_2)^2 = (p_3 + p_4)^2, \\t &= (p_1 - p_3)^2 = (p_2 - p_4)^2, \\u &= (p_1 - p_4)^2 = (p_2 - p_3)^2,\end{aligned}\tag{1.1}$$

where p_1, p_2 are the four-momenta of incoming particles and p_3, p_4 are the four-momenta of outgoing particles. Furthermore, the letters s, t, u stand for the channels corresponding to Feynman diagrams [1]. In the reaction $1 + 2 \rightarrow 3 + 4$ the Mandelstam variable s is square of the total center of mass energy and the variable t is the four momentum transfer squared [2].

In high energy physics, the Cartesian coordinate system is used, where z coordinate is parallel to the beam direction and x, y are in the plane perpendicular to the beam direction. A very often used kinematical variable is rapidity y , which is related to the longitudinal momentum of final-state particle p_z and is defined as:

$$y = \frac{1}{2} \ln \frac{E + p_z c}{E - p_z c},\tag{1.2}$$

where E is the energy of particle and c denotes the speed of light. The useful feature of rapidity is that it transforms additively under a Lorentz boost along z coordinate. The rapidity y is often replaced by the pseudorapidity η since it is hard to measure the rapidity y for highly relativistic particles. The pseudorapidity is defined as:

$$\eta = -\ln \tan \frac{\theta}{2},\tag{1.3}$$

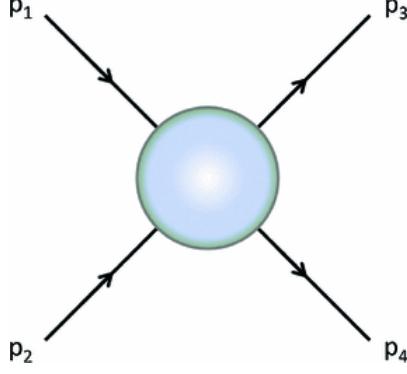


Fig. 1.1: The illustration of the two-to-two particle scattering process. Two incoming particles with momenta p_1 and p_2 are scattered to two outgoing particles with momenta p_3 and p_4 . Taken from Ref. [3].

where the angle θ is a polar angle and specifies the direction of particle motion with respect to the z coordinate. The pseudorapidity is identically equal to rapidity for massless particles and its is approximately equal to rapidity for highly relativistic particles [2].

1.2 Regge theory and Pomeron

The traditional theoretical framework for diffraction is Regge theory, developed in late 1960s. Regge theory describes hadronic reactions at high energies in terms of the exchange of objects called reggeons or Regge trajectories. Each family of bound states or resonances corresponds to a single Regge trajectory.

Regge theory is based on analytical properties of scattering amplitudes \mathcal{A} and thus on properties of the S-matrix. S-matrix, scattering matrix, is a unitary linear operator transforming the initial state $|i\rangle$ of a scattering process into the final state $|f\rangle$:

$$S|i\rangle = |f\rangle. \quad (1.4)$$

The S-matrix elements can be written as:

$$S_{if} \equiv \langle f|S|i\rangle = \delta_{if} + i(2\pi)^4 \delta^4(p_f - p_i) A_{if}(s, t), \quad (1.5)$$

where $A_{if}(s, t)$ is the relativistic scattering amplitude and $p_{i,f}$ are the four-momenta of initial and final states [2].

One of the main Regge ideas was to parametrize A_{if} as a function of the orbital angular momentum l , $\mathcal{A}(l, t)$, where l is a continuous complex variable. The $\mathcal{A}(l, t)$ in the complex l -plane has singularities, poles, whose locations vary with t :

$$l = \alpha(t). \quad (1.6)$$

These poles are known as Regge poles, or reggeons. The functions $\alpha(t)$ are called Regge trajectories and are usually approximated by the Taylor series around $t = 0$ as:

$$\alpha(t) = \alpha(0) + \alpha' \cdot t, \quad (1.7)$$

where $\alpha(0)$ is intercept and α' is a slope of the trajectory.

Values of t such that $\alpha(t)$ is a non-negative integer correspond to the squared mass of a bound state or resonance having that spin. Chew and Frautschi plotted the spins of low lying mesons against their squared masses and noticed that the mesons lie in a straight line as show in Fig. 1.2. A Regge trajectory $\alpha(t) = 0.5 + 0.9 t$ [GeV²], the ρ -trajectory, was obtained by fit. From the intercept and the optical theorem [2], one can obtain the asymptotic behaviour of the total cross-section of simple reggeon exchange [4]:

$$\sigma_{\text{tot}} \propto s^{(\alpha(0)-1)}. \quad (1.8)$$

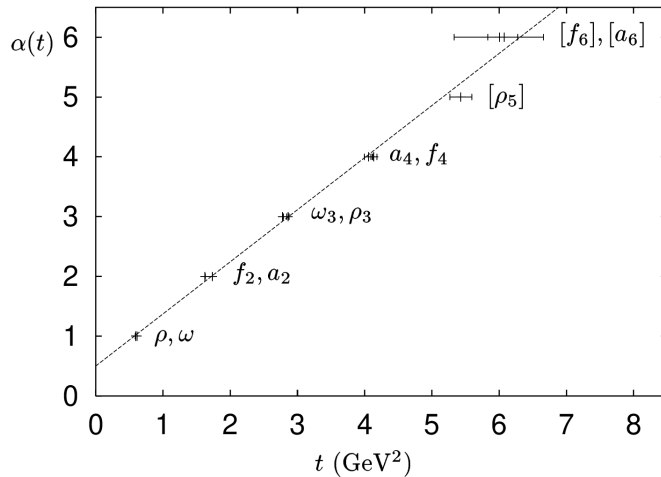


Fig. 1.2: The Chew-Frautschi plot. Particle spins $\alpha(t)$ plotted against their squared masses t . The straight line is $\alpha(t) = 0.5 + 0.9 t$ [GeV²]. Taken from Ref. [5].

For the ρ -trajectory, the cross-section for a process with isospin exchange falls as s increases, since $\alpha_\rho(0) < 1$. In 1956, Pomeranchuk and Okun proved that the cross-section vanishes asymptotically for any scattering process with charge exchange, the Pomeranchuk theorem. However it was observed experimentally that the total cross-sections for pp and $p\bar{p}$ did not fall as s increases for $\sqrt{s} > 10$ GeV, see Fig. 1.3. Thus, the process must be dominated by the exchange of reggeon with vacuum quantum numbers lying on a new trajectory, whose intercept is greater than one. The new trajectory was called the Pomeron [2, 4, 5].

1.2.1 Pomeron

The Pomeron was named after Soviet theoretical physicist I. Ya. Pomeranchuk. Its existence was postulated by Chew, Frautschi and Gribov in 1961. However the exact nature of the Pomeron still remains elusive. It is known that it is a colour singlet object with internal quantum numbers of the vacuum 0^{++} and intercept $\alpha_{IP}(0) \simeq 1$. The Pomeron trajectory does not correspond to any known particle. In perturbative QCD, in the leading order the Pomeron is represented by a pair of gluons since two gluons is the minimal number of gluons needed to reproduce the Pomeron quantum numbers. This representation was proposed by Low and Nussinov [2, 5].

¹The Pomeron can have also even spin, 2^{++} .

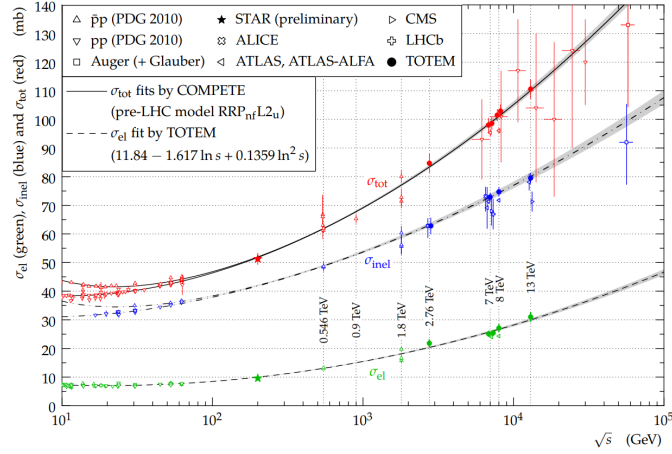


Fig. 1.3: The world data of p, p and p, \bar{p} elastic and total scattering cross-section. Taken from Ref. [6].

It is believed that diffractive processes at high energies are occurring via the exchange of the Pomeron and the exchange of other reggeons with vacuum quantum numbers is suppressed at high energy. Thus in Regge theory the diffractive reactions are those dominated by Pomeron exchange. The Pomeron trajectory:

$$\alpha_{IP}(t) = 1.08 + 0.25 t [\text{GeV}^2] \quad (1.9)$$

with a 2^{++} glueball candidate, measured by the WA91 Collaboration [7], can be seen in Fig. 1.4. The Pomeron trajectory $\alpha_{IP}(t)$ was obtained by fitting the elastic scattering data [2, 5].

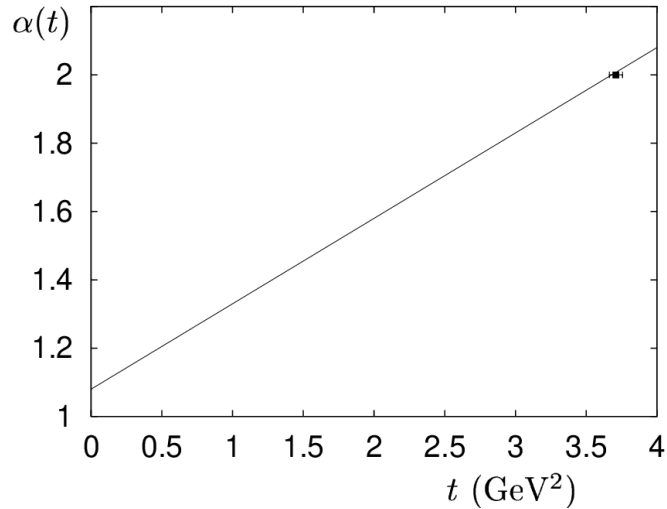


Fig. 1.4: The Pomeron trajectory $\alpha_{IP}(t) = 1.08 + 0.25 t [\text{GeV}^2]$ with a 2^{++} glueball candidate. Taken from Ref. [5].

1.2.2 Glueball

QCD predicts the existence of meson-like objects which contain only gluons as a consequence of their self-interaction [8]. These mesons, gluon bound states, are

called glueballs. Despite its theoretical predictions, the glueball has not yet been observed. As a potential source of glueballs has been regarded the production of mesons in the central region of proton-proton scattering. The experimental confirmation of its existence would be yet another strong support for the validity of the QCD theory [5, 9, 10].

Lattice QCD calculations, well-established non-perturbative approach to solving the QCD theory of quarks and gluons, have predicted the lowest-lying scalar glueball state in the mass range of $1500 - 1700 \text{ MeV}/c^2$ and tensor and pseudo-scalar glueballs in $2000 - 2500 \text{ MeV}/c^2$. Experimentally measured glueball candidates for the scalar glueball states are the $f_0(1500)$ and the $f_0(1710)$ [5, 9, 10].

1.3 Central exclusive production

Central exclusive particle production is a process $A + B \rightarrow A + X + B$, where the colliding particles A and B emerge intact, and for which all particles in the final state are fully measured. A process is called central when the produced state X is produced at central (small) rapidity and is well separated from outgoing particles by rapidity gap $\Delta y > 3$. The diagram of central exclusive production through double Pomeron exchange (DIPE) compare to the diagram of elastic scattering in proton-proton collisions with $\eta - \phi$ representations can be seen in Fig. 1.5 [9].

CEP does not involve valance quarks and the properties of the produced state X should be independent of the colliding hadrons. Furthermore, the produced state X must have positive parity and charge parity, total charge equal to zero and even spin. As Robson first suggested in 1977 [11], the DIPE is expected to be the dominant CEP mechanism at high \sqrt{s} . The first hadron collider, with sufficiently high enough \sqrt{s} , was the CERN Intersecting Storage Rings [12], where DIPE was first observed. Since then, DIPE has been studied at numerous colliders such as the Super Proton–Antiproton Synchrotron at CERN, the Tevatron at Fermilab, the Relativistic Heavy Ion Collider at Brookhaven National Laboratory and last but not least at the Large Hadron Collider at CERN [13].

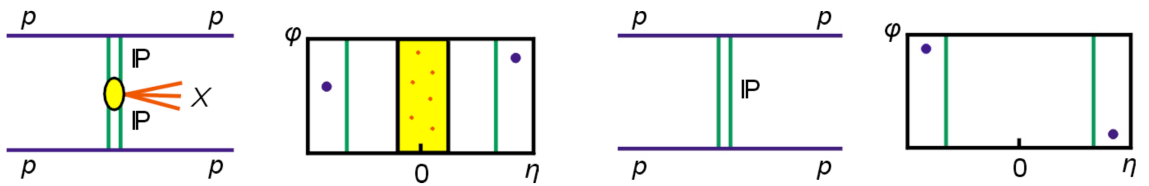


Fig. 1.5: The diagram of central exclusive production through double Pomeron exchange with $\eta - \phi$ representation (left). The diagram of elastic scattering with $\eta - \phi$ representation (right). Taken from Ref. [10].

Chapter 2

Status of central exclusive production

The most recent results in CEP with $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ final states was obtained by the STAR experiment at RHIC [14]. The sample of 15.8 pb^{-1} data collected in pp collisions at $\sqrt{s} = 200 \text{ GeV}$ was used. The DIPE is expected to be dominant at that \sqrt{s} , since the contribution from vector meson photoproduction and photon-photon processes are not significant and additionally suppressed at $-t > 0.04 \text{ GeV}^2/c^2$. Unlike the CMS Collaboration, the STAR experiment uses Roman Pot detectors, allowing efficient measurement of diffractively scattered protons. This enables full control of collision kinematics and verification of the exclusivity by measuring the momenta of all final state particles.

Differential cross-sections for CEP of charged particle pairs $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ as functions of the invariant mass of the pair with small squared four momentum transfers $0.04 < -t_1, -t_2 < 0.2 \text{ GeV}^2/c^2$ of the forward protons were measured and they can be seen in Fig. 2.1 and in Fig. 2.2, respectively. The results were compared to MC model predictions, namely: DiMe, GenEx [15] and PYTHIA8 [16] MBR [17]. The MC models predict only a continuum production under the resonances, therefore only a roughly comparison of shape of non-resonant production and overall normalization can be made. The GenEx is event generator based on a phenomenological models of continuum production mechanism of $\pi^+\pi^-$ or K^+K^- pairs as is the DiMe model. The differences between the DiMe and GenEx model predictions are attributed to the absorption effects, since the GenEx does not take the absorption corrections explicitly as DiMe. The MBR model implemented in PYTHIA8 generates exclusive charged particle pairs through fragmentation and hadronization of the central state.

The results are restricted to the fiducial region for the forward protons defined as:

$$\begin{aligned} (p_x + 0.3 \text{ GeV}/c)^2 + p_y^2 &< 0.25 \text{ GeV}^2/c^2, \\ 0.2 \text{ GeV}/c &< |p_y| < 0.4 \text{ GeV}/c, \\ p_x &> -0.2 \text{ GeV}/c, \end{aligned} \tag{2.1}$$

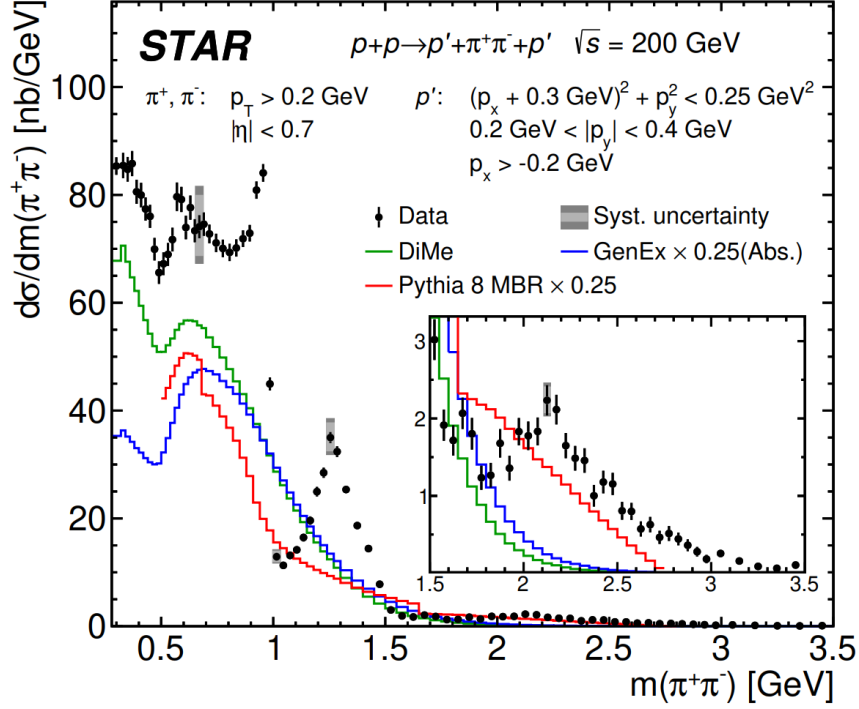


Fig. 2.1: Differential cross-section as a function of invariant mass of $\pi^+\pi^-$ compared to predictions from MC models. Taken from Ref. [14].

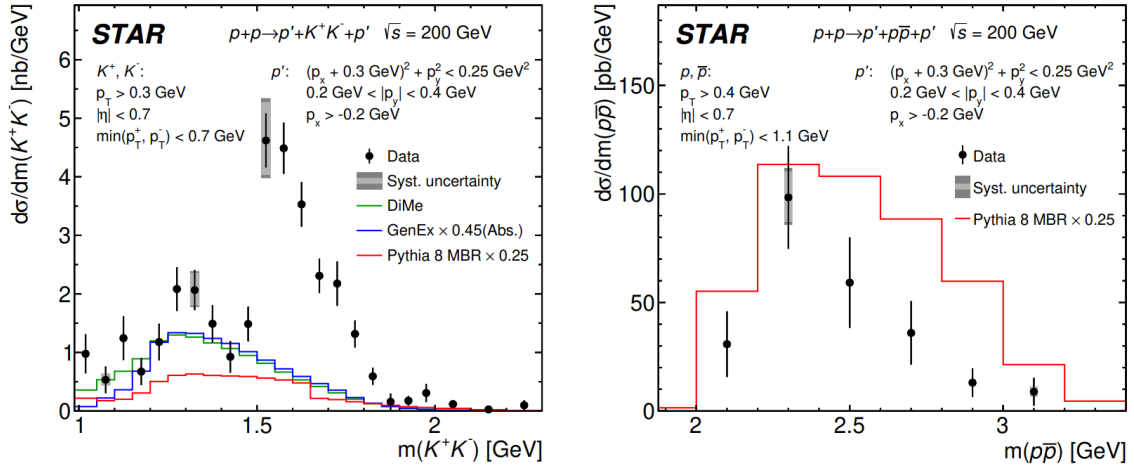


Fig. 2.2: Differential cross-sections as a function of invariant mass of K^+K^- (left) and $p\bar{p}$ (right) compared to predictions from MC models. Taken from Ref. [14].

where p_x , p_y are momenta of the forward protons measured in the Roman Pots. The selection of this fiducial region ensures high geometrical acceptance and track reconstruction efficiency.

In general, the differential cross-section as a function of invariant mass of $\pi^+\pi^-$ shows the main features as the spectrum measured by the AFS Collaboration, the drop at $1 \text{ GeV}/c^2$ and the peak about $1.3 \text{ GeV}/c^2$. These features are attributed to $f_0(980)$ and $f_2(1270)$ resonances discussed above. Furthermore, another resonance is observed at higher invariant masses about $2.2 \text{ GeV}/c^2$. The DiMe model roughly describes the shape of the non-resonant production and it gives better prediction for normalization compare to GenEx or MBR model.

The cross-section for the K^+K^- shows a peak at $1.3 \text{ GeV}/c^2$ and very pronounced peak at $1.55 \text{ GeV}/c^2$. The peaks are attributed to $f_2(1270)$ and $f_2'(1525)$ resonances that are expected in the DIPE. Since the $f_2(1270)$ resonance was measured in the $\pi^+\pi^-$ and K^+K^- production, a ratio of cross-sections in this mass region was measured. The ratio is equal to 18 and it is in good agreement to the branching fractions for decays into $\pi^+\pi^-$ and K^+K^- pairs from the Particle Data Group [18]. The cross-section for the $p\bar{p}$ does not show any resonances. Its shape is reasonably well described by the MBR model, the only available one, through the MBR model overestimate the data by a factor of eight.

The differential cross-sections for CEP of charged particle pairs $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ as a function of the pair rapidity and the difference in the azimuthal angles of the forward protons $\Delta\varphi$ can be seen in Fig. 2.3. The model predictions describe the shapes of measured distributions quite well. The differential cross-sections show a strong suppression about 90° due to fiducial cuts applied to the forward protons. The fiducial cuts are also mainly responsible for a hole at $0.5 \text{ GeV}/c^2$ in the cross-section for the $\pi^+\pi^-$ shown in Fig. 2.1, since the cross-section were measured in two $\Delta\varphi$ regions: $\Delta\varphi < 90^\circ$ and $\Delta\varphi > 90^\circ$. The scattered protons with $\Delta\varphi < 90^\circ$ create a recoil to the central system allowing measurement at low invariant mass region. Figure 2.4 shows the differential cross-sections for CEP of charged particle pairs of $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ as functions of the invariant mass of pairs in two regions of the $\Delta\varphi$. The enhancement of cross-sections at low invariant mass in the $\Delta\varphi < 90^\circ$ compared to $\Delta\varphi > 90^\circ$ region is clearly visible. Another differences can be seen: the $f_2(1270)$ resonance in the $\pi^+\pi^-$ cross-section in $\Delta\varphi < 90^\circ$ is significantly suppressed and the $f_0(980)$ resonance is enhanced compared to the $\Delta\varphi > 90^\circ$ region. These observations suggest vertex factorization breaking that was first reported by the WA91 experiment.

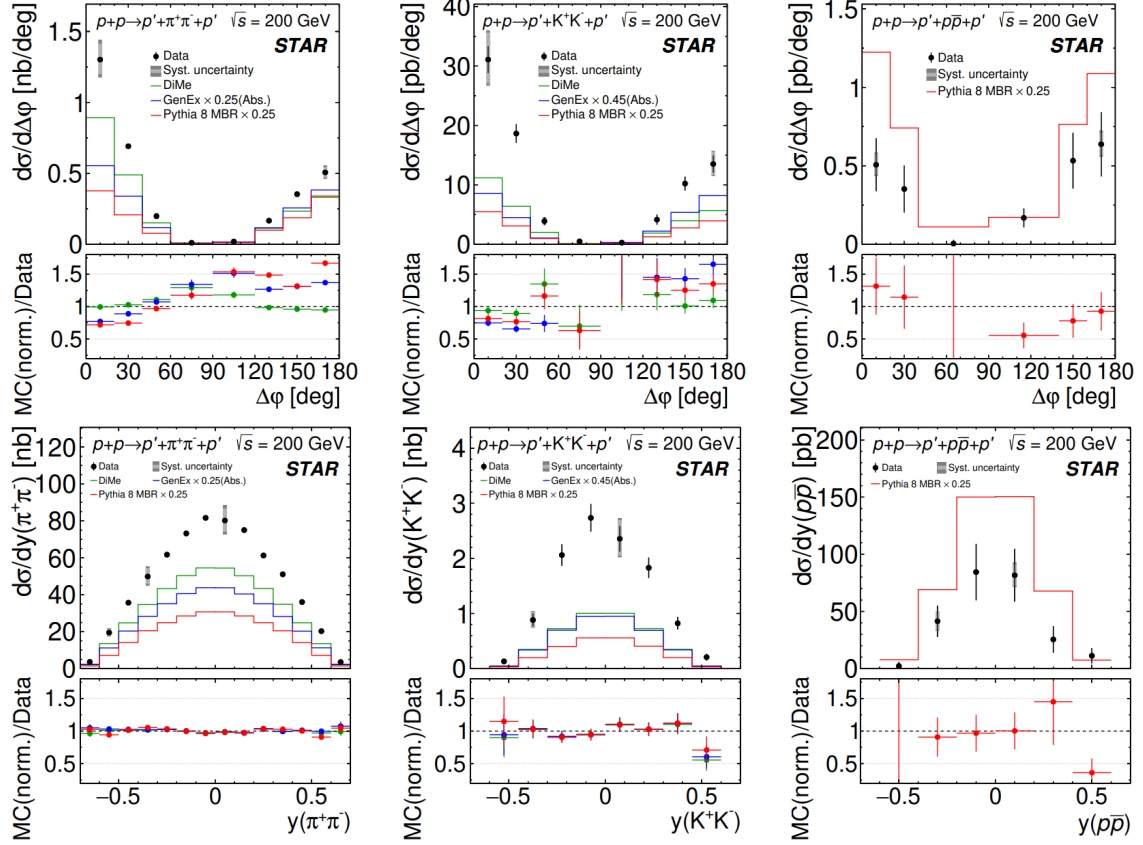


Fig. 2.3: Differential cross-sections for CEP of charged particle pairs $\pi^+\pi^-$ (left), K^+K^- (middle) and $p\bar{p}$ (right) as a function of the difference of azimuthal angles of the forward scattered protons $\Delta\varphi$ (upper row) and the pair rapidity (bottom row) compared to predictions from MC models. Taken from Ref. [14].

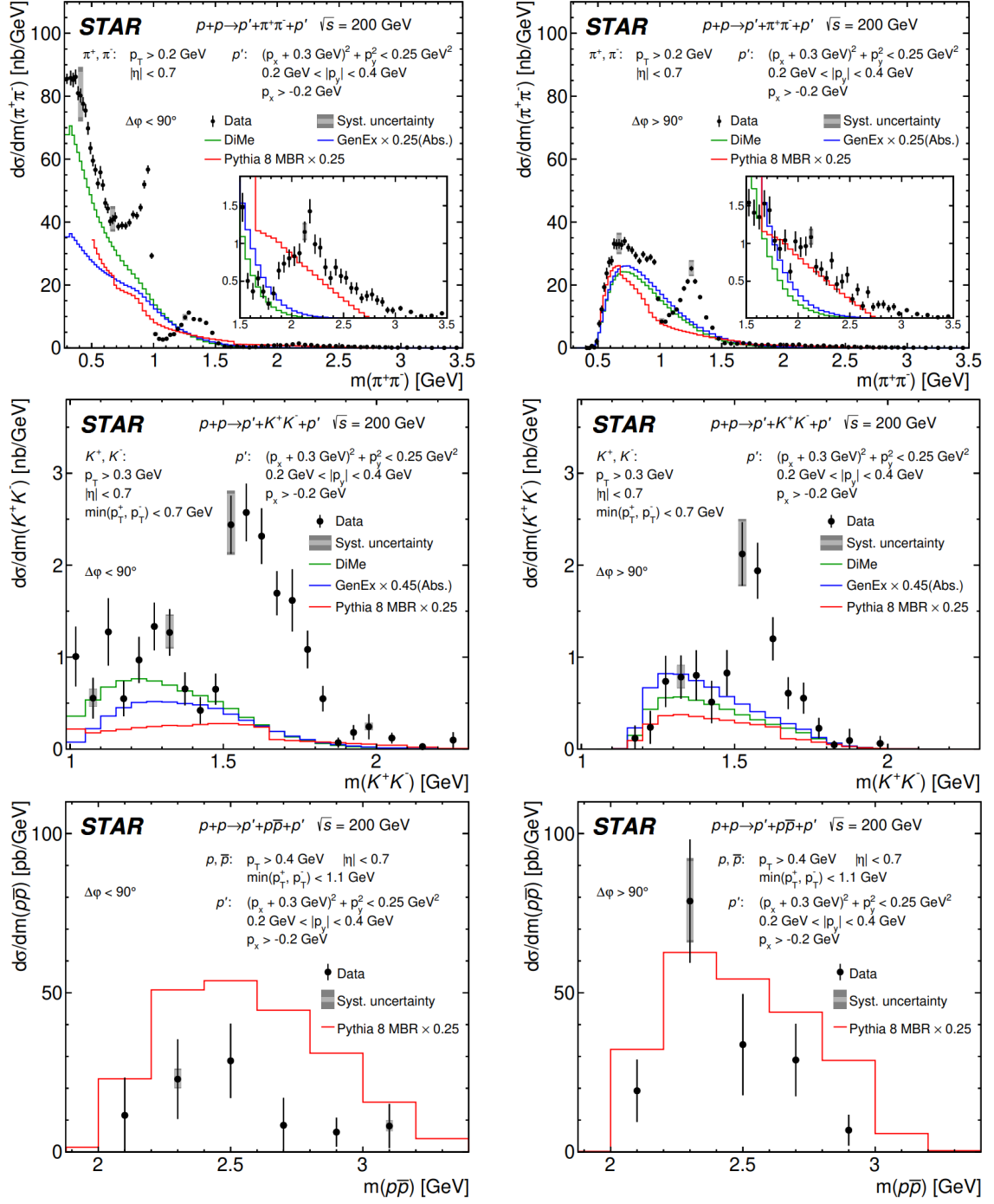


Fig. 2.4: Differential cross-sections for CEP of charged particle pairs of $\pi^+\pi^-$ (upper row), K^+K^- (middle row) and $p\bar{p}$ (bottom row) as a function of the invariant mass of the pair in two regions of the difference of azimuthal angles of the forward scattered protons: $\Delta\phi < 90^\circ$ (left) and $\Delta\phi > 90^\circ$ (right) compared to predictions from MC models. Taken from Ref. [14].

Chapter 3

STAR experiment at RHIC

3.1 Relativistic Heavy Ion Collider

In 2000, the Relativistic Heavy Ion Collider started operations at Brookhaven National Laboratory (BNL). RHIC is one of the biggest scientific facilities in the U.S. It was the first machine in the world capable of colliding heavy ions at very high energies, thus allowing the studies of the Quark-Gluon Plasma. Furthermore, it is the only major accelerator capable of colliding polarized proton beams, which allows measurements aimed at studying the spin structure of the proton [19–21].

Physics program at RHIC is very rich since RHIC is capable of colliding protons or different pairs of ions as U^{92+} , Au^{79+} , Ru^{44+} , Zr^{40+} and Cu^{29+} . Furthermore, it is able to produce asymmetrical collisions, for example protons or deuterons with the gold ions. A summary of RHIC operating modes from 2000 to 2019 is shown in Fig. 3.1. In addition, RHIC also has operated in a fixed target mode. The first STAR fixed-target data from Au-Au collisions at $\sqrt{s_{NN}} = 3.0$ and 7.2 GeV were taken in 2018 as a part of the Beam Energy Scan I program. In 2020, RHIC continued with the Beam Energy Scan II program with Au-Au collisions. Also, in 2020 BNL was selected for eRHIC upgrade, the high energy electron-ion collider (EIC). The EIC will be a particle accelerator that collides polarized electrons with protons or nuclei and it will allow to study the strong force in a great detail. The EIC will be the first accelerator of its kind in the world. The first data taking is planned in 2030 [22].

RHIC consists of two rings with circumference of 3834 m where particles are accelerated in opposite directions, clockwise (blue) or anticlockwise (yellow). The rings intersect in 6 points called interaction points (IP), where experiments can be located. In 2000, there were four experiments: STAR, PHENIX, PHOBOS and BRAHMS. In 2002, the PP2PP experiment was added. The PP2PP was designed to study pp elastic scattering from $\sqrt{s} = 60$ GeV to $\sqrt{s} = 510$ GeV using Roman Pot systems to detect scattered protons. In 2009 the Roman Pots of the PP2PP experiment were incorporated into the STAR experiment, which is the only active experiment at present. Nevertheless, a new experiment sPHENIX [24] is being prepared with planned launch in 2022 [19–21].

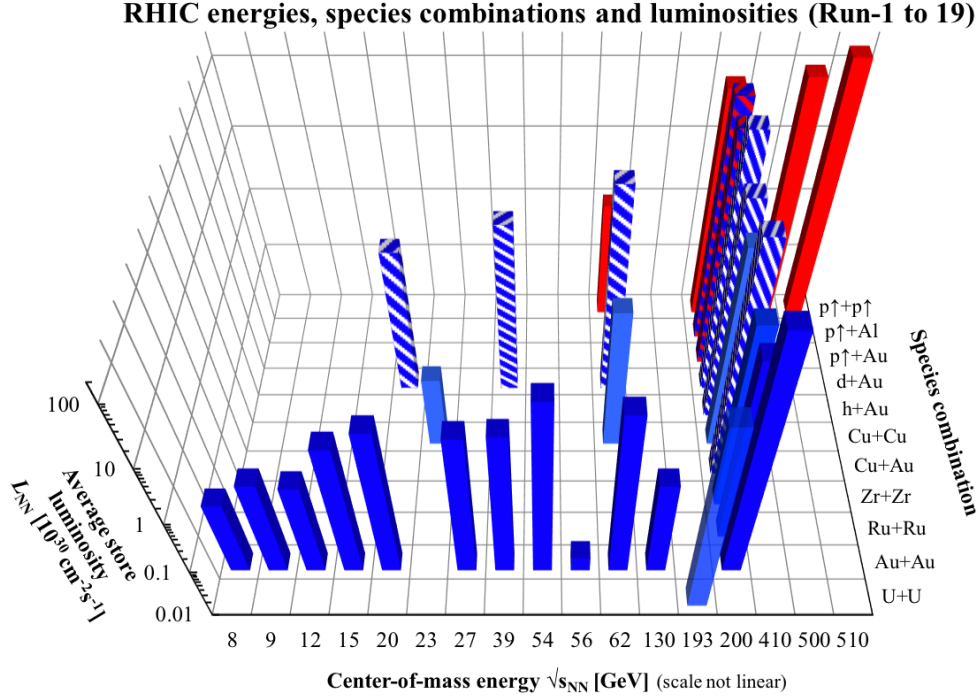


Fig. 3.1: The summary of RHIC energies, species combinations and luminosities from 2000 to 2019. Taken from Ref. [23].

3.2 Solenoidal Tracker at RHIC

The Solenoidal Tracker at RHIC (STAR) is a large-acceptance multi-purpose particle detector designed to study the strongly interacting matter at high temperature and high energy density. It is massive detector weighing 1200 tons consisting of many sub-detectors, such as the Time Projection Chamber (TPC), the Time of Flight detector (TOF), the Beam-Beam Counter (BBC), the Vertex Position Detector (VPD), and the Barrel Electromagnetic Calorimeter (BEMC). Another important part of the experiment is a large solenoidal magnet covering full azimuthal angle and $|\eta| < 1$ in longitudinal direction and generating 0.5 T solenoidal magnetic field parallel to the beam axis. The STAR experiment with its main sub-detectors is shown in Fig. 3.2 [25].

The STAR experiment is very suitable for measuring CEP processes, because of its unique capabilities such as:

1. high-resolution tracking of centrally produced charged particles in the TPC,
2. precise particle identification through the measurement of dE/dx and time of flight,
3. forward rapidity $2.1 < |\eta| < 5.0$ covered by the BBC detectors to ensure rapidity gaps,
4. Silicon Strip Detectors in Roman Pots for measurement of forward protons.

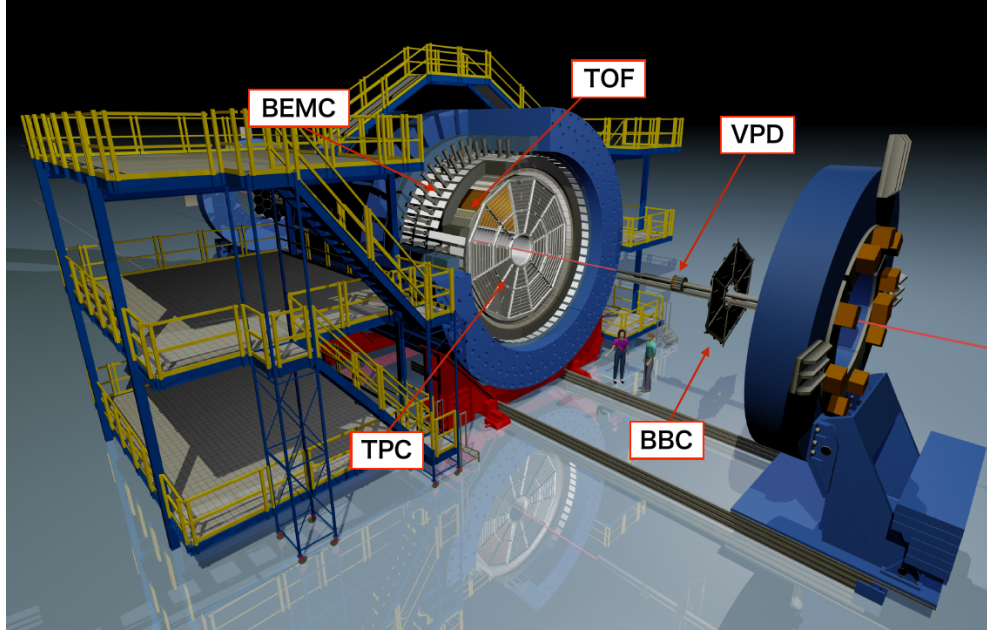


Fig. 3.2: The schematic view of the STAR experiment. Main sub-detectors, including the TPC, the TOF, the BBC, the VPD and the BEMC are highlighted. Taken from Ref. [26].

3.2.1 Time Projection Chamber

The TPC is the heart of the STAR experiment. Its main purpose is precise particle identification and particle tracking. The schematic view of the TPC can be seen in Fig. 3.3. It is a large gas-filled cylindrical detector with inner diameter 1.0 m, outer diameter 4.0 m and is 4.2 m long. It covers full azimuthal angle and pseudorapidity $|\eta| < 1.0$. The fill gas is a P10 gas consisting of 90% argon and 10% methane. The P10 gas is kept at 200 Pa above the atmospheric pressure. The P10 is very convenient for the TPC since a drift velocity of electrons is insensitive to small fluctuation in temperature and pressure [25, 27–29].

3.2.2 Time of Flight detector

The main purpose of the TOF is to extend the PID capabilities of the TPC for particles with higher p_T . Furthermore, the TOF system is used in the trigger of the STAR experiment to select charged particle multiplicity in the central rapidity for various physics processes. For the CEP events it helps triggering on low multiplicity events in the TPC [30].

3.3 Roman Pots

The Roman Pot (RP) technology is commonly used to detect and measure forward protons from elastic or inelastic collisions, scattered at small angles and moving

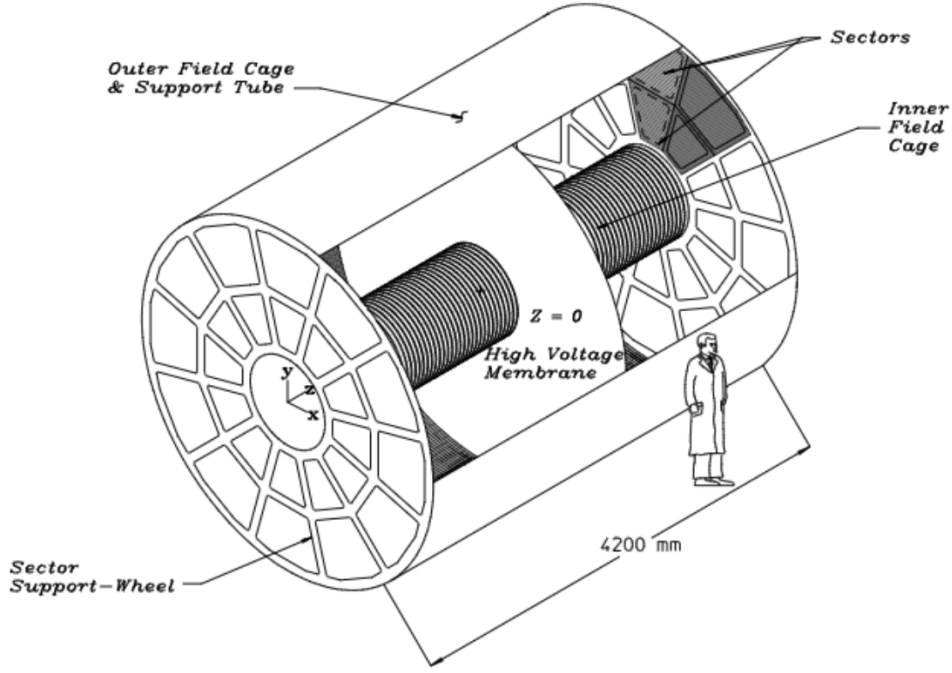


Fig. 3.3: The schematic view of the TPC. A person is shown for comparison. Taken from Ref. [27].

inside the accelerator beam pipe. Its main advantage is capability to measure scattered particles close to the beam at a distance of few millimetres from the beam with precision about $30 \mu\text{m}$ [31,32].

The RPs at the STAR experiment is system of forward detectors located on each side of the IP. There are two RP stations at a distance of 15.8 m and 17.6 m. Each station consists of two RP vessels, one below and one above the beam axis. The current layout of the RP system, called Phase II*, is shown in Fig. 3.4. Each vessel houses a detector package made of a Silicon Strip Detector (SSD) package, which measures position of the scattered proton and a scintillation counter for trigger. The RP vessel and SSD package are shown in Fig. 3.5. The RP detectors are located between DX and D0 dipole magnets and they enable the reconstruction of the particle momentum by measuring the track angle and its position since the constant and uniform magnetic field of the DX magnet works as a spectrometer [32,33].

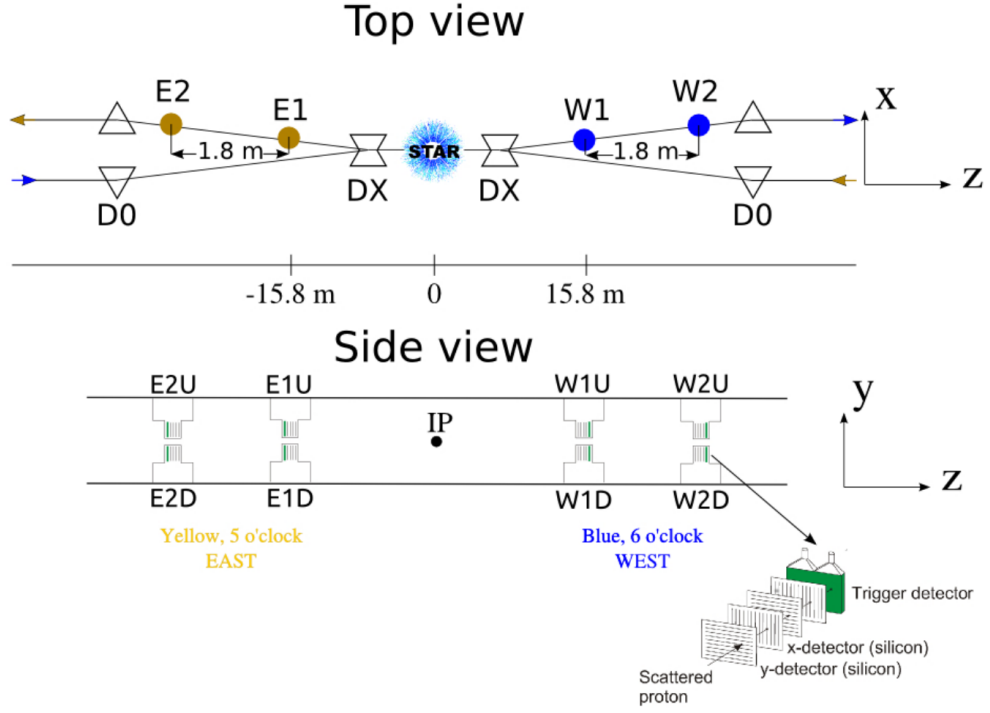


Fig. 3.4: The Roman Pot Phase II* layout. Top view with highlighted Roman Pot stations E1, E2, W1, W2 and dipole magnets DX, D0. Side view with depicted Roman Pots. Taken from Ref. [6].

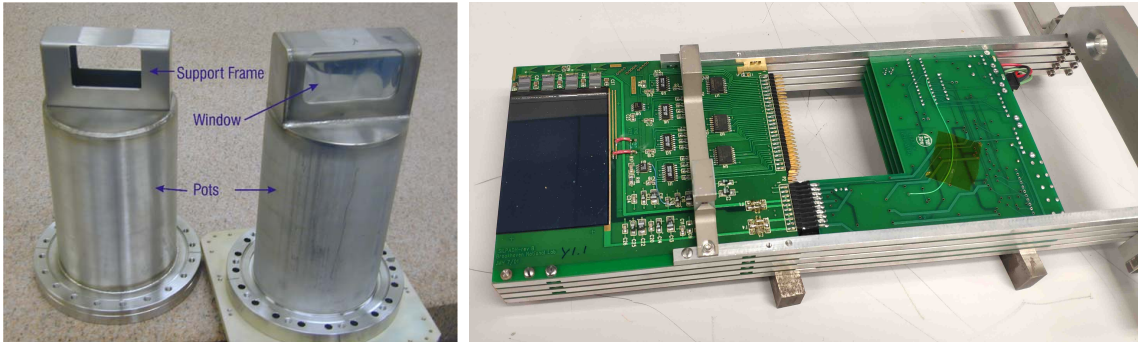


Fig. 3.5: The Roman pot vessels and the SSD package. Left: a photo of the two Roman Pots vessels, taken from Ref. [31]. Main parts are indicated. Right: a photo of the Silicon Strip Detector package, taken by myself.

Chapter 4

Analysis of CEP in pp collisions at $\sqrt{s} = 510$ GeV

In the following chapters, the analysis of the first measurement of CEP in pp collisions at $\sqrt{s} = 510$ GeV obtained by the STAR experiment is described. The analysis is focused on CEP in $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^-\pi^+\pi^-$ channels. The analysis is still in progress and the results presented in this work marked with “THIS THESIS” label have not yet been approved by the STAR collaboration for public presentation. Results with a “STAR preliminary” label were approved by the STAR collaboration and were presented at the 40th International (virtual) Conference on High Energy Physics [34].

The analysis was done using the ROOT framework and it was calculated on the STAR farm. In this and following chapters, the natural units are used.

4.1 Data preparation and quality assurance

In 2017, the STAR experiment collected data of pp collisions at $\sqrt{s} = 510$ GeV. The data were processed and stored in MuDst files, which are produced from the raw data collected during the data taking. The MuDst files contain all information about trajectories of centrally produced particles and outgoing protons. Since the MuDst files are greater than 622 TB, they are quite impractical for the analysis. For this reason, in STAR the picoDst files, containing only the most important information about the events and individual tracks, are produced, reducing the size by a factor of 100. However, the standard picoDst format is not suitable for this analysis, since it does not contain the information about outgoing protons from the Roman Pot detectors. Thus, a *star-upcDst* framework was developed to simplify analyses that are related to forward and ultra peripheral collisions (UPC) physics. My contributions were an implementation of all parts related to Roman Pot detectors and the testing the framework. Finally, I produced picoDst files, called UPC picoDst, using the developed *star-upcDst* framework. Furthermore, I validated the UPC picoDst production by comparing it with another version of picoDst files which was made by our colleagues from the AGH University of Science and Technology, Cracow, Poland.

Although, I have done the data quality assurance, the data sample that was used for this analysis does not contain all corrections, e.g. improved space charge calibration in the TPC, which was not implemented in the MuDst files. Hence, a new data reproduction is planned.

4.2 Two hadron production

In this section, the analysis on CEP of charged hadrons pairs, namely $\pi^+\pi^-$, K^+K^- and $p\bar{p}$, is described.

4.2.1 Data sample and event selection

More than 622×10^6 CEP triggers from pp collisions at $\sqrt{s} = 510$ GeV were analysed. The CEP trigger was used to ensure signals in the inelastic or in the elastic combinations in RP stations. They are shown in Fig. 4.1. Furthermore, at least two hits in the TOF were required to ensure at least two in-time tracks in the TPC and a veto on signal in the BBC was required to ensure the rapidity gaps typical of CEP events.

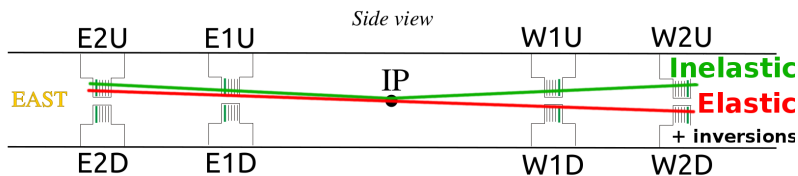


Fig. 4.1: An illustration of the inelastic and the elastic combinations of reconstructed tracks in RP stations.

First, information from RP stations was checked in the offline analysis. Since I am looking for process, where two hadrons are produced and the colliding protons stay intact, only events with one proton on each side of the interaction point were selected. Thus, only the elastic or the inelastic combinations of reconstructed tracks in RP stations were allowed. Furthermore, all eight silicon planes were required to be used in the proton reconstruction, see Fig. 4.2 (left), to ensure good quality of the proton track. In addition, the reconstructed proton is required to have transverse momenta (p_x , p_y) inside a fiducial region to ensure high geometrical acceptance. The fiducial region is defined as:

$$\begin{aligned} (p_x + 0.6 \text{ GeV})^2 + p_y^2 &< 1.25 \text{ GeV}^2, \\ 0.4 \text{ GeV} &< |p_y| < 0.8 \text{ GeV}, \\ p_x &> -0.27 \text{ GeV}, \end{aligned} \tag{4.1}$$

and it can be seen in Fig. 4.2 (right) with distributions of reconstructed protons momenta.

Second, since I am looking for two hadron production, only events with exactly two TPC tracks matched with two TOF hits originating from the same vertex were selected. Events satisfying this selection criterium plus the proton selection criteria

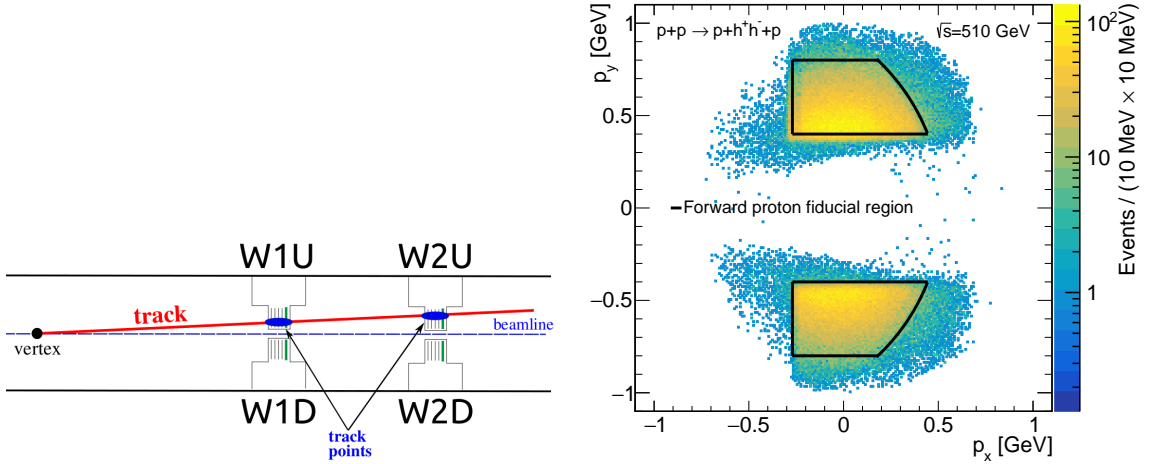


Fig. 4.2: Left: an illustration of forward proton reconstruction on the west upper branch of RP stations. Right: distributions of reconstructed protons momenta with fiducial region indicated by black lines.

are marked as "2 TPC-TOF tracks". Then, the opposite sign of two central tracks was required. Such events are marked as "Total charge 0". Finally, a cut on p_T^{miss} was applied to ensure exclusivity of the event, where the p_T^{miss} is a sum of the transverse momentum of the measured particles. For CEP processes the p_T^{miss} should be equal to zero because of the conservation of momentum. Thus, events with p_T^{miss} less than 100 MeV are called exclusive events. A comparison of three distributions of the p_T^{miss} with different selection criteria applied with indicated exclusivity cut is shown in Fig. 4.3 (left). A peak from the exclusive events can be seen in low p_T^{miss} region.

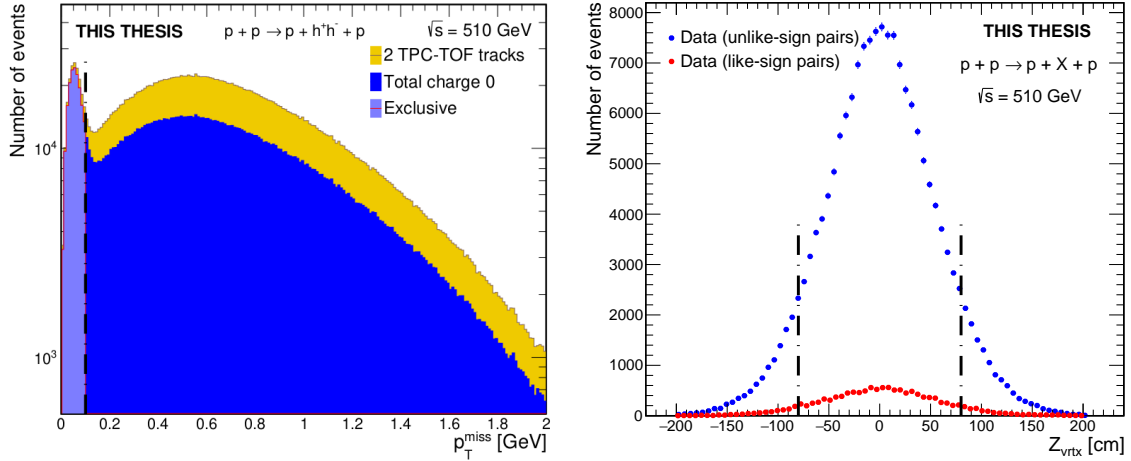


Fig. 4.3: Distributions of the sum of the transverse momentum of the measured particles p_T^{miss} for events with different selection cuts applied, described in the text (left) and the z -position of the vertex (right) for exclusive events. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties.

Furthermore, a cut on the z -position of the vertex was applied to ensure high geometrical acceptance for the central tracks in the entire fiducial phase space. The z -position of the vertex was required to be within the distance of 80 cm from the

interaction point. The distribution of the z -position of the vertex for exclusive events with applied cuts indicated by black dot-dash lines can be seen in Fig. 4.3 (right). The distribution shows the typical shape. Moreover, standard STAR good quality track criteria were applied on central tracks. Namely, a minimum of 25 hits in the TPC used for track reconstruction (N_{hits}^{fit}) to achieve good momentum resolution and a minimum of 15 hits in the TPC used to calculate dE/dx ($N_{hits}^{dE/dx}$) to ensure good dE/dx resolution, hence good PID.

Figure 4.4 shows distributions of N_{hits}^{fit} and $N_{hits}^{dE/dx}$ of central tracks for exclusive events with the additional cut on z -position of the vertex and the cut on $N_{hits}^{dE/dx}$, respectively. Distributions of N_{hits}^{fit} and $N_{hits}^{dE/dx}$ show expected shapes, although they are more symmetrical compare to another data sets measured at the STAR experiment. This discrepancy, which could be due to the space charge calibration, is a subject of further studies.

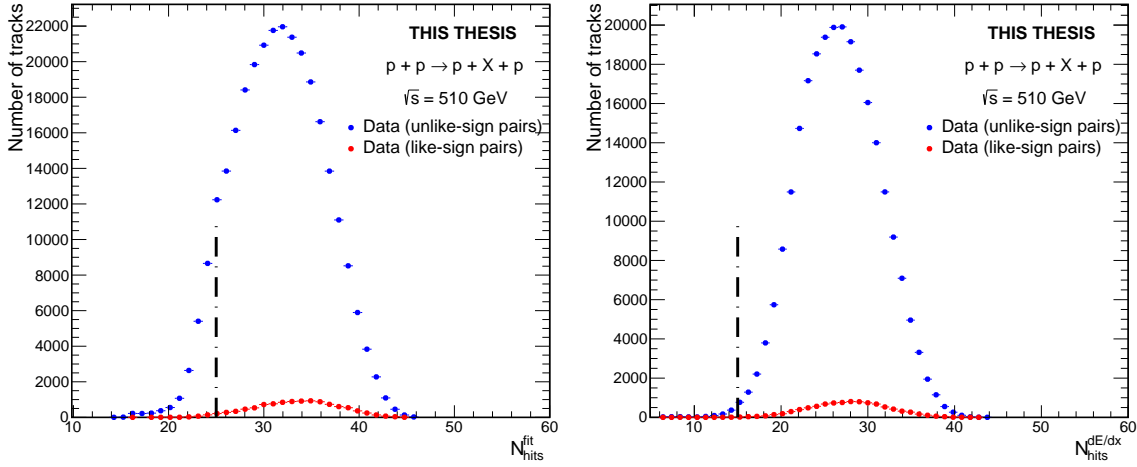


Fig. 4.4: Distributions of the number of hits in the TPC used for track reconstruction N_{hits}^{fit} (left) and the number of hits in the TPC used to calculate ionization energy loss $N_{hits}^{dE/dx}$ (right) of central tracks for exclusive events with additional cuts described in the text. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties.

Next, cuts to ensure good match between tracks and the vertex were applied. Especially, the distance of closest approach (DCA) between the central track and primary vertex in the z -direction ($DCA(z)$) had to be smaller in absolute value than 1 cm and in the transverse plane ($DCA(xy)$) smaller than 1.5 cm. Distributions of $DCA(z)$ and $DCA(xy)$ of central tracks for exclusive events with additional cuts described above can be seen in Fig. 4.5 and show the expected behaviour. Then, cut on the tracks pseudorapidity $|\eta| < 0.7$ was applied to ensure that the particles were within the TOF acceptance. The distribution of pseudorapidity of central tracks for exclusive events with additional cuts described above shows an asymmetric shape, see Fig. 4.6 (left). The discrepancy was studied and it comes from runs with TPC and TOF inefficiencies. Such inefficiency can be seen in Fig. 4.6 (right), where an averaged distributions of azimuthal angle vs. pseudorapidity of central TPC track matched with TOF for day 137 is shown.

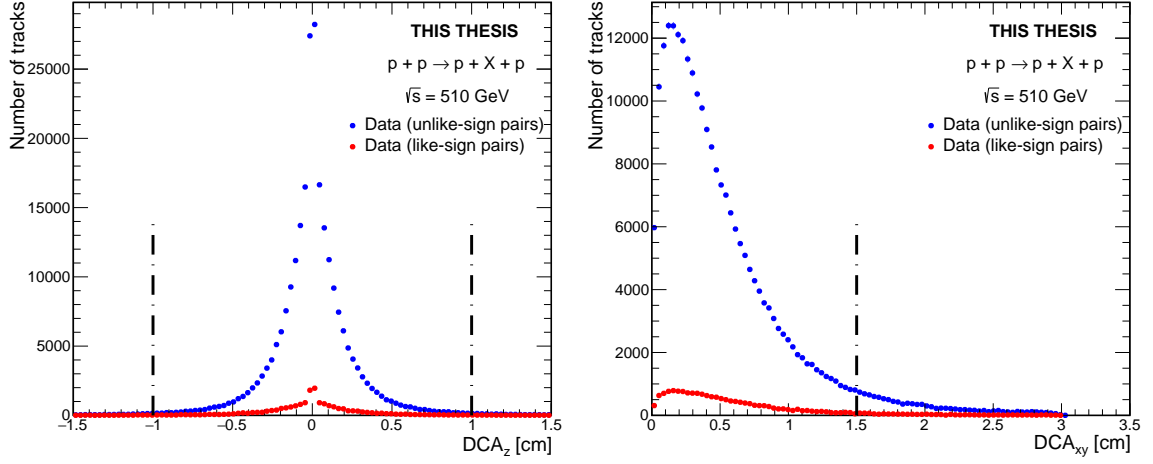


Fig. 4.5: Distributions of the distance of closest approach between the central track and primary vertex in the z -direction $DCA(z)$ (left) and in the transverse plane $DCA(xy)$ (right) for exclusive events with additional cuts described in the text. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties.

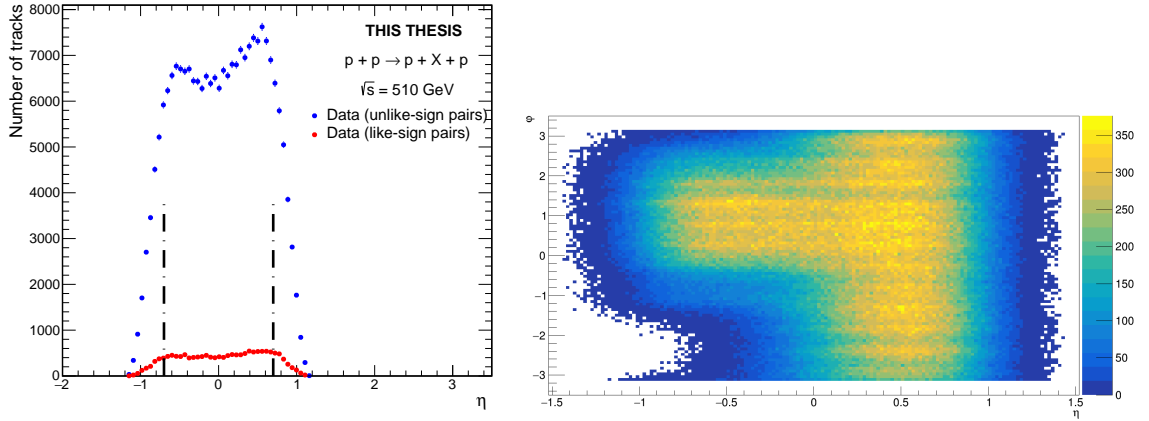


Fig. 4.6: Left: the distribution of the pseudorapidity η of central tracks for exclusive events with additional cuts described in the text. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties. Right: An averaged distributions of azimuthal angle φ vs. pseudorapidity η of central TPC track matched with TOF for day 137.

Finally, the last selection cut was applied on the four momentum transfer squared at the proton vertices t to ensure a comparable t distribution with Graniitti model that is discussed in the chapter 8. The cut was selected in such a way to minimize data loss. The cuts indicated by black dot-dash lines with the distribution for exclusive events with additional cuts described above is shown in Fig. 4.7

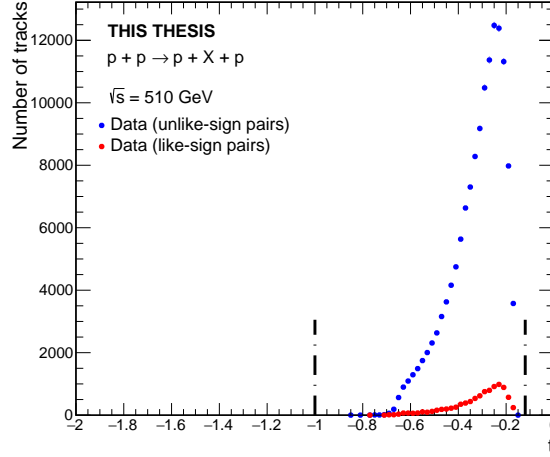


Fig. 4.7: Distributions of the four momentum transfer squared t at the proton vertices for exclusive events with additional cuts described in the text. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties.

After all the above selection criteria and PID cuts, described in the subsection 4.2.2, there are 62077 $\pi^+\pi^-$, 1697 K^+K^- and 125 $p\bar{p}$ CEP event candidates. The event selection is summarized in Fig. 4.8, where numbers of CEP event candidates remaining after the application of each event cut can be seen, and in the following list:

1. exactly two tracks in Roman Pots inside the p_x , p_y fiducial region with all eight silicon planes used in reconstruction,
2. exactly two primary TPC tracks matched with two TOF hits and originating from the same vertex,
3. total charge of central tracks equals to zero,
4. $|z\text{-position of vertex}| < 80$ cm,
5. $N_{hits}^{fit} \geq 25$,
6. $N_{hits}^{dE/dx} \geq 15$,
7. $|DCA(z)| < 1$ cm,
8. $DCA(xy) < 1.5$ cm,
9. $|\eta| < 0.7$,
10. $0.12 \text{ GeV}^2 < -t < 1.0 \text{ GeV}^2$,
11. $p_T^{miss} < 100$ MeV,
12. particle identification criteria.

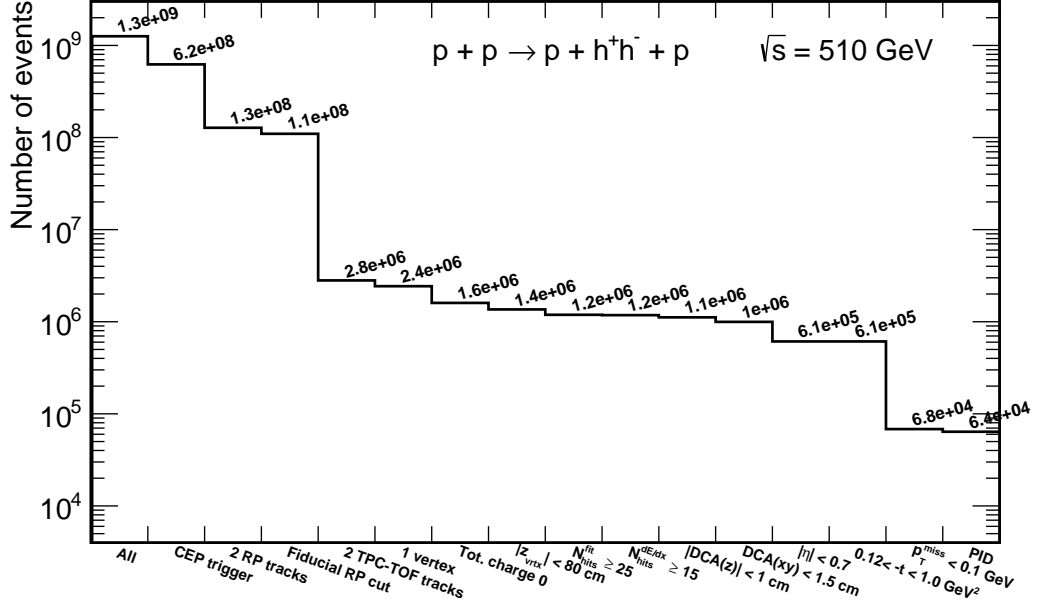


Fig. 4.8: Numbers of CEP event candidates remaining after the application of each event selection cut. The event selection cuts are described in the text.

4.2.2 Particle identification

Particle identification is based on the measurement of the ionization energy loss and the time of flight information. Since I am trying to distinguish $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs, the combined information from the TPC ($\chi_{dE/dx}^2$) is introduced and defined as:

$$\chi_{dE/dx}^2(XX) = (n\sigma_X^{\text{track1}})^2 + (n\sigma_X^{\text{track2}})^2, \quad (4.2)$$

where $n\sigma_X$ corresponds to the number of standard deviations between the measured and the theoretical energy loss for a particle type X (π, K, p) at a given momentum and superscripts track1 and track2 stand for the first and the second track in the pair.

In the CEP events, we are not able to reconstruct the time of the collision t_0 due to lack of produced forward particles. Since m_{TOF}^2 method is used. The method is based on the simple algebra and on the assumption that two central particles, produced in the same vertex, have the same mass. Thus, their masses squared are equal, $m_1^2 = m_2^2 = m_{\text{TOF}}^2$.

The energy loss of charged particles as a function of their momentum with expected values for pions, kaons, protons and deuterons is shown in Fig. 4.9 (left). Pions are dominant, as is expected in pp inelastic collisions and at RHIC energies with DIPE. The distributions of m_{TOF}^2 for $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs from exclusive events is shown in Fig. 4.9 (right). The pairs were determined only from the dE/dx information. First, the protons hypothesis was verified:

$$\chi_{dE/dx}^2(pp) < 9 \ \& \ \chi_{dE/dx}^2(KK) > 9 \ \& \ \chi_{dE/dx}^2(\pi\pi) > 9. \quad (4.3)$$

If a pair satisfied the hypothesis above, the pair was determined as $p\bar{p}$ pair. Otherwise, the kaons hypothesis was checked:

$$\chi_{dE/dx}^2(KK) < 9 \ \& \ \chi_{dE/dx}^2(pp) > 9 \ \& \ \chi_{dE/dx}^2(\pi\pi) > 9. \quad (4.4)$$

If pair was not recognized as K^+K^- or $p\bar{p}$ pair, then it was determined as $\pi^+\pi^-$ pair if:

$$\chi_{dE/dx}^2(\pi\pi) < 12. \quad (4.5)$$

Such selective PID technique was used to minimize possible misidentification. As it can be seen in Fig. 4.9, the number of $\pi^+\pi^-$ pairs is much larger than K^+K^- or $p\bar{p}$ pairs. Therefore, if a small sample of $\pi^+\pi^-$ pairs is identified as K^+K^- or $p\bar{p}$ pairs then it can produce significant false contribution to invariant mass distributions of K^+K^- or $p\bar{p}$ pairs.

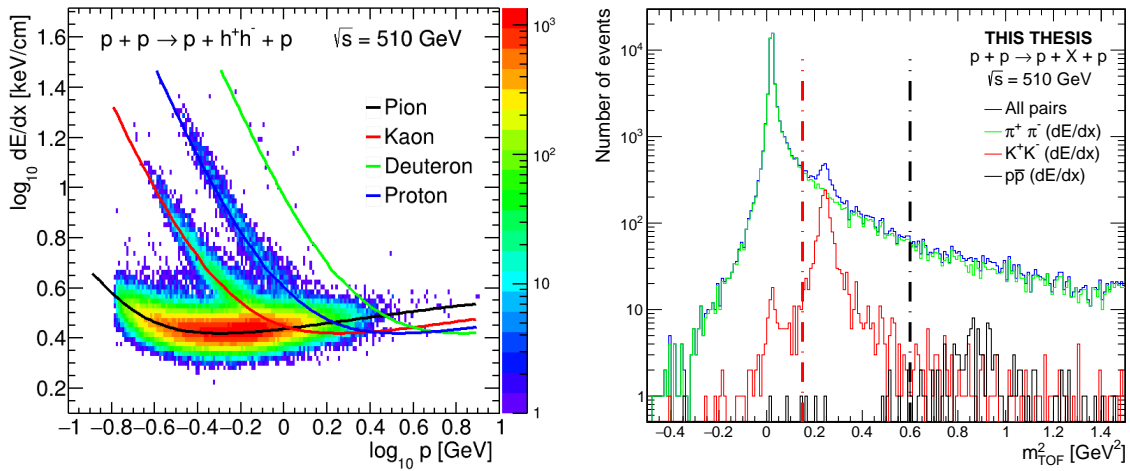


Fig. 4.9: Left: the energy loss of charged particles as a function of their momentum for exclusive events with indicated expected values for pions, kaons, protons and deuterons by coloured curves. Right: the distribution of m_{TOF}^2 for exclusive events. Distributions of m_{TOF}^2 for $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs determined just from the energy loss is shown as well. Dot-dash lines indicate m_{TOF}^2 PID cuts.

Furthermore, the m_{TOF}^2 cut was added to kaons and protons hypothesis. Kaons were required to have $m_{\text{TOF}}^2 > 0.15 \text{ GeV}^2$ and protons $m_{\text{TOF}}^2 > 0.6 \text{ GeV}^2$. These cuts are illustrated as colour dot-dash lines in Fig. 4.9 (right) and the values correspond to real mass squared for kaons $m_{\text{Kaons}}^2 \approx 0.24 \text{ GeV}^2$ and for protons $m_{\text{Protons}}^2 \approx 0.88 \text{ GeV}^2$ taking into account the detector resolution. In Fig. 4.9 (right), the peaks of kaons and protons about their real mass squared can be seen. Pions misidentified as kaons, using only the dE/dx information, can be seen as well.

In addition, cuts on transverse momenta p_T of identified central tracks were imposed: $p_T > 0.2 \text{ GeV}$ for pions, $p_T > 0.3 \text{ GeV}$ for kaons and $p_T > 0.4 \text{ GeV}$ for protons to ensure high track reconstruction efficiency. Furthermore, it was required that the lower p_T of the track in the identified pair was below 0.7 GeV for kaons and 1.1 GeV for protons to ensure high pair identification efficiency. These cuts were taken from the CEP analysis in pp collisions at $\sqrt{s} = 200 \text{ GeV}$ [14], discussed in chapter 2. An optimization of these values in pp collisions at $\sqrt{s} = 510 \text{ GeV}$ is subject of further studies.

Chapter 5

Simulations and acceptance corrections

Single particle efficiency corrections were estimated using pure STARsim, the STAR simulator based on GEANT 3. Since real data, like zero-bias events are not used, we call them acceptance corrections. I generated about one million events for each particle (π^+ , π^- , K^+ , K^- , p and \bar{p}). In each event, six particles were generated and divided into six azimuthal ranges $(0, \frac{\pi}{3})$, $(\frac{\pi}{3}, \frac{2\pi}{3})$, \dots with uniform transverse momentum distribution $0.1 < p_T < 3.5$ GeV and with uniform pseudorapidity distribution $|\eta| < 1.1$. The events were processed through STARsim with vertex spread in z coordinate described by a Gaussian function with the width $\sigma_z = 50$ cm. The width was estimated from the distribution of z -position of the vertex from the data. The distribution of z -position of the vertex of unlike-sign pairs, with like-sign background subtracted, was fitted by Gaussian function, see Fig. 5.1 (left).

The simulated events were reconstructed with the same software as the data. The acceptance corrections were obtained by dividing reconstructed tracks by MC generated tracks in each bin of a phase space. The phase space was given by z -position of the vertex, transverse momentum and pseudorapidity. The MC tracks are tracks generated by myself, as described above and they were an input to the STARsim. The reconstructed tracks are output from the STARsim. The same quality cuts as in the data analysis were imposed on the reconstructed tracks. Furthermore, a requirement of matching reconstructed track to original MC track was applied. The matching was done in (φ, η) plane using δ^2 variable defined as:

$$\delta^2 = (\varphi_{\text{MC}} - \varphi_{\text{reco}})^2 + (\eta_{\text{MC}} - \eta_{\text{reco}})^2. \quad (5.1)$$

The distribution of δ^2 and the used matching criterium $\delta^2 < 0.001$ indicated as dot-dash line can be seen in Fig. 5.1 (right).

As a quality check, the uncorrected invariant mass distribution of $\pi^+\pi^-$ pairs was compared to three corrected distributions. The first corrected distribution was obtained using efficiency corrections, obtained from embedding, from pp collisions at $\sqrt{s} = 200$ GeV [14], the second distribution was calculated using acceptance corrections described above and the last distribution was obtained using acceptance

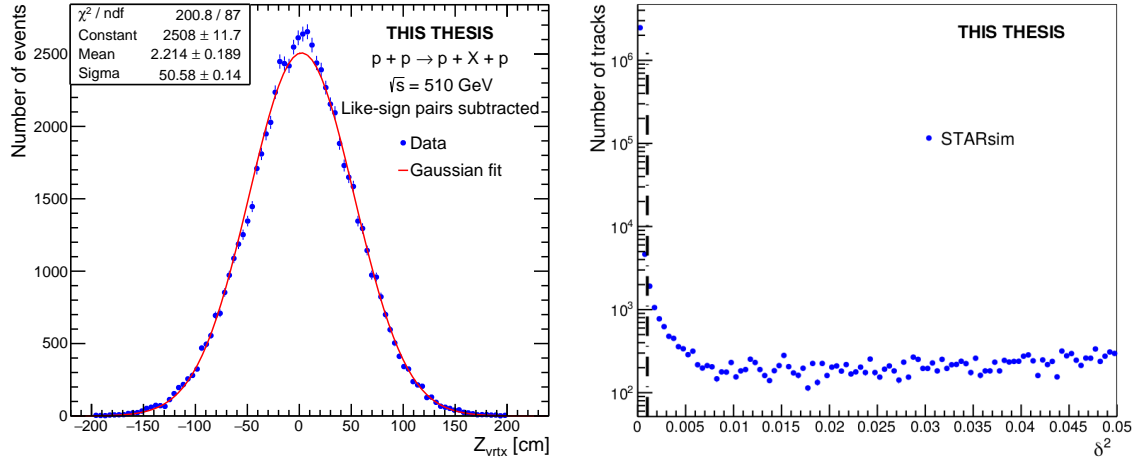


Fig. 5.1: Left: the distribution of z -position of the vertex with a Gaussian fit. Right: the distribution of δ^2 , defined in the text, for π^+ MC events. Error bars represent the statistical uncertainties.

corrections, described above, averaged over z -position of the vertex and pseudorapidity. The comparison is shown in Fig. 5.2. All distributions were normalized so the area under histogram is equal to one. Some differences can be seen, however overall shape of the distributions are the same. The efficiency corrections from pp collisions at $\sqrt{s} = 200$ GeV were not used, since there are differences between $\sqrt{s} = 200$ GeV and 510 GeV data taking, mainly bigger pile-up in 510 GeV and different tracking algorithm in the TPC.

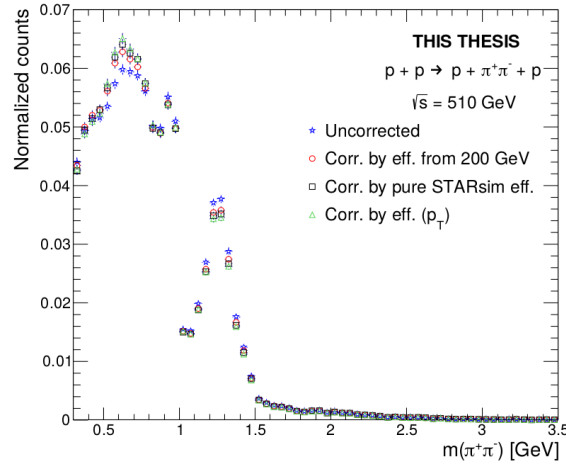


Fig. 5.2: The comparison of normalized uncorrected and corrected invariant mass of $\pi^+\pi^-$ pairs distributions for three different acceptance corrections. Error bars represent the statistical uncertainties.

Chapter 6

Results in two hadron channel

In this chapter, the results for $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs shall be shown. They include invariant mass spectra, both uncorrected and acceptance corrected, and distributions of various kinematical variables, namely the difference of azimuthal angles of the forward protons $\Delta\varphi$ and the pair rapidity y , and dependence of the invariant mass spectra on the kinematics of the forward protons. Those results shall be helpful in distinguishing between various phenomenological models.

The acceptance uncorrected and corrected invariant mass spectra of centrally exclusively produced $\pi^+\pi^-$ pairs are shown in Fig. 6.1. Both spectra show the expected features, discussed in chapter 2, a drop at about 1 GeV and a peak consistent with the $f_2(1270)$. A broad structure at about 0.6 GeV can be seen in spectra of both unlike-sign pairs and like-sign pairs, some background from a non-exclusive background based on like-sign pairs is seen. An estimate of the non-exclusive background based on the p_T^{miss} cut is subject of further studies. A more detailed analysis of the observed structures in the spectra, based on the difference of azimuthal angles of the forward protons $\Delta\varphi$ will be discussed later.

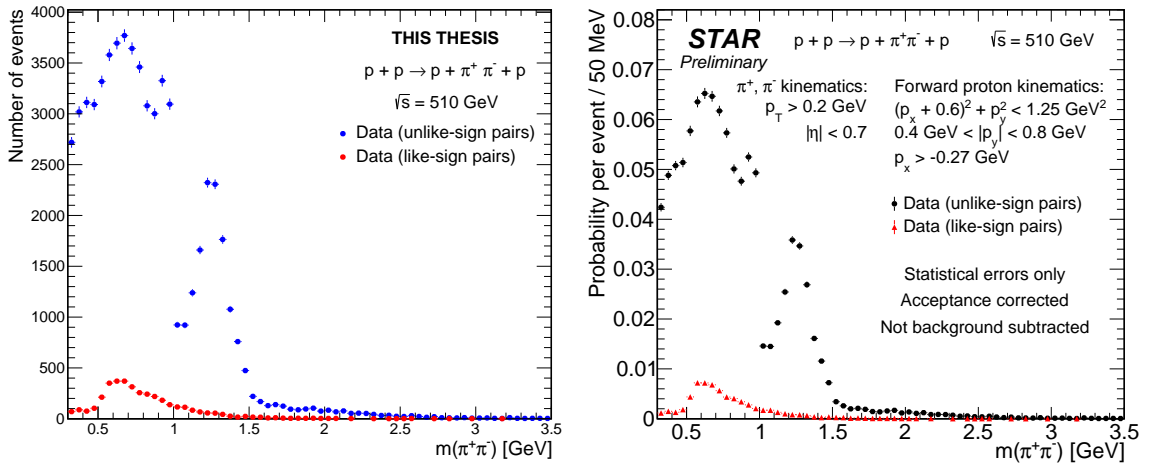


Fig. 6.1: Left: the uncorrected invariant mass spectrum of exclusively produced $\pi^+\pi^-$ pairs. Right: the acceptance corrected invariant mass spectrum of exclusively produced $\pi^+\pi^-$ pairs. Error bars represent the statistical uncertainties.

Figure 6.2 shows the acceptance uncorrected and corrected invariant mass spectra of centrally exclusively produced K^+K^- pairs. Both spectra show a very pronounced peak at above 1 GeV, a possible structure at 1.3 GeV and a peak at about 1.5 GeV consistent with the $f_2'(1525)$. The peak at 1.3 GeV could be due to the $f_2(1270)$ and the peak at 1 GeV could be due to the $\phi(1020)$. Because the latter is close to the K^+K^- mass threshold it needs more studies to make any conclusions. The peaks at 1.3 GeV and 1.5 GeV are consistent with the results from pp collisions at $\sqrt{s} = 200$ GeV, see Fig. 2.2. Because a different t -range was covered at $\sqrt{s} = 200$ GeV, a different physics might be observed at $\sqrt{s} = 510$ GeV.

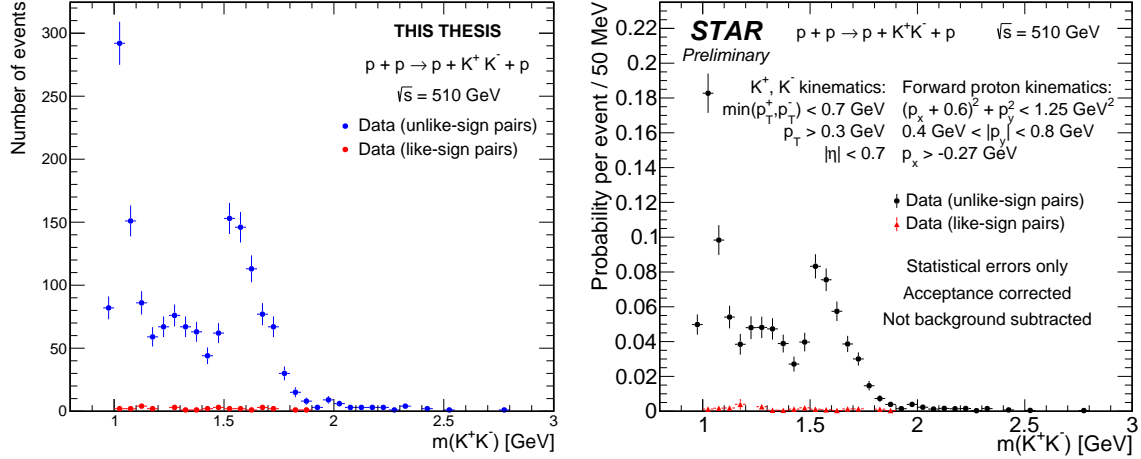


Fig. 6.2: Left: the uncorrected invariant mass spectrum of exclusively produced K^+K^- pairs. Right: the acceptance corrected invariant mass spectrum of exclusively produced K^+K^- pairs. Error bars represent the statistical uncertainties.

The invariant mass spectrum of centrally exclusively produced $p\bar{p}$ pairs for selected events can be seen in Fig. 6.3 (left). The acceptance corrected spectrum of $p\bar{p}$ pairs is shown Fig. 6.3 (right). Both spectra do not show any resonances. This observation is consistent with the measurement at $\sqrt{s} = 200$ GeV [14].

Next, we shall examine the dependence of the mass spectra on the forward proton kinematics. A commonly used variable is a difference between the azimuthal angle of the forward protons $\Delta\varphi$. Acceptance corrected distributions of $\Delta\varphi$ for selected events with exclusively produced $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs is shown in Fig. 6.4 (left). The $\Delta\varphi$ distributions show a strong suppression about 90° and peak at 0° and 180° , favouring parallel and anti parallel proton momenta. Another kinematical variable of interest is the rapidity y of the produced pair, see Fig. 6.4 (right). The y distributions show wide peaks at zero rapidity. All distributions show similar shape as measured at $\sqrt{s} = 200$ GeV, see Fig. 2.3.

As shown in chapter 2, the invariant mass distributions at $\sqrt{s} = 200$ GeV were strongly affected by the choice of the $\Delta\varphi$ region. Hence, the acceptance corrected invariant mass distributions of centrally exclusively produced $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs were differentiated into two $\Delta\varphi$ regions: $\Delta\varphi < 90^\circ$ and $\Delta\varphi > 90^\circ$ and are shown in Figs. 6.5 and 6.6. The distributions show an enhancement at low invariant masses in the $\Delta\varphi < 90^\circ$ compared to $\Delta\varphi > 90^\circ$ region. This could be due to

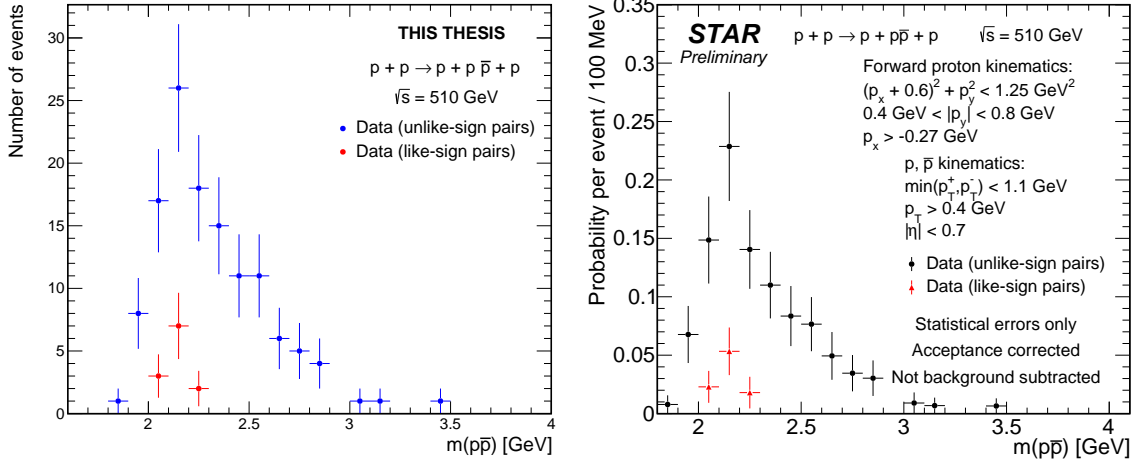


Fig. 6.3: Left: the uncorrected invariant mass spectrum of exclusively produced $p\bar{p}$ pairs. Right: the acceptance corrected invariant mass spectrum of exclusively produced $p\bar{p}$ pairs. Error bars represent the statistical uncertainties.

different Pomeron dynamics and due to a recoil, created by forward protons with $\Delta\varphi < 90^\circ$, to the central system allowing measurement at low invariant mass region.

The invariant mass spectrum of $\pi^+\pi^-$ pairs within $\Delta\varphi < 90^\circ$ shows a suppressed peak consistent with the $f_2(1270)$ compared to the $\Delta\varphi > 90^\circ$. The suppression is consistent with the results discussed in chapter 2, see Fig. 2.4. Furthermore, there is no resonances at about 0.6 GeV for $\Delta\varphi < 90^\circ$, thus the broad structure seen in Fig. 6.1 at about 0.6 GeV could be due to superposition of two shifted continuum production.

In the invariant mass spectrum of K^+K^- pairs, the peaks at 1.3 GeV and at about 1.5 GeV consistent with the $f_2(1270)$ and the $f_2'(1525)$, respectively, can be seen only in the $\Delta\varphi > 90^\circ$. In addition the peak at about 1 GeV is totally suppressed in the $\Delta\varphi > 90^\circ$. The results show much stronger dependence of the invariant mass of K^+K^- pairs on the choice of the $\Delta\varphi$ region compared to the results at $\sqrt{s} = 200$ GeV, see Fig. 2.4.

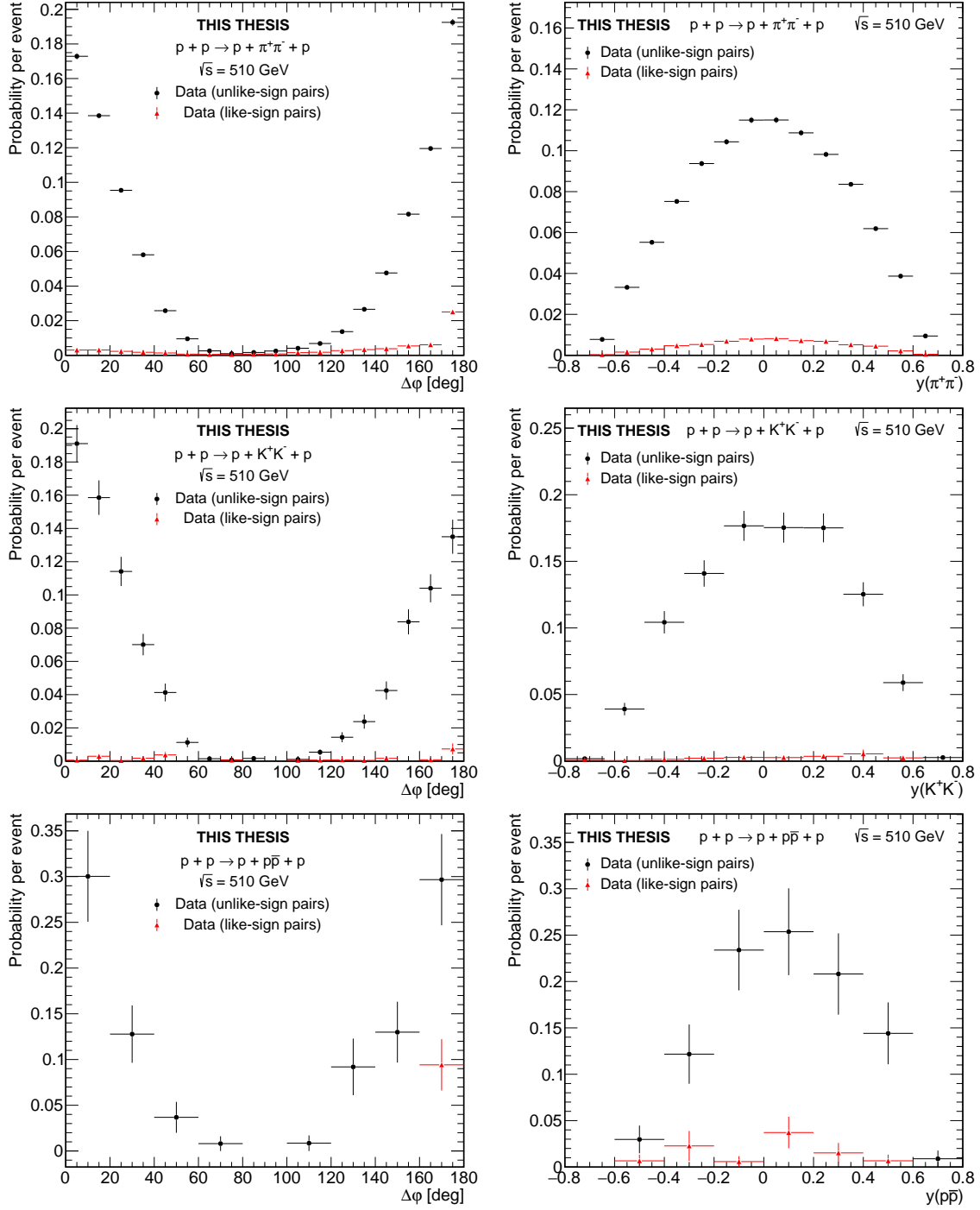


Fig. 6.4: Acceptance corrected distributions of the difference of azimuthal angles of the forward protons $\Delta\varphi$ (left) and the pair rapidity (right) for events with exclusively produced $\pi^+\pi^-$ (upper row), K^+K^- (middle row) and $p\bar{p}$ (bottom row) pairs. Error bars represent the statistical uncertainties.

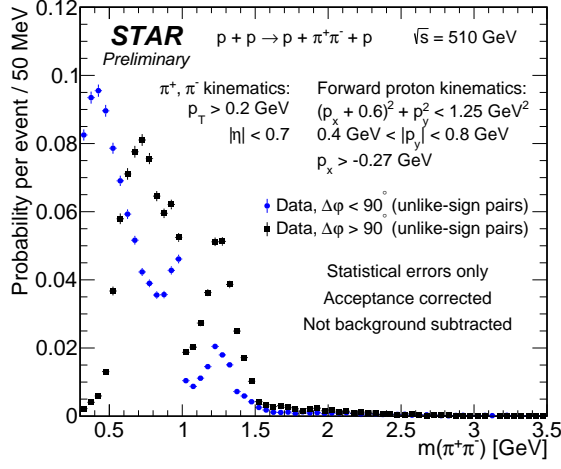


Fig. 6.5: Acceptance corrected invariant mass spectra of exclusively produced $\pi^+\pi^-$ pairs in two regions of the difference of azimuthal angles of the forward protons: $\Delta\varphi < 90^\circ$ and $\Delta\varphi > 90^\circ$. Error bars represent the statistical uncertainties.

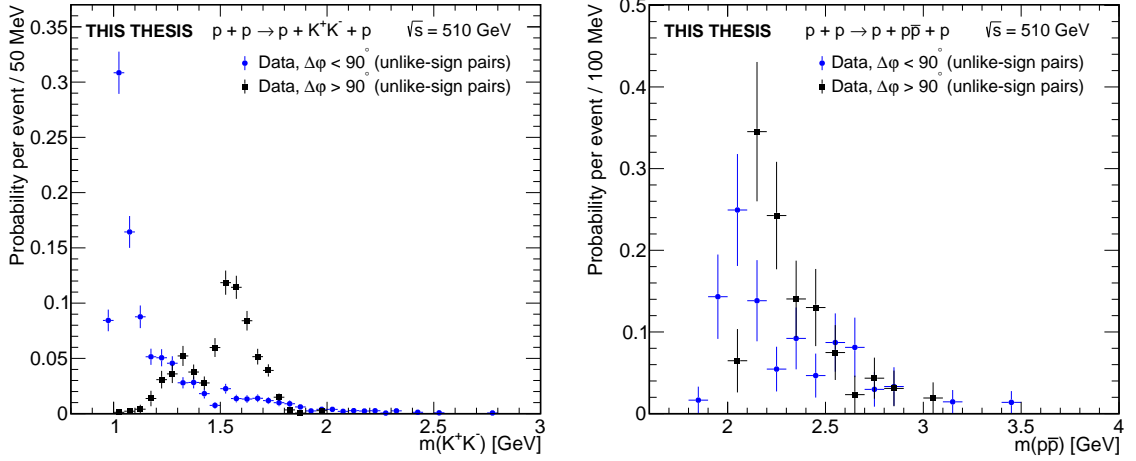


Fig. 6.6: Acceptance corrected invariant mass spectra of exclusively produced K^+K^- (left) and $p\bar{p}$ (right) pairs in two regions of the difference of azimuthal angles of the forward protons: $\Delta\varphi < 90^\circ$ and $\Delta\varphi > 90^\circ$. Error bars represent the statistical uncertainties.

Chapter 7

$\pi^+\pi^-\pi^+\pi^-$ production

The same data sample was used to examine the four-pion channel, the $\pi^+\pi^-\pi^+\pi^-$ production. More than 622×10^6 CEP triggers from pp collisions at $\sqrt{s} = 510$ GeV were analysed. The same event selection as in the two hadron production was applied, except four primary TPC tracks matched with four TOF hits were required instead of two, since I am looking for the $\pi^+\pi^-\pi^+\pi^-$ channel. The requirement of four TOF hits was used to ensure four in-time tracks in the TPC, therefore to eliminate pile-up tracks, since the TOF is a much faster detector than the TPC.

The summary of the event selection is shown in Fig. 7.1, where numbers of CEP event candidates remaining after the application of each event cut can be seen. In the four-hadron channel, the production of $\pi^+\pi^-\pi^+\pi^-$ is expected to be dominant and in addition pions are dominant in two hadron production. Therefore, only a $n\sigma_\pi < 3$ cut for PID was used for each track. Thus, $p_T > 0.2$ GeV cut was imposed for all four particles to ensure high track reconstruction efficiency. After all the above selection cuts, there are 758 $\pi^+\pi^-\pi^+\pi^-$ CEP event candidates.

Figure 7.2 (left) shows three distributions of p_T^{miss} for $\pi^+\pi^-\pi^+\pi^-$ event candidates with different selection cuts applied. The first distribution marked as “4 TPC-TOF” tracks satisfies following criteria:

1. exactly two tracks in Roman Pots inside the p_x , p_y fiducial region with all eight silicon planes used in reconstruction,
2. exactly four primary TPC tracks matched with four TOF hits and originating from the same vertex.

The second distribution marked as “Total charge 0” satisfies the two criteria above and in addition a total charge of central particles is required to equal to zero. Finally, the last distribution marked as “Exclusive” has to meet the three criteria above plus the exclusivity cut $p_T^{miss} < 100$ MeV. The distribution shows a sign of a peak from exclusive events, although it is very suppressed compared to the two hadron production. The requirement of 4 hits in TOF reduces a phase space, where particles can be produced and so it suppresses the statistic.

The acceptance uncorrected and corrected invariant mass spectra of centrally exclusively produced $\pi^+\pi^-\pi^+\pi^-$ pairs for selected events are shown in Fig. 7.3. The

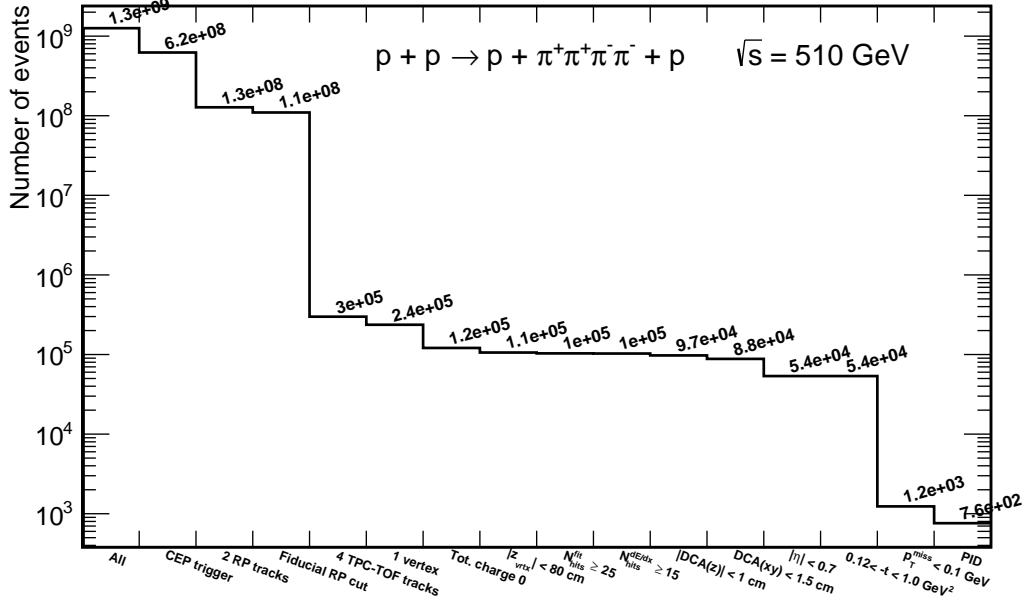


Fig. 7.1: Numbers of CEP event candidates remaining after the application of each event cut for $\pi^+\pi^-\pi^+\pi^-$ channel.

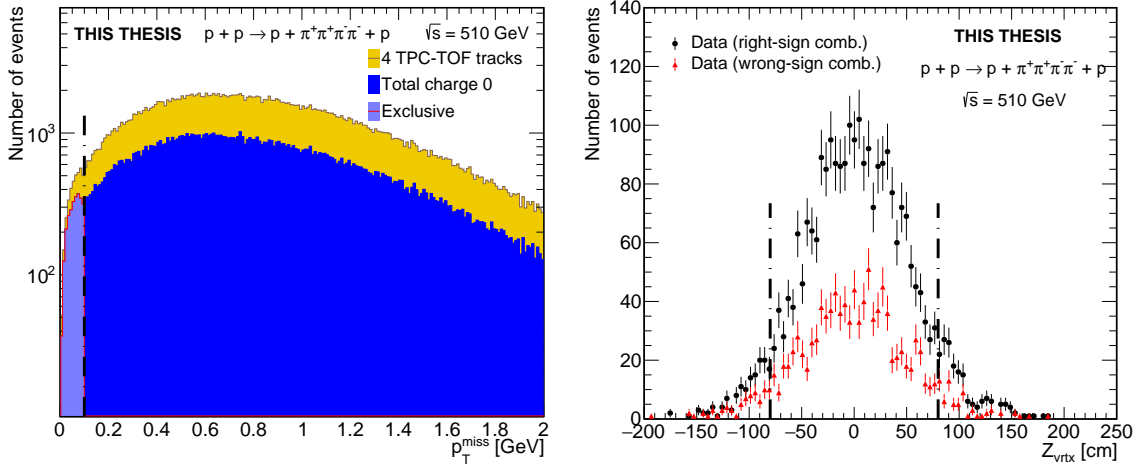


Fig. 7.2: Distributions of the sum of the transverse momentum of the measured particles p_T^{miss} for events with different selection cuts applied, described in the text (left) and z -position of the vertex (right) for exclusive $\pi^+\pi^-\pi^+\pi^-$ events. Black dot-dash lines indicate the applied cut. Error bars represent the statistical uncertainties.

corrected spectra shows a rise at about 1.3 GeV and drop at about 2.2 GeV, these features are consistent with the measurement done by WA102 Collaboration [35]. Further studies need to be done to make more detailed conclusions.

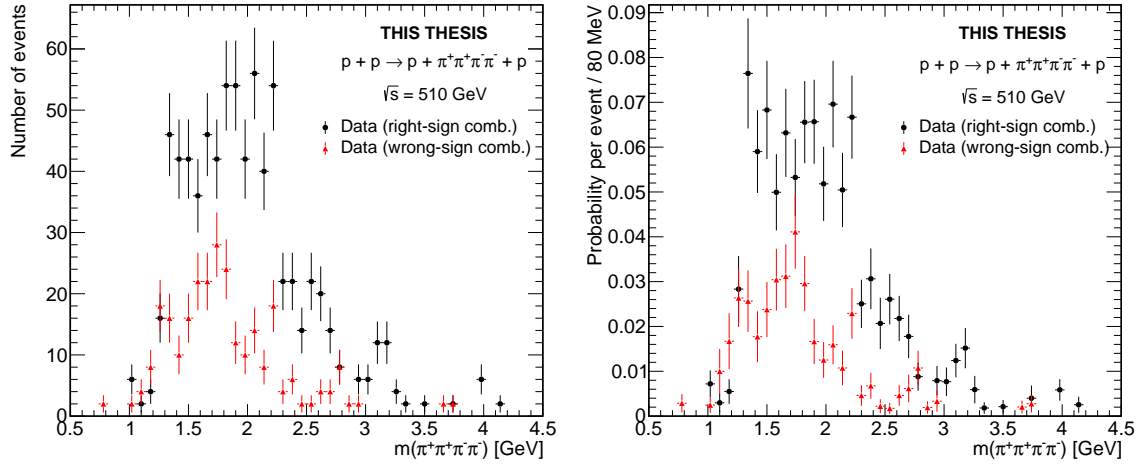


Fig. 7.3: Left: the uncorrected invariant mass spectrum of exclusively produced $\pi^+\pi^-\pi^+\pi^-$. Right: the acceptance corrected invariant mass spectrum of exclusively produced $\pi^+\pi^-\pi^+\pi^-$. Error bars represent the statistical uncertainties.

Chapter 8

Comparison with Graniitti Monte Carlo model

I compare results with Graniitti Monte Carlo event generator [36], which is a MC generator designed for high energy diffraction physics. Unlike the MC generator discussed in chapter 2, Graniitti combines a parametrized resonance production with a continuum production. Furthermore, it combines forward and central spin correlations together with event-by-event eikonal screening loop and forward proton excitation kinematics.

Graniitti is focused on the low-mass domain of CEP processes using the S -matrix and where the glueballs are expected. In addition to the Pomeron-Pomeron interaction, Graniitti also includes the photon-photon and the photon-Pomeron interactions.

I simulated 10^5 CEP events for each channel presented in this thesis. Namely, I simulated $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^-\pi^+\pi^-$ using Graniitti with continuum and resonances production in pp collisions at $\sqrt{s} = 510$ GeV. The following resonances were used in the model: the $f_0(980)$, the $\phi(1020)$, the $f_2(1270)$, the $f_0(1500)$, the $f'_2(1525)$ and the $f_0(1710)$. The same kinematic cuts to the MC events were applied as in the analysis. Since Graniitti predicts pure signal, the background in the data was estimated from like-sign pairs distributions. Then, it was subtracted from unlike-sign distributions. Obtained distributions were normalized so the area under the histogram was equal to one and compared with the Graniitti normalized distributions.

The acceptance corrected invariant mass spectrum of exclusively produced $\pi^+\pi^-$ with subtracted like-sign background compared to the Graniitti output can be seen in Fig. 8.1 (left). Although differences can be seen, the main features are described, the drop at about 1 GeV and peak consistent with the $f_2(1270)$. Figure 8.1 (right) shows a comparison for K^+K^- channel. Graniitti describes all three peaks, that can be assigned to $\phi(1020)$, $f_2(1270)$ and $f'_2(1525)$.

The Graniitti output compared to the invariant mass spectrum of exclusively produced $p\bar{p}$ pairs with like-sign background subtracted is shown in Fig. 8.2 (left). Differences are seen, however overall shape of the distributions agree quite well.

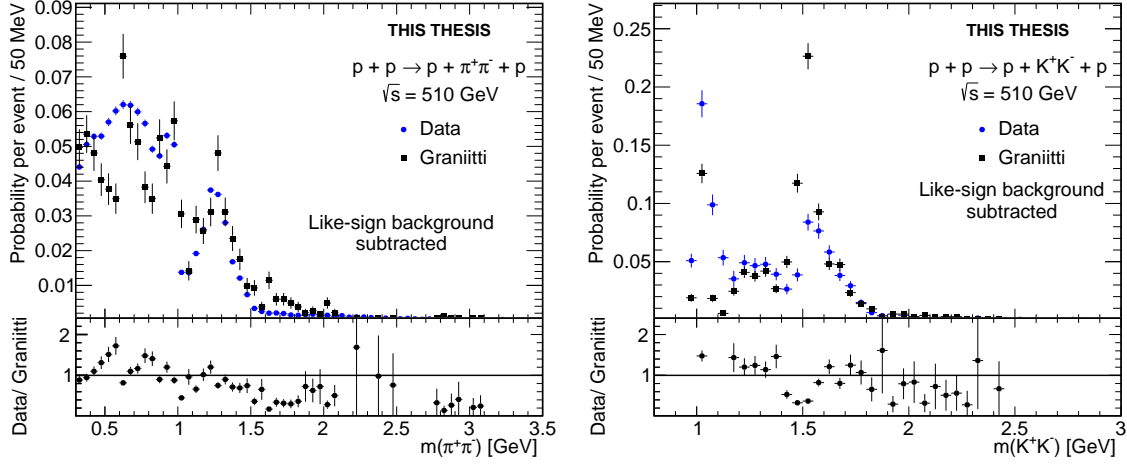


Fig. 8.1: The acceptance corrected invariant mass spectra of exclusively produced $\pi^+\pi^-$ (left) and K^+K^- (right) compared to Graniitti. The like-sign background was subtracted. Error bars represent the statistical uncertainties.

Comparison for $\pi^+\pi^-\pi^+\pi^-$ channel can be seen in Fig. 8.2 (right). A discrepancy at the low mass can be seen, nevertheless the data shows large statistical errors and more studies need to be done to make any conclusions.

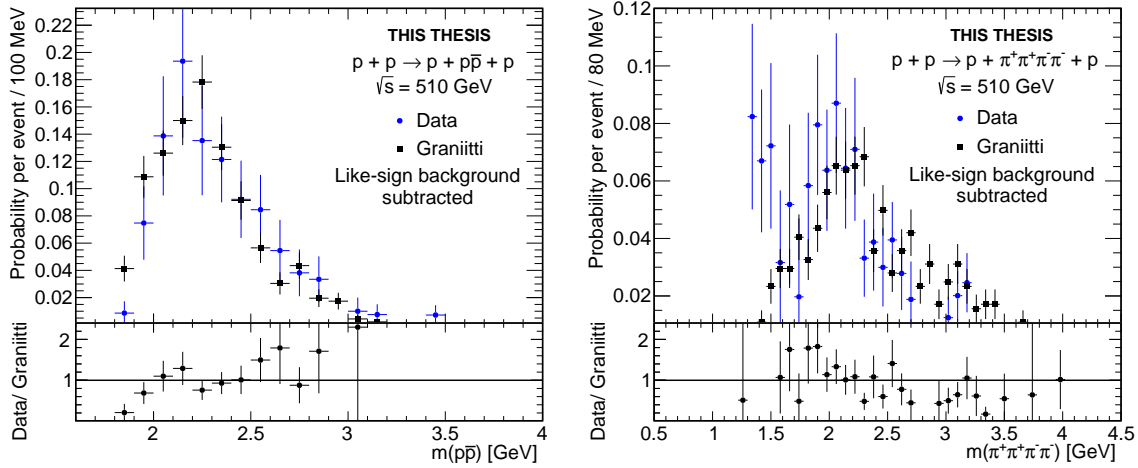


Fig. 8.2: The acceptance corrected invariant mass spectra of exclusively produced $p\bar{p}$ (left) and $\pi^+\pi^-\pi^+\pi^-$ (right) compared to Graniitti. The like-sign background was subtracted. Error bars represent the statistical uncertainties.

Comparisons of the difference of azimuthal angles of the forward protons $\Delta\varphi$ for $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^-\pi^+\pi^-$ channels is shown in Figs. 8.3 and 8.4. The overall shape of the distributions is described quite well. Acceptance corrected distributions of the pair rapidity compared to Graniitti can be seen in Figs. 8.5 and 8.6. Discrepancies can be seen and they are subject of further studies.

As we discussed earlier, the choice of $\Delta\varphi$ region can strongly affect the invariant mass distribution. Therefore acceptance corrected invariant mass spectra of exclusively produced $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs compared to Graniitti predictions were examined in two $\Delta\varphi$ regions. The comparison can be seen in Fig. 8.7.

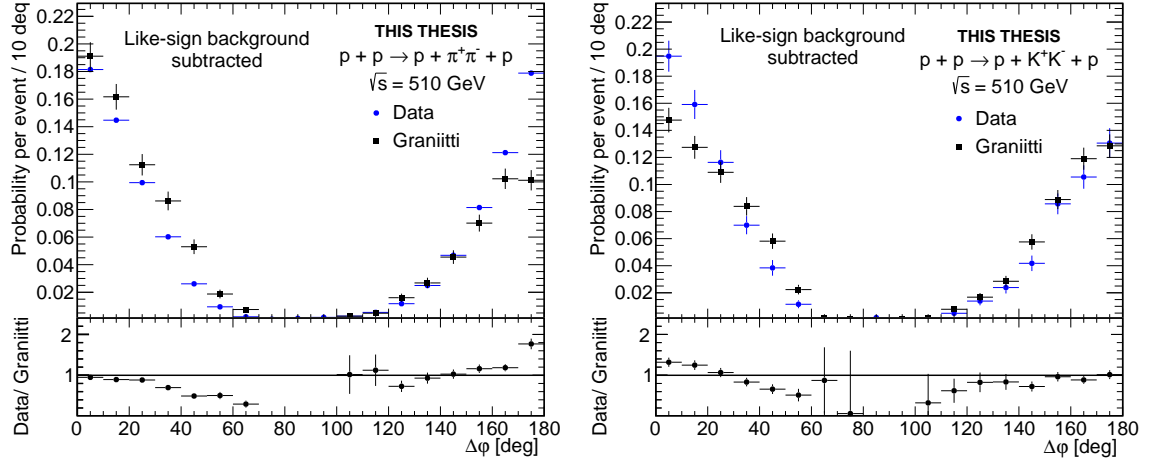


Fig. 8.3: Acceptance corrected distributions of the difference of azimuthal angles of the forward protons $\Delta\varphi$ for events with exclusively produced $\pi^+\pi^-$ (left) and K^+K^- (right) events compared to Graniitti. The like-sign background was subtracted. Error bars represent the statistical uncertainties.

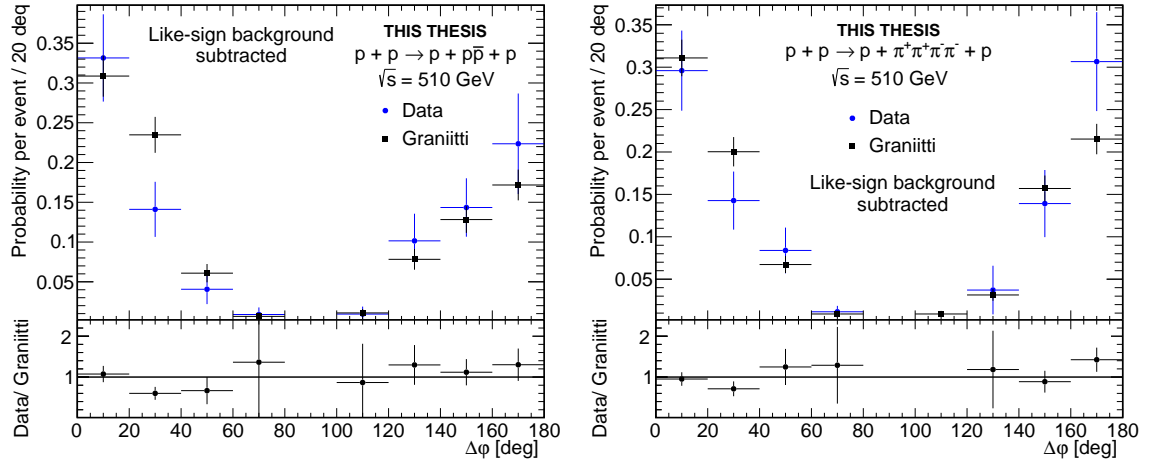


Fig. 8.4: Acceptance corrected distributions of the difference of azimuthal angles of the forward protons $\Delta\varphi$ for events with exclusively produced $p\bar{p}$ (left) and $\pi^+\pi^-\pi^+\pi^-$ (right) events compared to Graniitti. The like-sign background was subtracted. Error bars represent the statistical uncertainties.

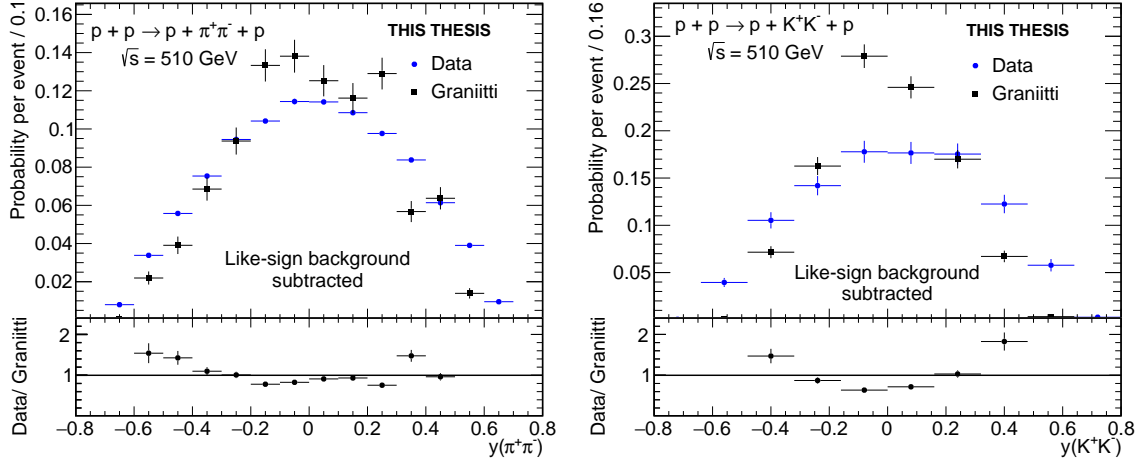


Fig. 8.5: Acceptance corrected distributions of the pair rapidity for events with exclusively produced $\pi^+\pi^-$ (left) and K^+K^- (right) events compared to Graniitti. The like-sign was background subtracted. Error bars represent the statistical uncertainties.

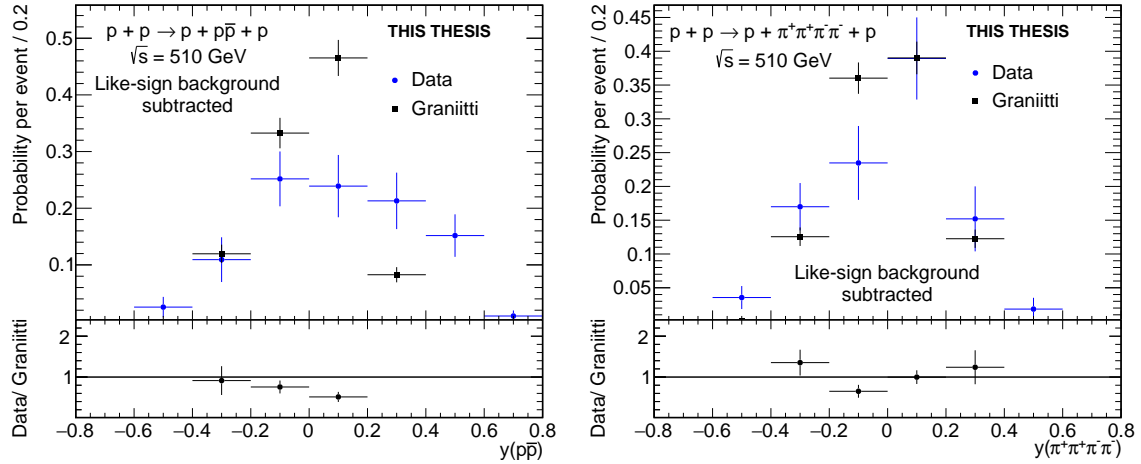


Fig. 8.6: Acceptance corrected distributions of the pair rapidity for events with exclusively produced $p\bar{p}$ (left) and $\pi^+\pi^-\pi^+\pi^-$ (right) events compared to Graniitti. The like-sign background was subtracted. Error bars represent the statistical uncertainties.

The invariant mass spectrum of $\pi^+\pi^-$ pairs in $\Delta\varphi > 90^\circ$ is described well by the Graniitti prediction up to 1.5 GeV, especially the rise at above 0.5 GeV followed by the drop at 0.7 GeV with the peak and sharp drop at about 0.9 GeV and the peak consistent with the $f_2(1270)$. In $\Delta\varphi < 90^\circ$, Graniitti is able to predict overall shape of the invariant mass spectrum, in particular the rise in the beginning of the spectra followed by the sharp drop at 0.4 GeV. Furthermore, the peak at 0.9 GeV followed by the rapid drop and the peak about 1.2 GeV can be seen as well. However, Graniitti predicts also peaks at 0.6 and 1.6 GeV, which are not seen in the data. Those need to be examined further, as they can be due to interference effects and fiducial volume cuts effects.

The invariant mass distribution of K^+K^- pairs in $\Delta\varphi > 90^\circ$ is really well described by the Graniitti prediction. First, the rise at 1.1 GeV followed by the broad structure at about 1.3 GeV is seen. Then, the peak consistent with the $f_2'(1525)$ is clearly visible, although it should be noted that Graniitti predicts more pronounced peak than it is observed in the data. For $\Delta\varphi < 90^\circ$, Graniitti is able to describe the data up to 1.4 GeV, especially the peak at about 1 GeV followed by the broad structure at 1.2 GeV. Nevertheless Graniitti predicts the peak at about 1.5 GeV, which is not seen in the data. The invariant mass spectra of $p\bar{p}$ pairs are roughly described by Graniitti in both $\Delta\varphi$ regions.

Graniitti is the first MC model that includes both the continuum and the resonance production. Although it was tuned to LHC energies, it is able to described the data from pp collisions at $\sqrt{s} = 510$ GeV quite well. Some discrepancies between Graniitti predictions and the data are seen. Therefore, I hope and believe that Graniitti will be tuned to the newest results measured at the STAR experiment [14], since Graniitti predictions are very promising.

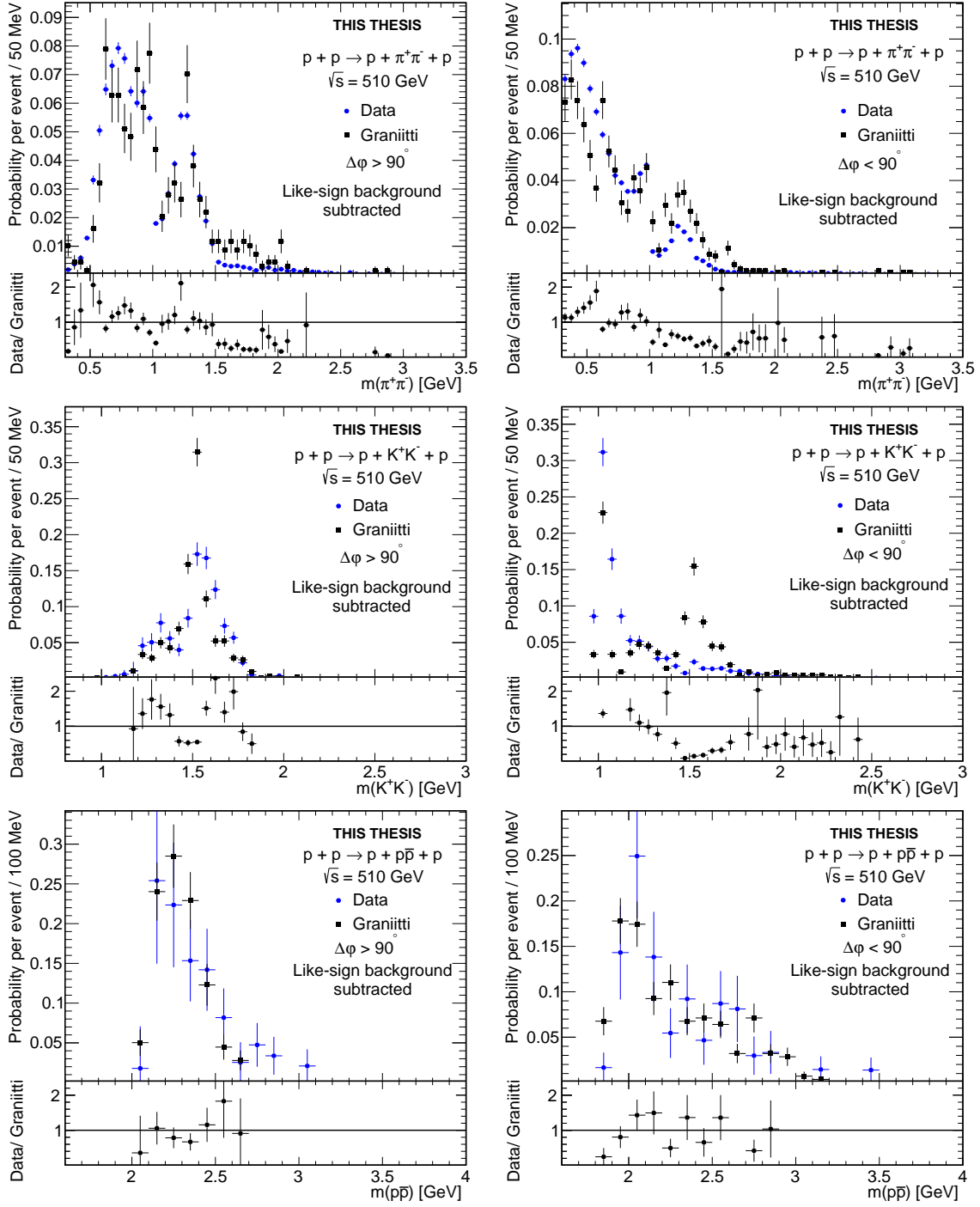


Fig. 8.7: Acceptance corrected invariant mass spectra of exclusively produced $\pi^+\pi^-$ (upper row), K^+K^- (middle row) and $p\bar{p}$ (bottom row) pairs compared to Graniitti and differentiated in two regions of the difference of azimuthal angles of the forward protons: $\Delta\varphi > 90^\circ$ (left) and $\Delta\varphi < 90^\circ$ (right). The like-sign background was subtracted. Error bars represent the statistical uncertainties.

Summary

In this thesis I studied the CEP of $\pi^+\pi^-$, K^+K^- , $p\bar{p}$ and $\pi^+\pi^-\pi^+\pi^-$ with the STAR detector at RHIC. The data from pp collisions at $\sqrt{s} = 510$ GeV collected by the STAR experiment in 2017 were analysed. Since the data contain triggers from DIPE in CEP, they are very promising for glueball search. The experimental confirmation of the existence of the glueball would be yet another strong support for the validity of the QCD theory.

The invariant mass spectra of the CEP $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs confirmed features seen in previous measurements, however an interesting feature as the peak at about 1 GeV in the K^+K^- spectrum is also seen and requires more investigation. The results are encouraging and there are further ongoing studies as: an efficiency study of Roman Pot detectors, new data reproduction with the improved space charge calibration in the TPC, using embedding to calculate efficiencies which will include the pile-up effects, and an estimate of the non-exclusive background based on the p_T^{miss} cut. The new MC event generator, Graniitti, was compared to the results and it gave promising predictions.

Some of the results of this work has been presented by the author at the 40th International (virtual) Conference on High Energy Physics in July 2020. Namely, the acceptance corrected invariant mass spectra of centrally exclusively produced $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs and the acceptance corrected invariant mass spectra of exclusively produced $\pi^+\pi^-$ pairs in two regions of the difference of azimuthal angles of the forward protons: $\Delta\varphi < 90^\circ$ and $\Delta\varphi > 90^\circ$.

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